



New Lipschitz estimates and long-time asymptotic behavior for porous medium and fast diffusion equations

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Joint work with Filippo Santambrogio

Gradient flows face to face

Lyon

13th September

$$\frac{\partial n}{\partial t} = \nabla \cdot (n \nabla (p + V)), \quad \text{in } (0, \infty) \times \Omega, \quad d \geq 2,$$

$$n(0, x) = n_0(x) \geq 0, n_0 \in L^1(\Omega)$$

- Ω is either \mathbb{R}^d or a convex bounded set (hom. Neumann boundary conditions)
- Constitutive law of the pressure:

$$p = P(n) := \text{sign}(\gamma)n^\gamma,$$

- $V : \mathbb{R}^d \rightarrow \mathbb{R}$ is either: $V = 0$, $V = |x|^2/2$, $|\nabla V|, D^2V \in L^\infty(\mathbb{R}^d)$

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Goal: study the quantity

$$u(t) := \max_x |p(t, x)|^b |\nabla p(t, x) + \nabla V(x)|^2$$

Strategy: exploit the pressure equation

$$\frac{\partial p}{\partial t} = \gamma p \Delta q + \nabla p \cdot \nabla q, \quad \text{with } q := p + V.$$

$$\frac{\partial n}{\partial t} = \frac{|\gamma|}{\gamma + 1} \Delta n^{\gamma+1}$$

- for $\gamma > 0$: slow diffusion or **porous medium equation**
- for $-2/d < \gamma < 0$: **fast diffusion equation**

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Aronson-Bénilan estimate:

$$\Delta p \geq -\frac{1}{(\gamma + \frac{2}{d})t}, \quad \text{for all } t > 0.$$

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Barenblatt solution:

$$\mathcal{B}(t, x) = t^{-\alpha d} F(xt^{-\alpha}), \quad F(\xi) = (C - k|\xi|^2)_+^{\frac{1}{\gamma}},$$

$$\alpha = \frac{1}{d\gamma + 2}, \quad k = \frac{\alpha\gamma}{2(\gamma + 1)}$$

Some remarks: the porous medium equation

Regularity:

- Hölder continuity of $n(t, x)$ in space and time
- the pressure $p(t, x)$ is actually Lipschitz after a *waiting time* (for $d \geq 2$)

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Bénilan's lecture notes: for $1 - \gamma^2(d - 1) > 0$

$$\begin{aligned} |\nabla n^{\gamma+1}|^2 &\leq \frac{nK_1}{t}, & K_1 &= K_1(\gamma, d, \|n_0\|_\infty), \\ u(t) = \max_x n |\nabla n^\gamma|^2 &\leq \frac{K_2}{t}, & K_2 &= K_2(\gamma, d, \|n_0\|_\infty) \end{aligned}$$

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Convergence to the self-similar profile:

$$\lim_{t \rightarrow \infty} \|n(t) - \mathcal{B}(t)\|_{L^1(\mathbb{R}^d)} = 0,$$
$$\lim_{t \rightarrow \infty} t^{\alpha d} \|n(t) - \mathcal{B}(t)\|_{L^\infty(\mathbb{R}^d)} = 0,$$

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Some literature: Aronson, Bénilan, Caffarelli, Friedman, Gil, Herrero, Vázquez, Wolanski (1970-1987), Carrillo, Toscani (2000), Carrillo, Jüngel, Markowich, Toscani, Unterreiter (2001), Vázquez's book (2006)

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- wait Nikita Simonov's talk

Main results

Goal: study the quantity $u(t) := \max_x |p(t)|^b |\nabla q(t)|^2$, where $q = p + V$

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- **Trivial V:** there exists a positive constant C such that

$$\max_x |\rho(t)|^b |\nabla \rho(t)|^2 \leq Ct^{-1-\gamma d(b+1)\alpha}, \text{ for all } t > 0,$$

where $\alpha = \frac{1}{d\gamma + 2}$. Moreover, the exponent is sharp.

