


# Gottesman-Knill theorem

- a  $n$ -qubit state is a vector that has  $2^n$  components
- its unitary evolution can be described by  $2^n \times 2^n$  unitary matrices


 exponentially large memory and extremely fast processors are required for large  $n$

- How many qubits can we follow on a computer with a 16GB memory?



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$$2 \cdot 32 \cdot (2^n)^2 < 16 \times 10^9 \times 8 \quad \therefore 2^{2n-1} < 10^9 \quad 1 \text{ byte} = 8 \text{ bit}$$

$$\therefore (2n - 1) \log_{10} 2 < 9 \quad \therefore n \leq 15 \quad \log_{10} 2 = 0.30$$

it is hard to simulate a large quantum system on a classical computer



# Gottesman-Knill theorem

- In the Heisenberg representation,  $n$ -qubit quantum circuits composed of Clifford gates can be described by following the evolution of at most  $2n$  logical operators  $\{\bar{X}_1, \dots, \bar{X}_n, \bar{Z}_1, \dots, \bar{Z}_n\}$
- since each logical operator is a Pauli product with  $\pm$ , it can be described by  $2n + 1$  bits  $(\pm(i^{\alpha_1\beta_1}X_1^{\alpha_1}Z_1^{\beta_1})\dots(i^{\alpha_n\beta_n}X_n^{\alpha_n}Z_n^{\beta_n}), \alpha_i, \beta_i \in \{0,1\})$

 the logical operators require at most  $2n(2n + 1)$  bits

Clifford circuits can be simulated on a classical computer in polynomial time of  $n$



# Gottesman-Knill theorem

Any quantum computer performing only Clifford group gates can be perfectly simulated in polynomial time on a classical computer.

- Clifford group gates do not provide a universal set of gates, e.g.,  $T$  is outside the Clifford group
- the theorem implies that quantum computation is only more powerful than classical computation when it uses gates outside the Clifford group



