

Gluing variations

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Beijing, July 2023

based on joint work with Wan Cong
arXiv:2302.06928 [gr-qc]

Questions

Throughout this work: vacuum spacetimes with $\Lambda \in \mathbb{R}$

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 0.$$

Can you

- 1 extend vacuum spacetimes?
- 2 glue together vacuum spacetimes?
- 3 realise *data on lower dimensional submanifolds* by embedding in a vacuum spacetime?
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Data on lower dimensional manifolds?

The Aretakis-Czimek-Rodnianski question

QUESTION (Aretakis, Czimek and Rodnianski (2021))

Can you find vacuum characteristic initial data interpolating between two sphere data sets?

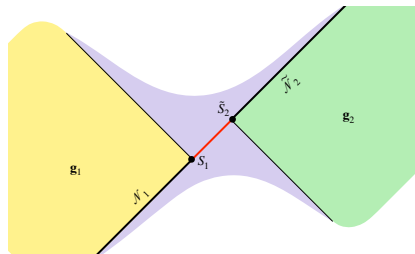


Figure: Gluing construction of Aretakis-Czimek-Rodnianski

Answer: “kind of”, with obstructions, for sphere data near spheres lying on a Minkowskian light cone



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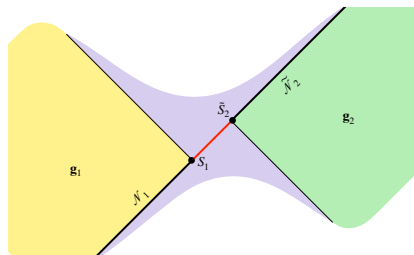


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ACR gluings: applications

Optimal asymptotics, control, third law

Theorem (Czimek & Rodnianski, arXiv:2210.09663)

*The Corvino-Schoen gluing can be done with **controlled mass, momentum, angular momentum and center of mass.***

Theorem (Kehle, Unger, arXiv:2211.15742)

*The “third law of black hole dynamics” is **wrong.***

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Gluing method

Nonlinear “superpositions”

- In **linear** theories, new initial data can be obtained by *adding* old ones



Alternative approach:

gluing



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6 functions g_{ij}

6 functions K_{ij}

= 12

minus (3-dimensional diffeomorphism + 1 choice of initial slice)

= 4

8 functions, 4 constraint equations



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Theorem (Corvino & Schoen, 2010)

*An asymptotically flat initial data set can be glued to a **Kerr one**, with nearby global charges across a sufficiently distant coordinate annulus.*



State-of-the-art variations on Corvino-Schoen

Gluing-in small black holes with $\Lambda = 0$

Theorem (Peter Hintz, arXiv:2210.13960)

*Let (Σ, g, K) be a vacuum initial data set and suppose that there are no Killing vectors near $p \in \Sigma$. For every $\epsilon > 0$ sufficiently small there exists a vacuum initial data set which coincides with (g, K) outside an ϵ -neighborhood of p and **coincides with a small Kerr black hole** inside the neighborhood.*

This can be done all over the place

State-of-the-art variations on Corvino-Schoen

Gluing-in small black holes with $\Lambda = 0$; the Hintz black hole sprinkler (compare Anderson, Corvino, Pasqualotto arXiv:2301.08238)

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Asymptotic gluing:

Gluing-in black holes with $\Lambda > 0$ (P. Hintz, arXiv:2001.10401)

Theorem 1.1. *Let $N \in \mathbb{N}$. For $i = 1, \dots, N$, fix points $p_i \in \partial M = \mathbb{S}^3 \subset \mathbb{R}^4$ and (subextremal) masses $0 < m_i < (3\Lambda)^{-1/2}$ such that the balance condition*

$$\sum_{i=1}^N m_i p_i = 0 \in \mathbb{R}^4. \quad (1.2)$$

holds. Then there exists a metric g solving the Einstein vacuum equation (1.1) in a neighborhood of ∂M with the following properties:

- (1) *in a neighborhood of p_i , g is isometric to a Schwarzschild–de Sitter black hole metric with mass m_i , containing future affine complete event and cosmological horizons;*
- (2) *outside a small neighborhood of $\{p_1, \dots, p_N\}$, $\cos^2(s)g$ is smooth down to $s = \pi/2$, and asymptotic to the rescaled de Sitter metric $\cos^2(s)g_{\text{dS}}$ at the rate $\cos^3(s)$.*