Assuming $n_0 \in L^1(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$, there exists a positive constant C such that

$$\max_x |\rho(t)|^b |\nabla \rho(t)|^2 \leq Ct^{-1}, \text{ for all } t > 0.$$

- **Quadratic V:** let $C_0 := \sup_x |\rho_0|^b |\nabla \rho_0 + x|^2 < \infty$, there exists a positive constant C such that

$$\max_x |\rho(t)|^b |\nabla q(t)|^2 \leq C_0 e^{-Ct}, \text{ for all } t > 0.$$

Assumptions on γ and b

For $V=0$, for $\gamma, b > 0$, resp. $\gamma < 0, b < -1$, we assume

$$(i) \gamma < \frac{1}{\sqrt{d-1}}, \text{ resp. } |\gamma| < \frac{2}{d},$$

$$(ii) 1 - \sqrt{1 - \gamma^2(d-1)} < \gamma b < 1 + \sqrt{1 - \gamma^2(d-1)}$$

For $V = |x|^2/2$, for $\gamma, b > 0$, resp. $\gamma < 0, b < -1$, we assume

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$$\begin{aligned} \gamma > 0 \quad \delta \text{ small} \\ \gamma < 0 \quad -\frac{2}{d} < \gamma < 0 \end{aligned}$$

$$b \in [c^1(\gamma, d), c^2(\gamma, d)]$$

$$\frac{1}{\delta}$$

Ideas of proof

Let $V = 0$, the pressure equation is

$$\frac{\partial p}{\partial t} = \gamma p \Delta p + |\nabla p|^2.$$

We will study $\max_x e^{\alpha(p(t,x))} |\nabla p(t,x)|^2$ and then choose $\alpha(p) = \ln|p|^b$.
Since we are at a maximum point

$$\nabla \left(e^{\alpha(p)} \frac{|\nabla p|^2}{2} \right) = 0, \quad \Delta \left(e^{\alpha(p)} \frac{|\nabla p|^2}{2} \right) \leq 0.$$

The first condition gives

$$D^2 p \nabla p = -\frac{\alpha'}{2} |\nabla p|^2 \nabla p,$$

while from the second we get

$$\Delta \left(\frac{|\nabla p|^2}{2} \right) \leq \frac{|\nabla p|^4}{2} (|\alpha'|^2 - \alpha'') - \alpha' \frac{|\nabla p|^2}{2} \Delta p.$$

$$\begin{aligned}
 & \frac{\partial}{\partial t} \left(e^\alpha \frac{|\nabla p|^2}{2} \right) \\
 &= e^\alpha \left(\alpha' \frac{|\nabla p|^2}{2} (\gamma p \Delta p + |\nabla p|^2) + \nabla p \cdot \nabla (\gamma p \Delta p + |\nabla p|^2) \right) \\
 &= e^\alpha \left(\gamma p \frac{\alpha'}{2} |\nabla p|^2 \Delta p + \frac{\alpha'}{2} |\nabla p|^4 + \gamma |\nabla p|^2 \Delta p + \gamma p \nabla p \cdot \nabla \Delta p + 2 \nabla p D^2 p \nabla p \right).
 \end{aligned}$$

Now we use $2 \nabla p \cdot \nabla \Delta p = \Delta(|\nabla p|^2) - 2|D^2 p|^2$, where $|D^2 p|^2 := \sum_{i,j=1}^d (\partial_{i,j} p)^2$

$$\gamma p \nabla p \cdot \nabla \Delta p \leq \gamma p \left(\frac{|\nabla p|^4}{2} (|\alpha'|^2 - \alpha'') - \alpha' \frac{|\nabla p|^2}{2} \Delta p \right) - \gamma p |D^2 p|^2$$

Therefore we have

$$\frac{\partial}{\partial t} \left(e^\alpha \frac{|\nabla p|^2}{2} \right) \leq e^\alpha \left(\frac{|\nabla p|^4}{2} (\gamma p |\alpha'|^2 - \gamma p \alpha'' - \alpha') + \gamma |\nabla p|^2 \Delta p - \gamma p |D^2 p|^2 \right)$$

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Ideas of proof

Let us denote $\lambda := -\frac{\alpha'}{2}|\nabla p|^2$, $\delta_i = (D^2q)_{i,i}$ for $i = 2, \dots, d$, $\delta := \sum_{i=2}^d \delta_i$.