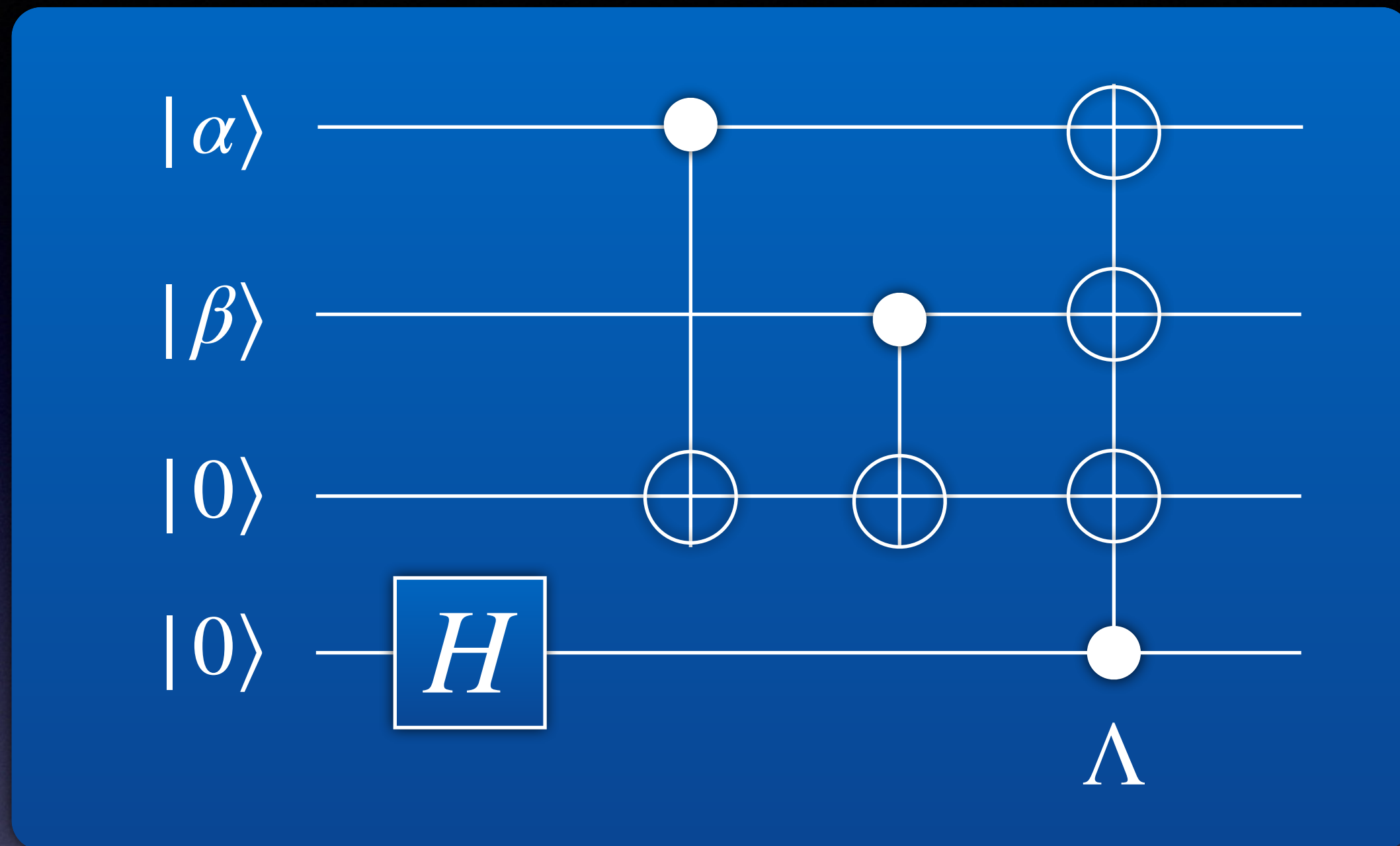
# Gottesman-Knill theorem

Any quantum computer performing only Clifford group gates can be perfectly simulated in polynomial time on a classical computer.

- however circuits using only Clifford group gates also have a number of important applications in the area of quantum communications - quantum error-correcting codes, quantum teleportation,...



# Quantum Error-Correcting codes



start  
