

RECONSTRUCTION OF INCLUSIONS AND CRACKS
IN CALDERÓN'S INVERSE CONDUCTIVITY PROBLEM

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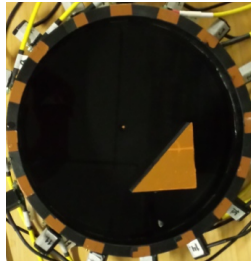
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Electrical Impedance Tomography

Inverse problem of EIT: From several current–voltage measurements on surface electrodes, reconstruct the electrical conductivity distribution of an object.



The continuum model

- $\Omega \subset \mathbb{R}^d$, $d \geq 2$, bounded Lipschitz domain with connected complement
 - $\Gamma \subseteq \partial\Omega$ arbitrarily small open boundary piece
 - Conductivity $\gamma \in L_+^\infty(\Omega) := \{\varsigma \in L^\infty(\Omega; \mathbb{R}) \mid \inf \varsigma > 0\}$
 - Boundary current density f
 - Electric potential $u = u_f^\gamma$
-

Continuum Model

$$\nabla \cdot (\gamma \nabla u) = 0 \quad \text{in } \Omega, \quad \nu \cdot (\gamma \nabla u)|_{\partial\Omega} = \begin{cases} f & \text{on } \Gamma, \\ 0 & \text{on } \partial\Omega \setminus \Gamma. \end{cases}$$

Local Neumann-to-Dirichlet map (current-to-voltage)

$$\Lambda(\gamma) : L_\diamond^2(\Gamma) \rightarrow L_\diamond^2(\Gamma), \quad f \mapsto u_f^\gamma|_\Gamma,$$

is compact and self-adjoint, with \diamond denoting a zero-mean condition on Γ .

Forward and inverse problems

■ Forward problem

- $\Lambda : \gamma \mapsto \Lambda(\gamma)$

■ Calderón's inverse conductivity problem

For which classes of coefficients do we have:

- $\Lambda(\gamma_1) = \Lambda(\gamma_2) \Rightarrow \gamma_1 = \gamma_2?$
- An algorithm to evaluate $\Lambda^{-1} : \Lambda(\gamma) \mapsto \gamma?$

■ Inclusion/obstacle detection

Let $\gamma = \gamma_0 + \gamma_D$ and $D = \text{supp}(\gamma_D)$.

- Can we reconstruct D from γ_0 and $\Lambda(\gamma)$?

Lemma 1 (Kang–Seo–Sheen)

For $f \in L^2_\diamond(\Gamma)$ and $\gamma_1, \gamma_2 \in L^\infty_+(\Omega)$ there are the following estimates:

$$\int_\Omega \frac{\gamma_2}{\gamma_1} (\gamma_1 - \gamma_2) |\nabla u_f^{\gamma_2}|^2 dx \leq \langle (\Lambda(\gamma_2) - \Lambda(\gamma_1))f, f \rangle_{L^2(\Gamma)} \leq \int_\Omega (\gamma_1 - \gamma_2) |\nabla u_f^{\gamma_2}|^2 dx$$

This implies the following intuitive relation between conductivity and power:

$$\gamma_1 \geq \gamma_2 \text{ a.e. in } \Omega \quad \Rightarrow \quad \Lambda(\gamma_2) \geq \Lambda(\gamma_1).$$

Monotonicity-based reconstruction

- **Tamburrino–Rubinacci:** Bounds for inclusions using monotonicity inequalities.
- **Harrach–Ullrich:** For closed set C with connected complement, and for γ piecewise analytic with $-\beta_L \leq \gamma - 1 \leq \beta_U$:

$$\text{supp}(\gamma - 1) \subseteq C \quad \text{if and only if} \quad \Lambda(1 - \beta_L \chi_C) \geq \Lambda(\gamma) \geq \Lambda(1 + \beta_U \chi_C).$$

Extreme inclusions

- For $\varsigma \in L_+^\infty(\Omega)$ and $C = C_0 \cup C_\infty$, let $\sigma = \sigma(\varsigma, C_0, C_\infty)$ denote

$$\sigma = \begin{cases} \varsigma & \text{in } \Omega \setminus C, \\ 0 & \text{in } C_0, \\ \infty & \text{in } C_\infty. \end{cases}$$