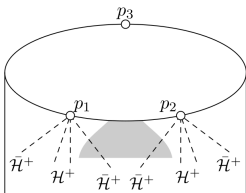


FIGURE 1.2. Illustration of Theorem 1.1. We glue SdS black holes into neighborhoods of the points p_i ; only two black holes are shown here. The

- The linearised prescribed scalar constraint equation at the Euclidean metric is

$$\delta R[h] \equiv \partial_i \partial_j (h^{ij} - h^k{}_k \delta^{ij}) = f.$$

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- using the **Bogovskii operator**, one can construct a solutions of this equation which **vanish** outside of the support of f
- this can be used to give a simple and sharper version of the Corvino-Schoen gluing:

Theorem (Mao, Oh, Tao, arXiv:2308.13031)

*An asymptotically flat initial data set can be glued to **any** Kerr (possibly Schwarzschild) **with longer** energy-momentum vector*



Mao-Tao's simplification of the Corvino-Schoen gluing

the original Corvino & Schoen theorem: to **some** nearby Kerr

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- using the **Bogovskii operator**, one can construct a solutions of this equation which **vanish** outside of the support of f
- and an alternative proof of the Carlotto-Schoen gluing:

Theorem (Mao & Tao, arXiv:2210.09437)

Simple proof of the Carlotto-Schoen theorem, including optimal decay, using spacelike gluing.

Screening the gravitational field

Carlotto-Schoen “exotic gluings” (2014)

- Remove a solid cone C_1 from Euclidean space; initial data $(\mathbb{R}^n, g = \delta, K_{ij} = 0)$



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- Remove a **solid cone** C_1 from Euclidean space; initial data $(\mathbb{R}^n, g = \delta, K_{ij} = 0)$
- Remove a **slightly larger cone** C_2 from an asymptotically flat initial data set (M, g_{ij}, K_{ij})



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Theorem (Carlotto and Schoen)

If the tip of C_2 is sufficiently far away there exists an initial data set which **coincides with** (M, g_{ij}, K_{ij}) **outside of C_2 and has Minkowskian data on C_1**



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Mao, Tao, arXiv:2210.09437: can be done with optimal $1/r$ decay using a Green function for δR supported in a cone.



ACR gluings: applications

“No third law”

CONJECTURE (“third law of black hole dynamics”, Bardeen, Carter & Hawking (1973))

A black hole with zero surface-gravity cannot be formed in a dynamical process.

zero surface-gravity \approx zero temperature

Theorem (Kehle & Unger, arXiv:2211.15742)

The third law is wrong for spherically symmetric solutions of the Einstein-Maxwell-charged-scalar-field equations.

Proof: use null gluing to an extreme Reissner-Nordström black hole.

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Theorem (Kehle & Unger (2023), arXiv:2304.08455)

Black holes can be formed in vacuum by focusing of gravitational waves.

Proof: null gluing of a Minkowskian light-cone to a Kerr black hole

Previous work: Christodoulou (2008), arXiv:0805.3880, 594 pages & Li and Yu (2015) 70 pages



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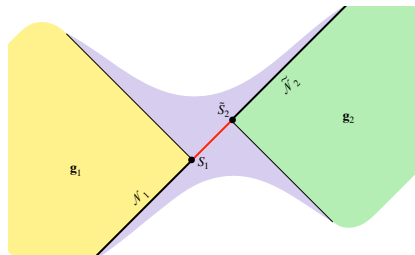


Figure: Gluing construction of ACR

Characteristic gluing: implicit function theorem together with

Theorem (Aretakis, Czimek & Rodnianski, arXiv:2107.02449)

The C^2 linearised characteristic gluing at $(3 + 1)$ -Minkowski is solvable up to a 10-dimensional space of obstructions.

$(3 + 1)$ -Minkowski: cross-section $\mathbf{S} \approx S^2$, $\Lambda = 0 = m$



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Theorem (Cong, PTC and Gray, arXiv:2401.04442)

The C^k linearised characteristic gluing at $(n + 1)$ -Birmingham - Kottler with $m \neq 0$ is solvable up to a space of obstructions of dimension $\leq n + 1$.

$(n + 1)$ -Birmingham - Kottler: cross-section \mathbf{S} compact Einstein spaces e.g. spheres, torus, higher genus; $\Lambda \in \mathbb{R}$, $m \in \mathbb{R}$.