*IIZI*    *IIZI*    *ZIZI*    *ZZZI*  
*IIIZ*    *IIIX*    *IIIX*    *IIIX*  
 stabilizers

*ZZZZ*  
*XXXX*

Exercise

Show

$$\Lambda(I \otimes I \otimes I \otimes X)\Lambda = X \otimes X \otimes X \otimes X$$

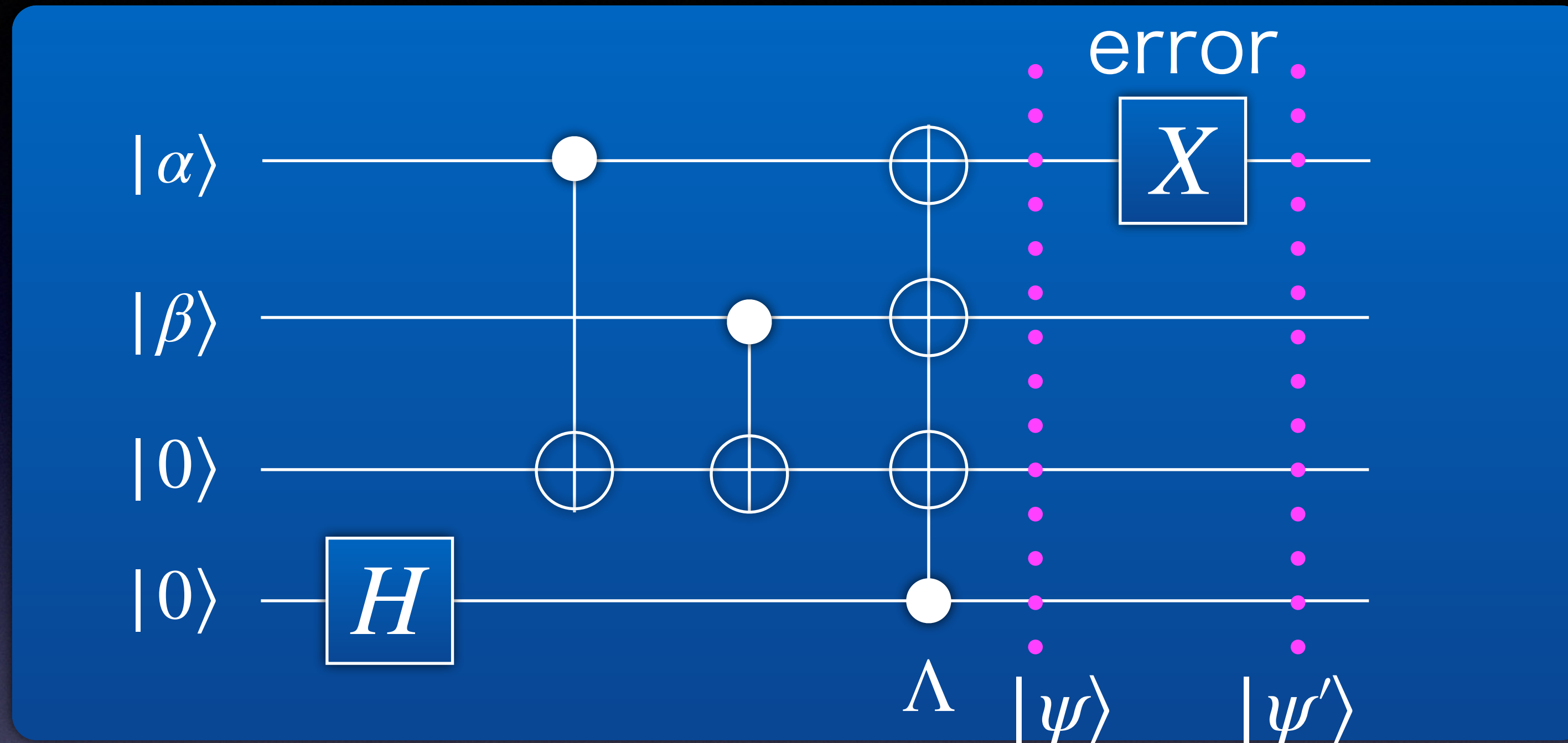
$$\Lambda(Z \otimes Z \otimes Z \otimes I)\Lambda = Z \otimes Z \otimes Z \otimes Z$$

using

$$\Lambda = I \otimes I \otimes I \otimes |0\rangle\langle 0| + X \otimes X \otimes X \otimes |1\rangle\langle 1|$$