Since $\gamma p > 0$, we have

$$\begin{aligned}\gamma|\nabla p|^2\Delta p - \gamma p|D^2p|^2 &\leq \gamma|\nabla p|^2(\lambda + \delta) - \gamma p\left(\lambda^2 + \sum_{i=2}^d \delta_i^2\right) \\ &\leq \gamma|\nabla p|^2(\lambda + \delta) - \gamma p\left(\lambda^2 + \frac{\delta^2}{d-1}\right) \\ &= -\gamma\frac{\alpha'}{2}|\nabla p|^4 + \gamma\delta|\nabla p|^2 - \gamma p\frac{|\alpha'|^2}{4}|\nabla p|^4 - \gamma p\frac{\delta^2}{d-1}.\end{aligned}$$

Using Young's inequality, we have

$$\gamma\delta|\nabla p|^2 \leq \gamma p\frac{\delta^2}{d-1} + \gamma\frac{|\nabla p|^4}{4p}(d-1),$$

and we finally find

$$\gamma|\nabla p|^2\Delta p - \gamma p|D^2p|^2 \leq \left(-\gamma\frac{\alpha'}{2} + \frac{\gamma(d-1)}{4p} - \gamma p\frac{|\alpha'|^2}{4}\right)|\nabla p|^4.$$

Let's come back to the main estimate

$$\frac{\partial}{\partial t} \left(e^\alpha \frac{|\nabla p|^2}{2} \right) \leq e^\alpha \left(|\nabla p|^4 \left(-\frac{\alpha'}{2} + \gamma p \frac{|\alpha'|^2}{4} - \gamma p \frac{\alpha''}{2} - \gamma \frac{\alpha'}{2} + \frac{\gamma(d-1)}{4p} \right) \right).$$

Choosing $\alpha(p) = b \ln |p|$ we find

$$\frac{\partial}{\partial t} \left(|p|^b \frac{|\nabla p|^2}{2} \right) \leq c_0 |p|^{b-1} |\nabla p|^4,$$

where c_0 is defined as

$$c_0 = -\frac{|b|}{2} + \frac{|\gamma|b^2}{4} + \frac{|\gamma|(d-1)}{4},$$

and is **negative**.

Let's come back to the main estimate

$$\frac{\partial}{\partial t} \left(|\rho|^b \frac{|\nabla \rho|^2}{2} \right) \leq |\rho|^b \left(|\nabla \rho|^4 \left(-\frac{\alpha'}{2} + \gamma \rho \frac{|\alpha'|^2}{4} - \gamma \rho \frac{\alpha''}{2} - \gamma \frac{\alpha'}{2} + \frac{\gamma(d-1)}{4\rho} \right) \right).$$

Choosing $\alpha(\rho) = b \ln |\rho|$ we find

$$\frac{\partial}{\partial t} \left(|\rho|^b \frac{|\nabla \rho|^2}{2} \right) \leq c_0 |\rho|^{b-1} |\nabla \rho|^4,$$

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$$\frac{\partial}{\partial t} \left(|\rho|^b \frac{|\nabla \rho|^2}{2} \right) \leq c_0 |\rho|^{-b-1} |\rho|^{2b} |\nabla \rho|^4,$$

Let us first prove the claimed result for $n_0 \in L^\infty(\mathbb{R}^d)$.

For $\gamma, b > 0$, we have $|\rho(t, x)|^{-b-1} \geq |\bar{\rho}|^{-b-1}$.

For $\gamma < 0, b \leq -1$, we have $|\rho(t, x)|^{-b-1} \geq |\bar{\rho}|^{-b-1}$.

Therefore, in both cases we find

$$\frac{\partial}{\partial t} \left(|\rho|^b \frac{|\nabla \rho|^2}{2} \right) \leq c_0 |\bar{\rho}|^{-b-1} |\rho|^{2b} |\nabla \rho|^4,$$

From this we get

$$u'(t) \leq -Cu^2(t),$$

$$u(t) \leq Ct^{-1}.$$

Ideas of the proof

Without asking $n_0 \in L^\infty(\mathbb{R}^d)$ we exploit the Aronson-Bénilan estimate to find

$$\|n(t)\|_\infty \leq Ct^{-\frac{d}{d\gamma+2}},$$

which is equivalent to

- $\max_x |p(t)| \leq Ct^{-\frac{d\gamma}{d\gamma+2}}$, for $\gamma > 0$,
- $\min_x |p(t)| \geq Ct^{-\frac{d\gamma}{d\gamma+2}}$, for $\gamma < 0$.