General topologies, higher dimensions, differentiability

Wan Cong, PTC, and Finnian Gray, arXiv:2401.04442

Obstructions arise from kernels of linear elliptic operators on the cross-section \mathbf{S} of the characteristic hypersurface; affected by *dimension* and *topology* of \mathbf{S} , e.g.:

C^2 -gluing with $m = 0, \Lambda = 0$	S^2	T^2	S^4
dim. of obstruction space	10	7	30

Both a non-vanishing *mass* m and a non-zero *cosmological constant* Λ provide additional degrees of freedom to remove some of the obstructions, e.g.:

C^k -gluing	$S^2, m = 0$	$S^2, m = 0$	$S^2, m \neq 0$	$\mathbf{S}, m \neq 0$
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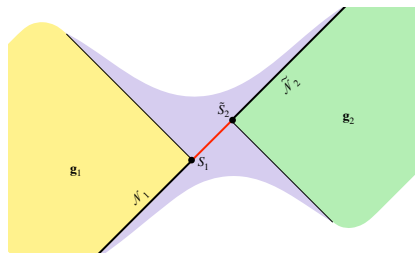


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Data on lower dimensional manifolds?

Vacuum jets

Let P be a submanifold of M .

Let $k \in \mathbb{N} \cup \{\infty\}$.

Definition

Let g be any smoothly differentiable metric defined in a neighborhood of P . The collection

$$j^k g := \{\partial_{\alpha_1} \cdots \partial_{\alpha_\ell} g_{\mu\nu} | P, 0 \leq \ell \leq k\}$$

will be called *jet of order k of g at P* .

Einstein equations

provide equations, differential and/or algebraic, for the jet of order k at P .

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A jet will be called **vacuum** if all such equations are satisfied

Data on lower dimensional manifolds?

Spacelike/timelike/null vacuum submanifold data

Definition

The collection of all vacuum jets of order k will be called **vacuum submanifold data of order k** and will be denoted by $\Psi[P, k]$.

A jet of order k of a metric g will be called $\left\{ \begin{array}{l} \text{spacelike,} \\ \text{timelike,} \\ \text{null,} \end{array} \right.$ if

the metric induced by g on P is.

We similarly define **spacelike/timelike/null/characteristic vacuum submanifold data of order k at P** .

QUESTION

Given a submanifold $P \subset M$ and a vacuum jet $j^k g$ in $\Psi[P, k]$, is there a vacuum metric on M which realises $j^k g$?

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Example: Vacuum jets at a point

Let P be a point, $P = \{p\}$

$j^k g$ = the coefficients of the Taylor series of a metric g at p .

vacuum \equiv algebraic conditions on the Taylor coefficients

Now: *in normal coordinates the Taylor coefficients can be expressed in terms of the Riemann tensor and its covariant derivatives.*

For example, using normal coordinates,

$$j^2 g|_p \approx \{g_{\mu\nu}|_p, W^{\alpha}{}_{\beta\gamma\delta}|_p\},$$

where $W^{\alpha}{}_{\beta\gamma\delta}$ has the symmetries of the Weyl tensor.

$\Psi[\{p\}, 2] \approx$ all such pairs

QUESTION

Is there a *vacuum* metric which realises $\Psi[\{p\},]$?



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$\Psi[\{p\}, 2] \approx$ all such pairs

QUESTION

Is there a *vacuum* metric which realises $\Psi[\{p\},]$?



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Data on lower dimensional manifolds

Example: Vacuum jets at a point

Let P be a point, $P = \{p\}$

$j^k g$ = the coefficients of the Taylor series of a metric g at p .

vacuum \equiv algebraic conditions on the Taylor coefficients

Now: *in normal coordinates the Taylor coefficients can be expressed in terms of the Riemann tensor and its covariant derivatives.*

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Is there a *vacuum* metric which realises $\Psi[\{p\}, \infty]$?

Codimension 1 vacuum spacelike hypersurface data

Vacuum spacelike constraint equations

Initial data surface Σ , **Riemannian** metric g_{ij} , $i, j = 1, \dots, n$, symmetric tensor K_{ij} (“initial time derivative of the metric”)

the **scalar constraint equation** (Λ is the *cosmological constant*):

$$R(g_{ij}) = 2\Lambda + |K|^2 - (\text{tr}K)^2,$$

and the *vector constraint equation*:

$$D_j K^j_k - D_k K^j_j = 0.$$

Fact

Spacelike vacuum hypersurface data

$$\Psi[\Sigma, \infty] \approx \Psi[\Sigma, k] \approx \Psi[\Sigma, 2] \approx \{ \text{all vacuum } (g, K) \}.$$

Proof: The Cauchy problem is well posed in, e.g., harmonic coordinates.