# Quantum Error-Correcting codes



$\Lambda$	$X$
$XXXX$	$XXXX$
$ZZZZ$	$-ZZZZ$

$$|\psi'\rangle = X_1 |\psi\rangle$$

no error

$$XXXX |\psi\rangle = |\psi\rangle$$

$$ZZZZ |\psi\rangle = |\psi\rangle$$

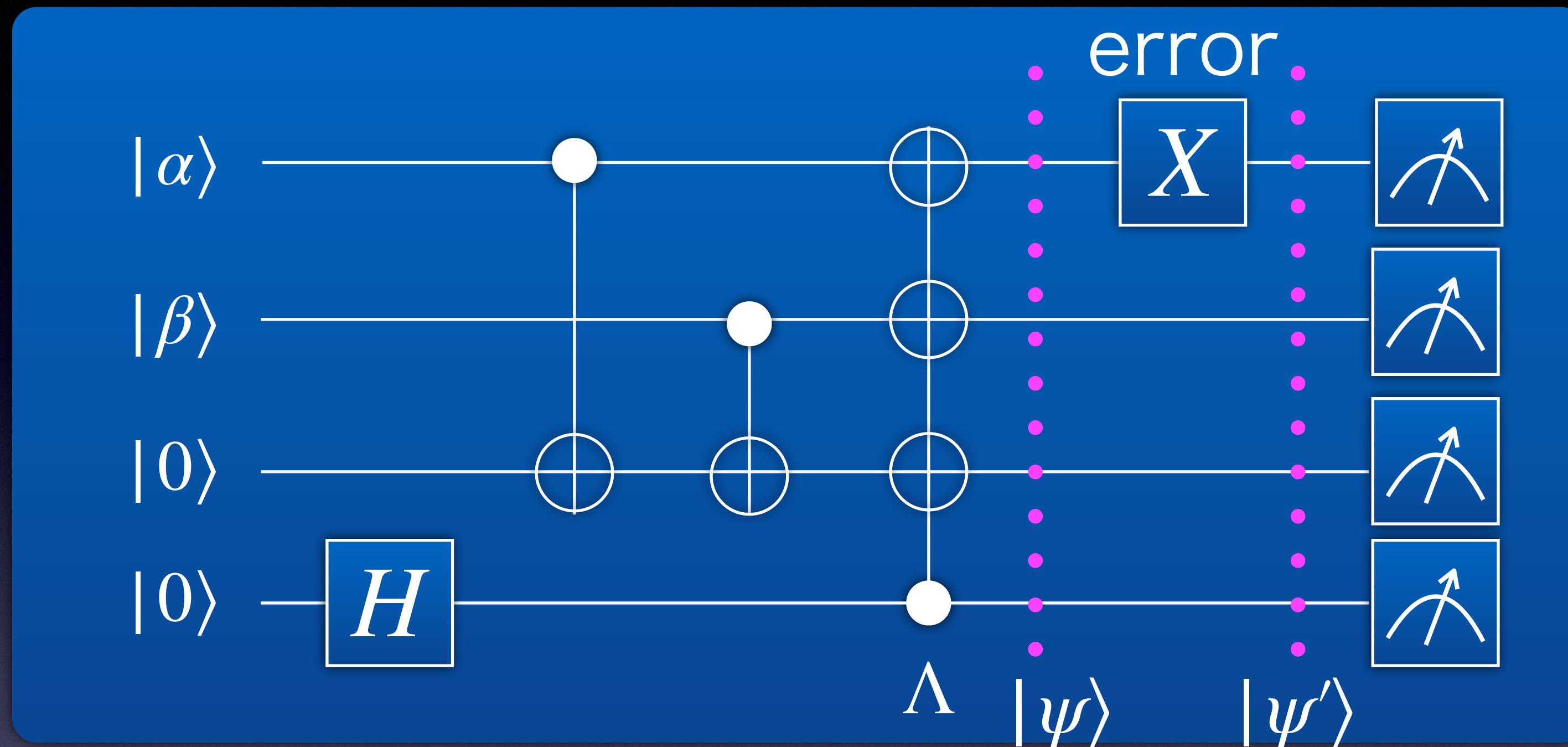
$X$  error

$$XXXX |\psi'\rangle = |\psi'\rangle$$

$$ZZZZ |\psi'\rangle = -|\psi'\rangle \quad \{ZZZZ, XIII\} = 0$$



# Quantum Error-Correcting codes



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outcome of  $ZZZZ$

+1

$X$  error

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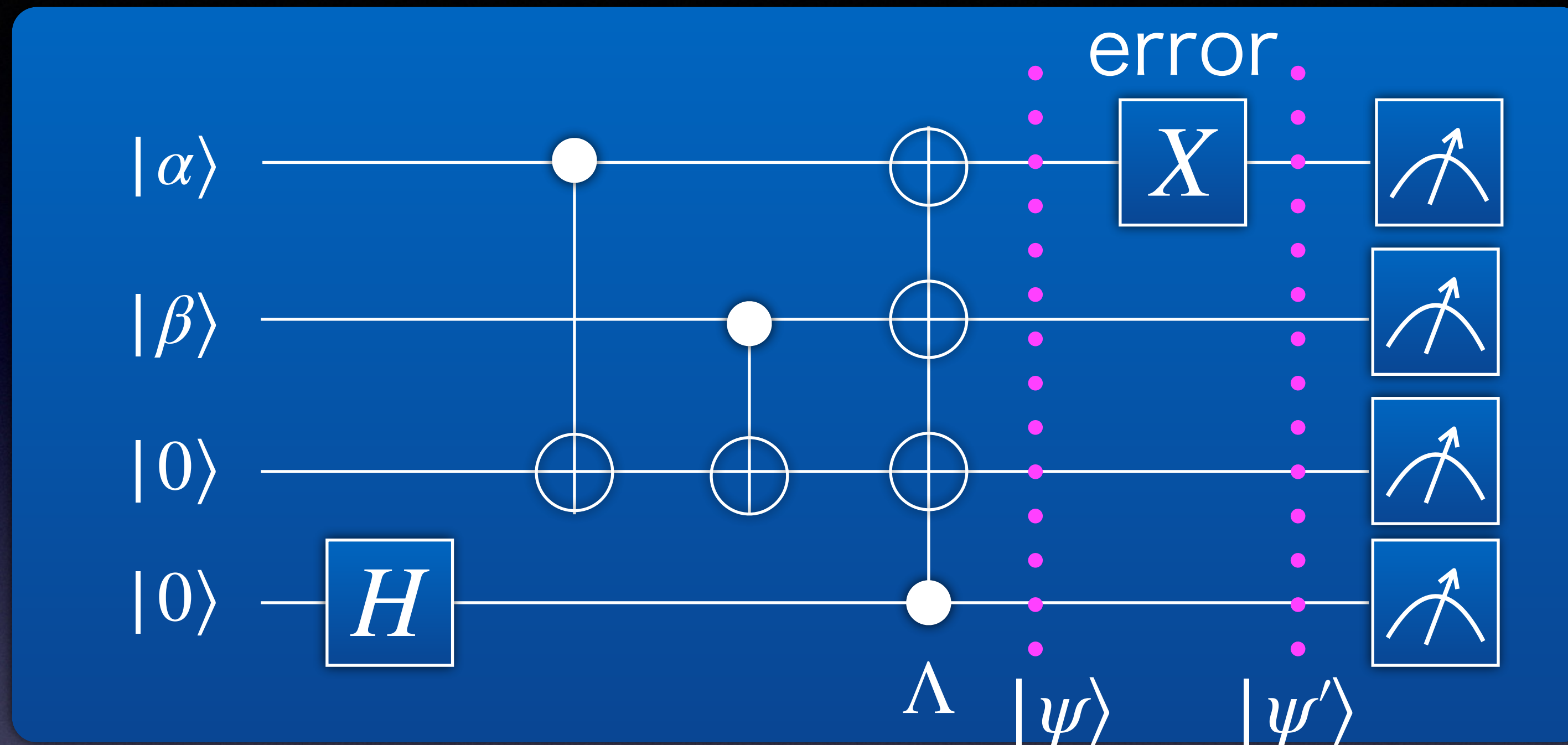
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# Quantum Error-Correcting codes