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Therefore, we have

$$\begin{aligned} \frac{\partial}{\partial t} \left(|p|^b \frac{|\nabla p|^2}{2} \right) &\leq c_0 |p|^{2b} |p|^{-(b+1)} |\nabla p|^4 \\ &\leq -Cp^{2b} |\nabla p|^4 t^{(b+1)\frac{d\gamma}{d\gamma+2}}, \end{aligned}$$

from which we find

$$u(t) \leq Ct^{-1-\gamma d(b+1)\alpha}.$$

Ideas of the proof

Let us consider the quadratic potential $V = |x|^2/2$, then

$$\frac{\partial p}{\partial t} = \gamma p \Delta q + \nabla p \cdot \nabla q, \text{ with } q = p + V.$$

With the same argument, we can find

$$\frac{\partial}{\partial t} \left(|p|^b \frac{|\nabla q|^2}{2} \right) \leq |p|^{b-1} (c_1 |\nabla q|^2 |\nabla p|^2 + c_2 |\nabla q \cdot \nabla p|^2) + c_3 |p|^b |\nabla q|^2,$$

where

$$c_3 = \frac{\gamma b d}{2} - 1.$$

The assumptions ensure $c_0 := c_1 + c_2 \leq 0$, and $c_1, c_3 < 0$. Therefore, we establish

$$\begin{aligned} \frac{\partial}{\partial t} \left(|p|^b \frac{|\nabla q|^2}{2} \right) &\leq c_0 |p|^b |\nabla q \cdot \nabla p|^2 + c_3 |p|^b |\nabla q|^2 \\ &\leq c_3 |p|^b |\nabla q|^2, \end{aligned}$$

and finally

$$u(t) \leq C_0 e^{-Ct}, \quad \forall t > 0.$$

Implications: rate of convergence to the Barenblatt

Time-dependent scaling:

$$\hat{n}(t, x) := \varphi(t)^d n(\psi(t), \varphi(t)x), \quad \varphi(t) = e^t, \quad \psi(t) = e^{(d\gamma+2)t}.$$

If $n(t, x)$ is a solution of

$$\frac{\partial n}{\partial t} = \frac{|\gamma|}{\gamma+1} \Delta n^{\gamma+1},$$

then $\hat{n}(t, x)$ satisfies

$$\frac{\partial \hat{n}}{\partial t} = \frac{|\gamma|}{\gamma+1} \Delta \hat{n}^{\gamma+1} + \nabla \cdot (\hat{n}x).$$

By our result:

$$\max_x |\hat{p}(t, x)|^b |\nabla \hat{p}(t, x) + x|^2 \leq C_0 e^{-Ct},$$

which translate into

$$\max_x |p(t, x)|^b |\nabla p(t, x) + xt^{-1}|^2 \leq C_0 t^\beta t^{-\frac{C}{d\gamma+2}},$$

with $\beta = -\alpha\gamma db - 2 + 2\alpha$.

Implications: rate of convergence to the Barenblatt

$$\max_x |\rho(t, x)|^b |\nabla \rho(t, x) + xt^{-1}|^2 \leq C_0 t^\beta t^{-\frac{c}{d\gamma+2}},$$

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For $\gamma < 0$, the pressure of the Barenblatt solution is

$$\mathcal{P}(t, x) = t^{-\alpha\gamma d} (C - k|x|^2 t^{-2\alpha}).$$

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Let us take $x \in C_{a,b}(t) := \{x; |x|t^{-\alpha} \in [a,b]\}$ for some constants $0 < a < b$.
For $n_0 \in \mathcal{X} \setminus \{0\}$ (see Nikita's talk)

$$n^{\gamma b}(t,x) \gtrsim t^{-\alpha\gamma db}, \quad |\nabla \mathcal{P}(t,x)|^2 \simeq t^{-2+2\alpha}.$$

Finally, we find

$$\frac{\|\nabla p(t) - \nabla \mathcal{P}(t)\|_{L^\infty(C_{a,b}(t))}^2}{\|\nabla \mathcal{P}(t)\|_{L^\infty(C_{a,b}(t))}^2} \lesssim t^{-\frac{c}{d\gamma+2}}, \quad \forall t > 0.$$

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Thank you for the attention!