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Extension theorem

Can one **extend vacuum** spacelike initial data on a **manifold with boundary beyond its boundary**? (while, of course, satisfying the vacuum constraint equations)

Theorem

Let (Σ, g, K) be spacelike vacuum initial data on a manifold *with boundary* $\partial\Sigma$.

There exists a manifold *without boundary* $\check{\Sigma}$ and vacuum initial data (\check{g}, \check{K}) on $\check{\Sigma}$ such that $\Sigma \subset \check{\Sigma}$, with

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Proof: I will give a sketch; for this we will need *characteristic vacuum hypersurface data* $\Psi[\mathcal{N}, \infty]$.

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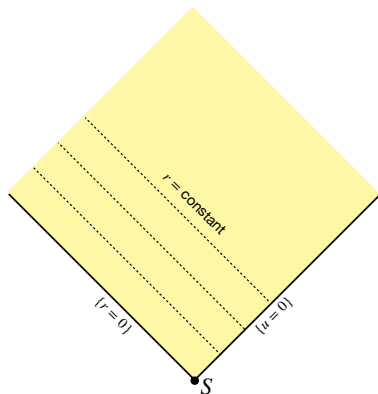
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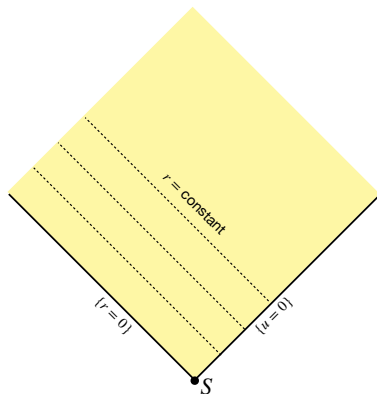
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Characteristic Cauchy problem



Transversally intersecting null hypersurfaces,

Characteristic Cauchy problem



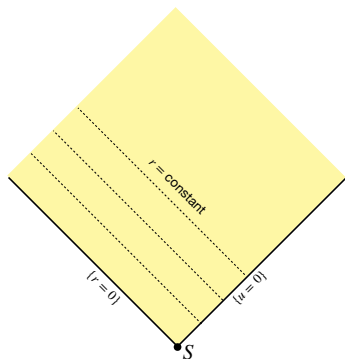
Transversally intersecting null hypersurfaces, or a light cone.

Characteristic Cauchy problem

Isenberg-Moncrief coordinates

The hypersurfaces $\mathcal{N} = \{u = 0\}$ and $\underline{\mathcal{N}} = \{r = 0\}$ are characteristic for the metric

$$g_{\mu\nu} dx^\mu dx^\nu = 2 \left(-du + u\alpha dr + u\beta_A dx^A \right) dr + g_{AB} dx^A dx^B.$$



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Theorem (4d: Rendall (1990);)

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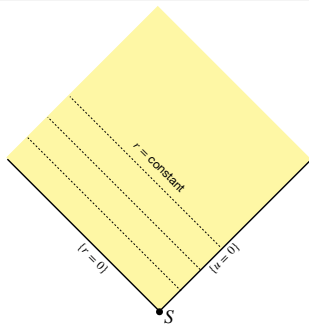
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Embedding a null hypersurface in a vacuum spacetime?



QUESTION

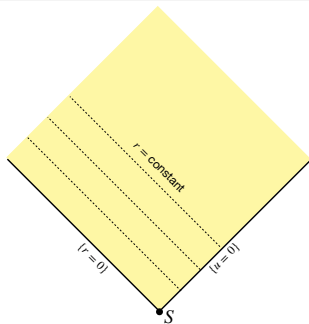
What about a vacuum metric in a **whole neighborhood** of $\mathcal{N} \cup \underline{\mathcal{N}}$? What about a **single** characteristic hypersurface?

Answer: one needs to understand data of order k on \mathcal{N}



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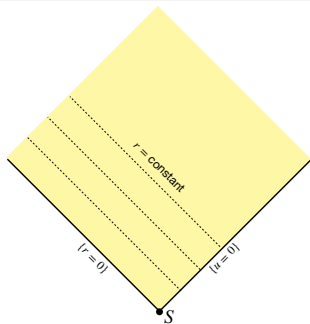
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Embedding a null hypersurface in a vacuum spacetime?