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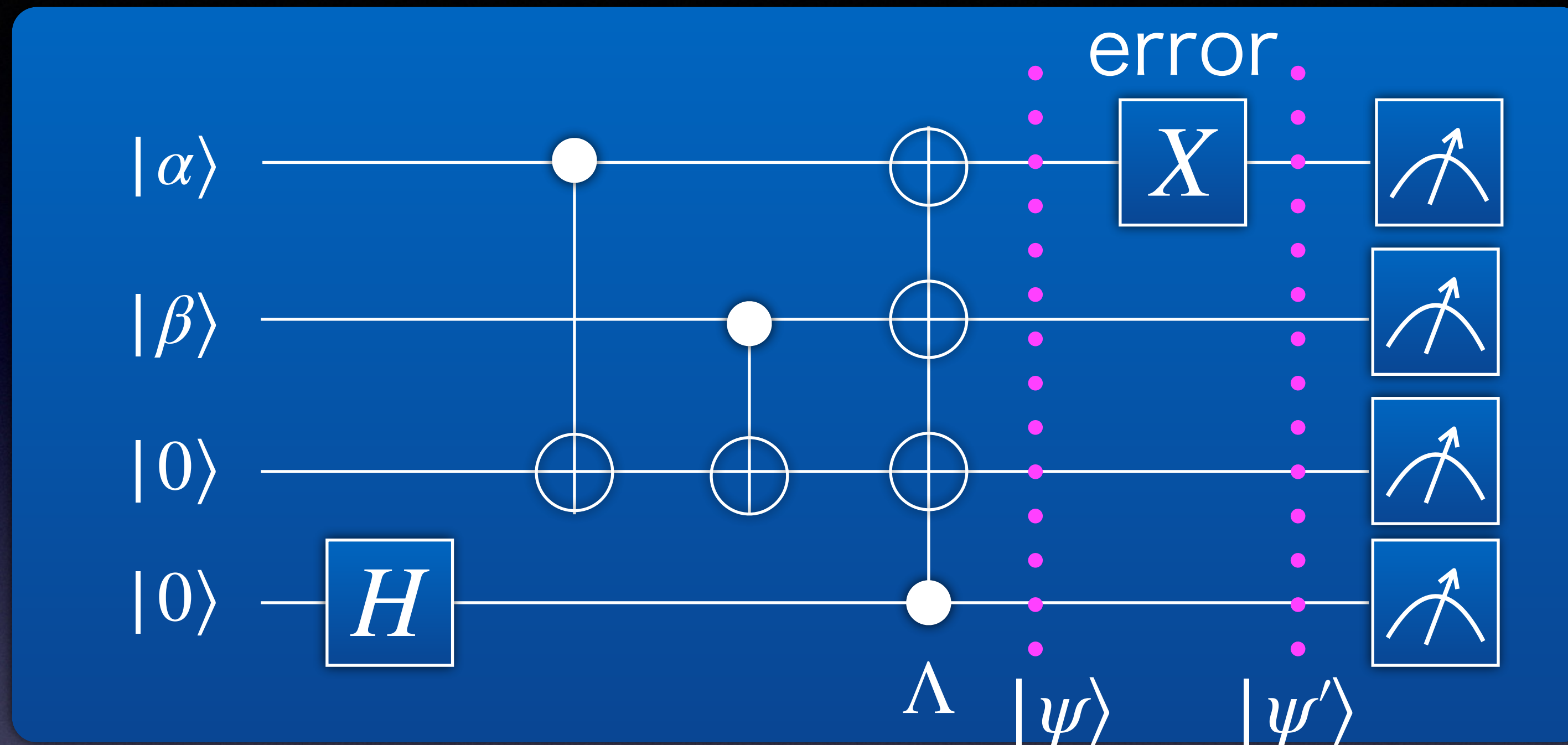
$$ZZZZ |\psi'\rangle = -|\psi'\rangle \quad \{ZZZZ, XIII\} = 0$$

-1

an error anticommuting with a stabilizer can be detected !



# Quantum Error-Correcting codes



$\Lambda$	$X$
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$ZZZZ$	$-ZZZZ$

$$|\psi'\rangle = X_1 |\psi\rangle$$

no error

$$XXXX |\psi\rangle = |\psi\rangle$$

$$ZZZZ |\psi\rangle = |\psi\rangle$$

+1

X error

$$XXXX |\psi'\rangle = |\psi'\rangle$$

$$ZZZZ |\psi'\rangle = -|\psi'\rangle \quad \{ZZZZ, XIII\} = 0$$

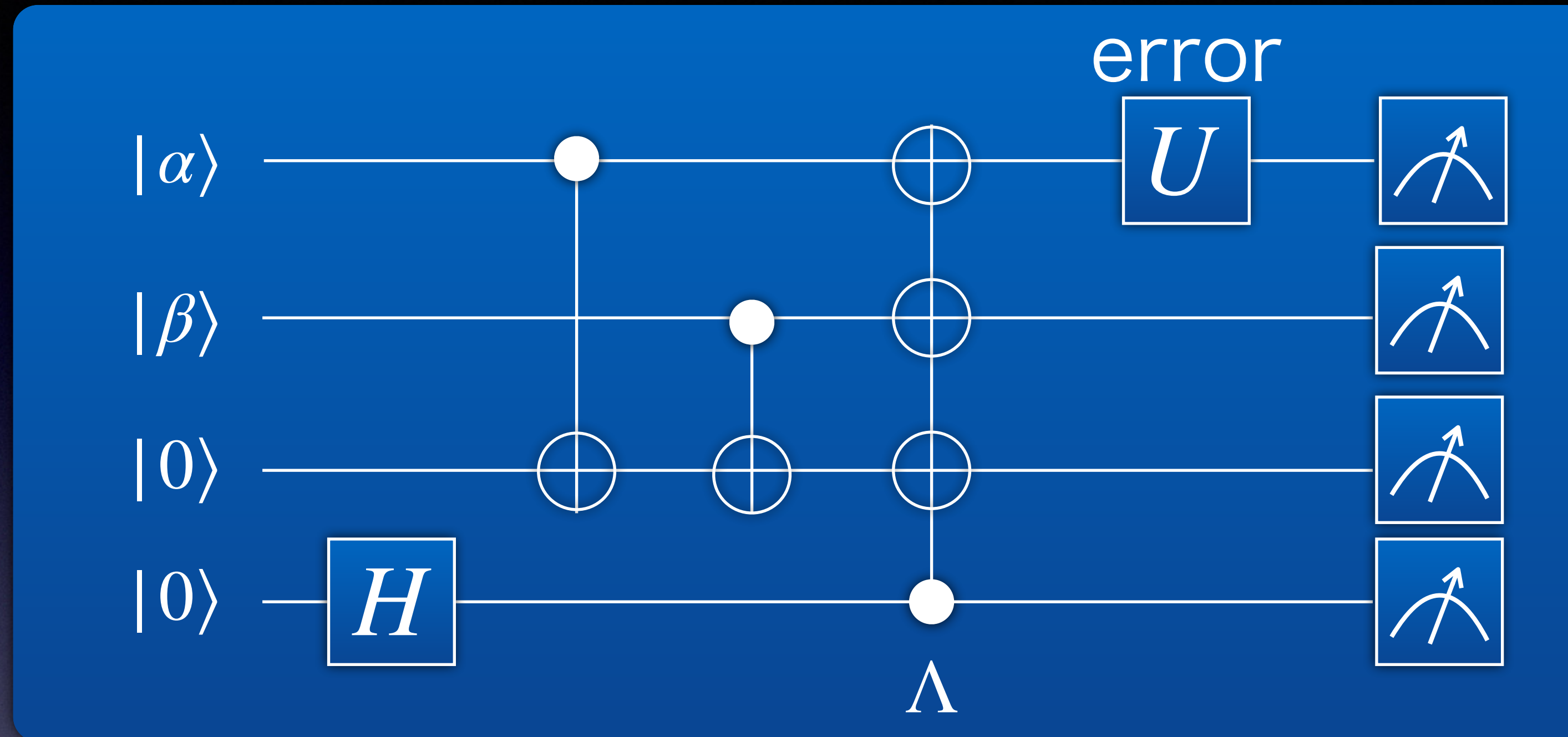
-1

outcome of  $ZZZZ$

the state does not collapse, because it's an eigenstate of the stabilizer!