Proposition

In the Isenberg-Moncrief coordinate system the vacuum characteristic initial data $\Psi[\mathcal{N}, k]$ can be reduced to

$$\Phi_{\text{IM}}[\mathcal{N}, k] := \{(\partial_{u^j} g_{AB}, \beta_A)_{0 \leq j \leq k} \text{ on } \mathbf{S} \text{ and } (g_{AB}, \alpha) \text{ on } \mathcal{N}\}, \quad (0.2)$$

where \mathbf{S} is a cross-section of \mathcal{N} . □

Proof: transverse derivatives of the metric on a characteristic hypersurface \mathcal{N} are determined uniquely by the above data through ODEs along the null geodesics threading \mathcal{N} or through algebraic equations.

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The “hand-crank construction”

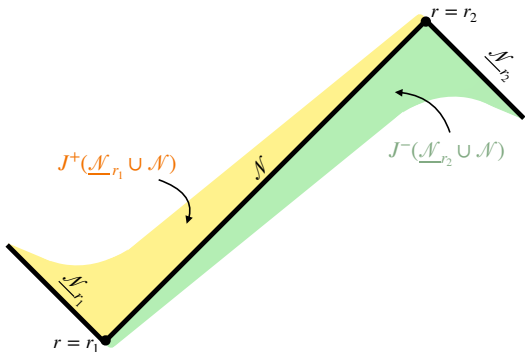


Figure: The “hand-crank construction”.

Corollary: Extending a characteristic future development

A Fledermaus = two cranks

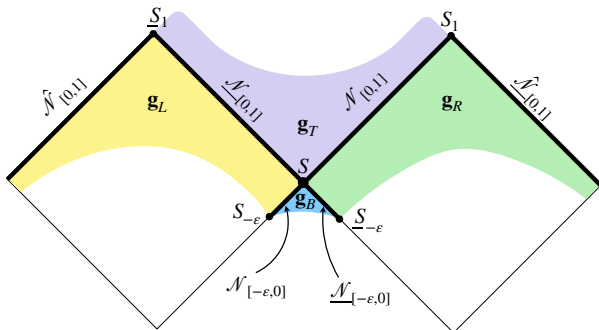


Figure: The Fledermaus.

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Proof:

Use the crank

Extension theorem

Theorem

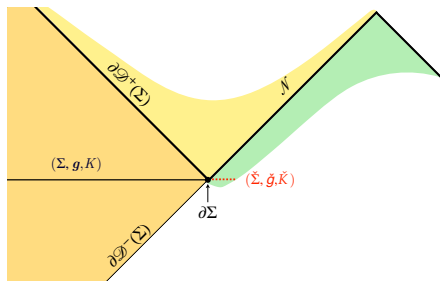
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Corollary: asymptotically flat extensions

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Every vacuum initial data set with spherical boundary and $\theta > 0$ can be extended to an asymptotically flat vacuum initial data set (?)

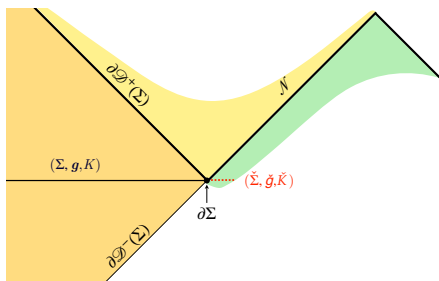
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Proof (?):

use an ACR
gluing to Kerr



Corollary: Embedding a truncated cone

Use the spacelike data extension

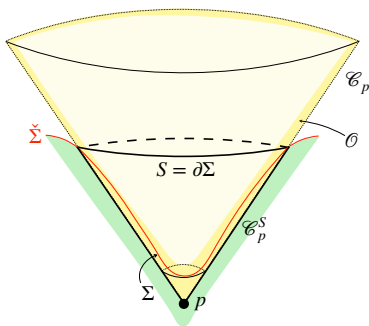


Figure: Extending a vacuum metric on a truncated future cone $J^+(p) \cap J^-(\mathcal{S})$ to a neighborhood thereof.

Corollary: Embedding data at a point

Use the embedding of a cone

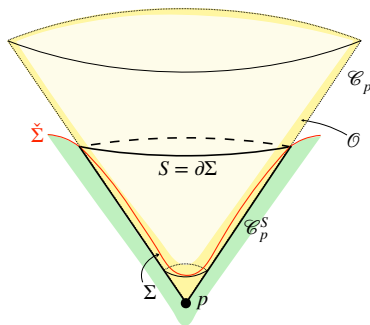


Figure: Extending data at a point