# Quantum Error-Correcting codes

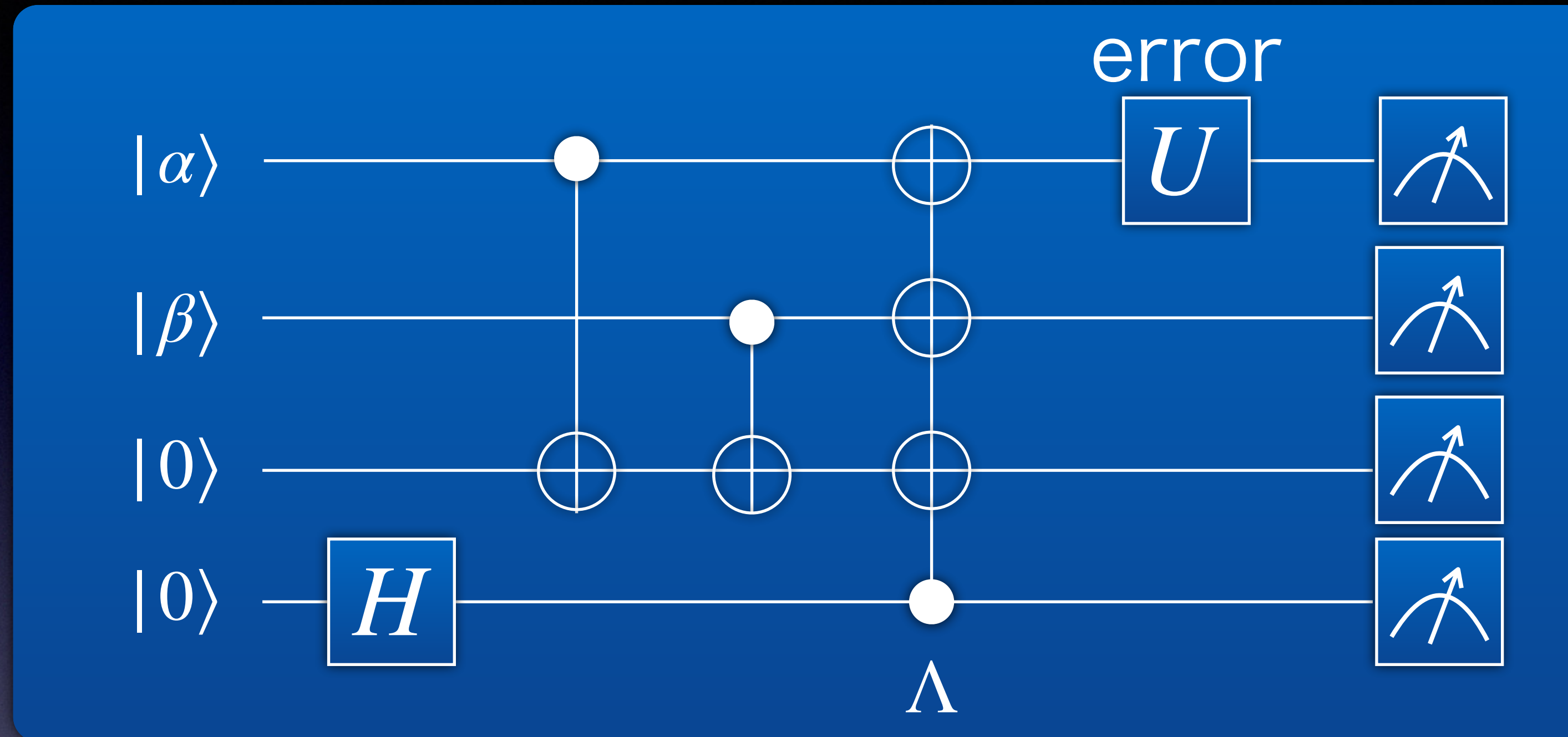


$$\begin{array}{cccc}
 U = X & Z & Y & \\
 XXXX & XXXX & -XXXX & -XXXX \\
 ZZZZ & -ZZZZ & ZZZZ & -ZZZZ
 \end{array}$$

- any 1-qubit Pauli error can be detected by stabilizers  $XXXX$  and  $ZZZZ$
- 2-qubit errors that commute with them cannot be detected 🙄  $XZ, YX$ , 🙅  $XX, YY, ZZ$



# Quantum Error-Correcting codes

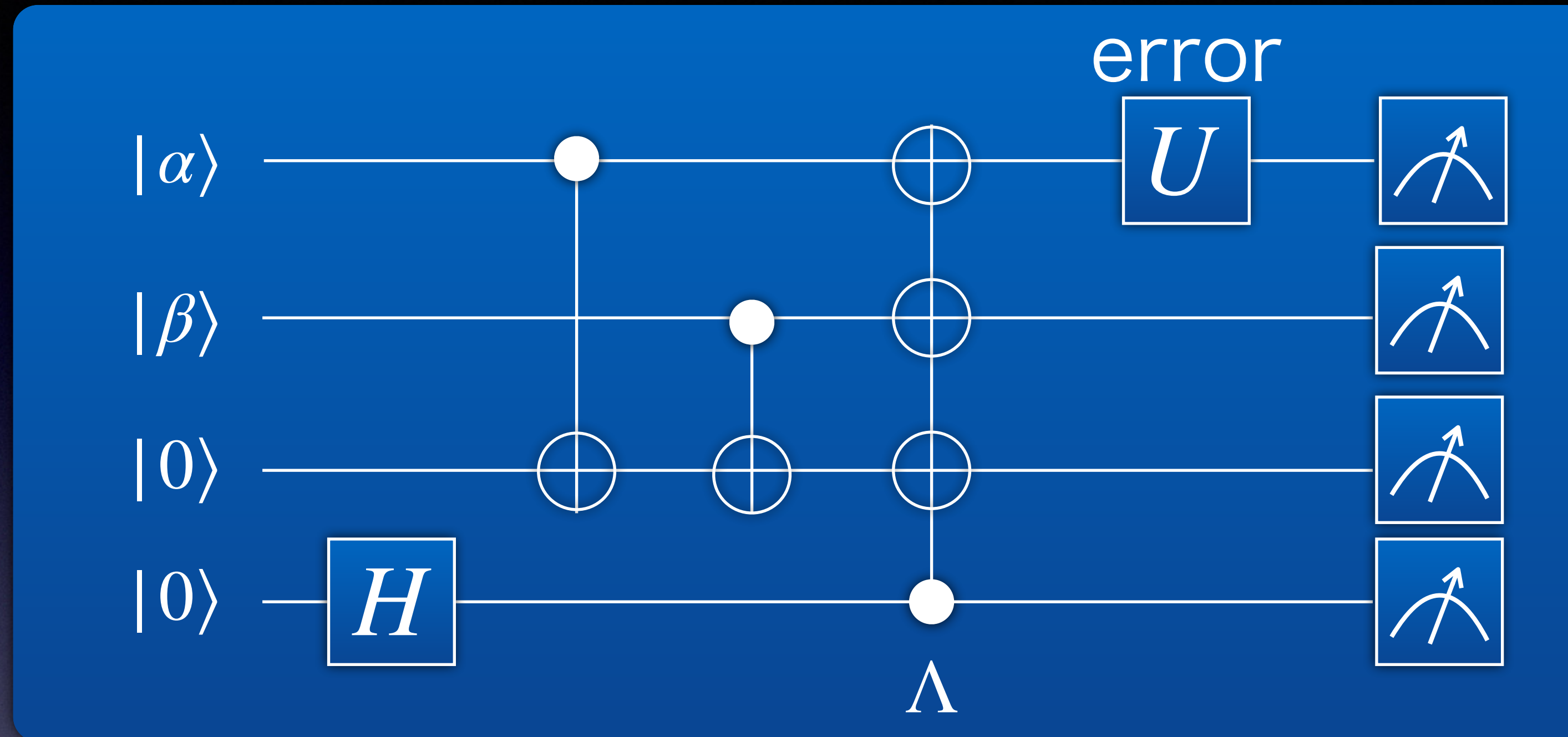


general error:  $U = aI + bX + cY + dZ$       $|\psi\rangle \rightarrow a|\psi\rangle + bX_1|\psi\rangle + cY_1|\psi\rangle + dZ_1|\psi\rangle$

	$ \psi\rangle$	$X_1 \psi\rangle$	$Z_1 \psi\rangle$	$Y_1 \psi\rangle$
$XXXX$	+1	+1	-1	-1
$ZZZZ$	+1	-1	+1	-1



# Quantum Error-Correcting codes



$$|\psi\rangle \rightarrow a|\psi\rangle + bX_1|\psi\rangle + cY_1|\psi\rangle + dZ_1|\psi\rangle$$

outcome of  $XXXX$       +1      +1      -1      -1

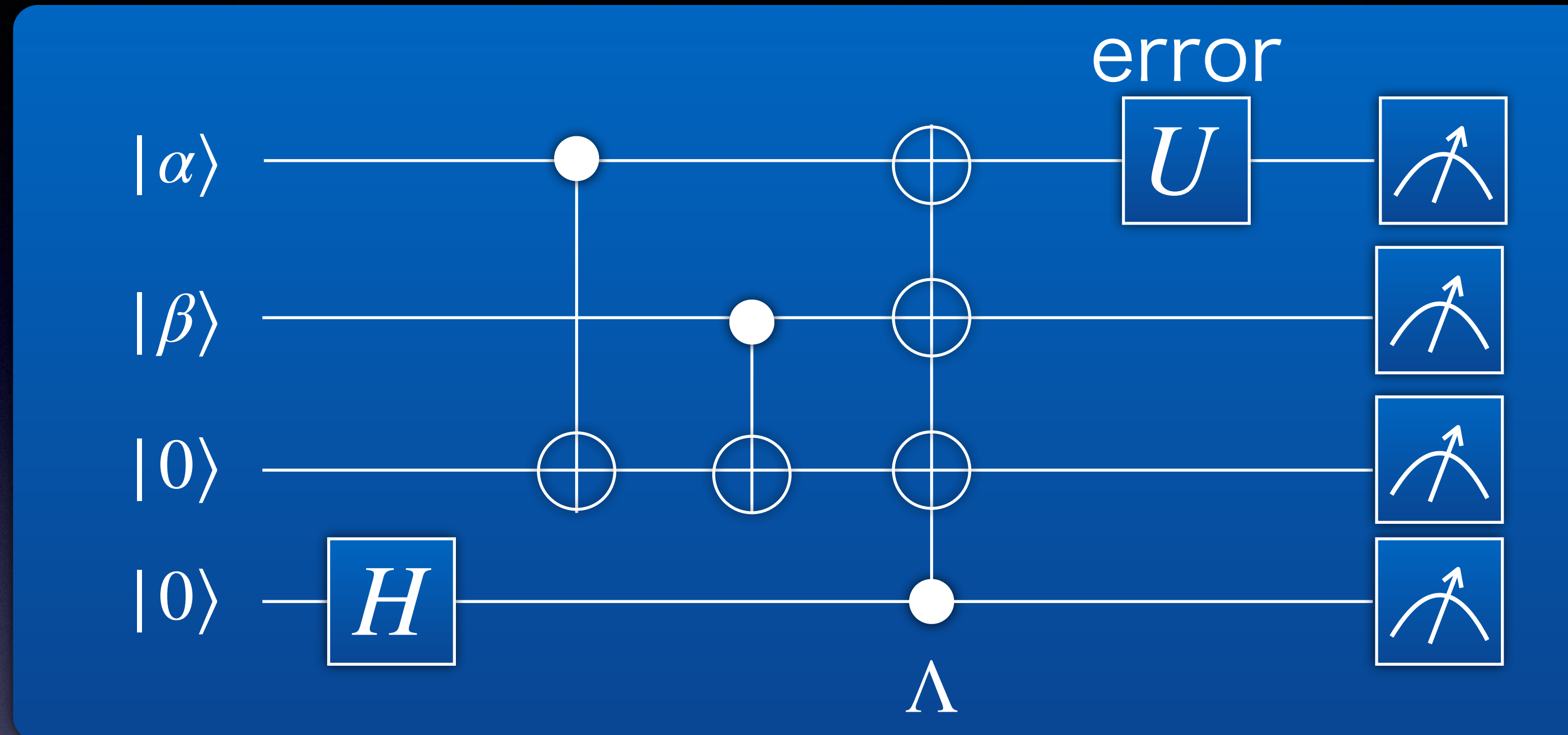
outcome of  $ZZZZ$       +1      -1      +1      -1

state after measurement       $|\psi\rangle$        $X_1|\psi\rangle$        $Z_1|\psi\rangle$        $Y_1|\psi\rangle$

measurement collapses the general error into a Pauli error



# Quantum Error-Correcting codes



$$|\psi\rangle \rightarrow a|\psi\rangle + bX_1|\psi\rangle + cY_1|\psi\rangle + dZ_1|\psi\rangle$$

outcome of  $XXXX$       +1      +1      -1      -1

outcome of  $ZZZZ$       +1      -1      +1      -1

state after measurement       $|\psi\rangle$        $X_1|\psi\rangle$        $Z_1|\psi\rangle$        $Y_1|\psi\rangle$

general error can be corrected by operating the detected Pauli error!