- The conductivity equation now becomes:

$$\begin{aligned} \nabla \cdot (\sigma \nabla u) &= 0 \text{ in } \Omega \setminus C, \\ \nu \cdot (\sigma \nabla u) &= \begin{cases} f & \text{on } \Gamma, \\ 0 & \text{on } \partial(\Omega \setminus C_0) \setminus \Gamma, \end{cases} \\ \nabla u &= 0 \text{ in } C_\infty^\circ, \\ \int_{\partial C_i} \nu \cdot (\sigma \nabla u) \, dS &= 0 \text{ for each component } C_i \text{ of } C_\infty. \end{aligned}$$

Convergence of ND maps

Theorem 2

Define the ϵ -truncated version of σ , with $\epsilon > 0$, by

$$\sigma_\epsilon = \begin{cases} \varsigma & \text{in } \Omega \setminus C, \\ \epsilon\varsigma & \text{in } C_0, \\ \epsilon^{-1}\varsigma & \text{in } C_\infty. \end{cases}$$

Then the following estimate holds

$$\|u_f^{\sigma_\epsilon} - u_f^\sigma\|_{H^1(\Omega \setminus C_0)} \leq K\sqrt{\epsilon}\|f\|_{L^2(\Gamma)},$$

with $K > 0$ independent of f and ϵ . As a direct consequence

$$\|\Lambda(\sigma_\epsilon) - \Lambda(\sigma)\|_{\mathcal{L}(L^2_\diamond(\Gamma))} \leq K\sqrt{\epsilon}.$$

Corollary: There is an H^1 -extension of u_f^σ onto the set C_0 , satisfying

$$\|u_f^{\sigma_\epsilon} - u_f^\sigma\|_{H^1(\Omega)} \leq K\sqrt{\epsilon}\|f\|_{L^2(\Gamma)}.$$

Inverse problem for detection of extreme inclusions

- We denote a set of *admissible test inclusions* by

$$\mathcal{A} = \{C \Subset \Omega \mid C \text{ is the closure of an open set,} \\ \text{has connected complement,} \\ \text{and has Lipschitz boundary } \partial C\}.$$

- Let $\gamma_0 \in L_+^\infty(\Omega)$ satisfy the *unique continuation principle* (UCP), and suppose $0 \leq \gamma \leq \infty$ is measurable and $D = \text{supp}(\gamma - \gamma_0) \in \mathcal{A}$.
- **Inverse problem:** Reconstruct D from knowledge of γ_0 and $\Lambda(\gamma)$.
- Some simplifying notation for some $C \in \mathcal{A}$:
 - $\Lambda_C^\emptyset = \Lambda(\sigma(\gamma_0, C, \emptyset))$
 - $\Lambda_\emptyset^C = \Lambda(\sigma(\gamma_0, \emptyset, C))$

General monotonicity method

- Suppose $D = D_0 \cup D_\infty \cup D_- \cup D_+$ is a disjoint union, with $D, D_0, D_\infty \in \mathcal{A}$ and D_\pm are measurable sets.
- Define $0 \leq \gamma \leq \infty$ by

$$\gamma = \begin{cases} 0 & \text{in } D_0, \\ \infty & \text{in } D_\infty, \\ \gamma_- & \text{in } D_-, \\ \gamma_+ & \text{in } D_+, \\ \gamma_0 & \text{in } \Omega \setminus D. \end{cases}$$

- We assume γ satisfies a technical assumption near ∂D (next slide).

Theorem 3 (Candiani–Dardé–Garde–Hyvönen)

For any $C \in \mathcal{A}$, then

$$D \subseteq C \quad \text{if and only if} \quad \Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C.$$

Technical assumption near ∂D

- For each $x \in \partial D$, and every open neighbourhood W of x , there exists a relatively open set $V \subset D$ that intersects ∂D , and $V \subset \tilde{D} \cap W$ for one set $\tilde{D} \in \{D_0, D_\infty, D_-, D_+\}$.
 - If $\tilde{D} = D_-$, there exists an open ball $B \subset V$ such that $\sup_B(\gamma_- - \gamma_0) < 0$.
 - If $\tilde{D} = D_+$, there exists an open ball $B \subset V$ such that $\inf_B(\gamma_+ - \gamma_0) > 0$.
- **In non-technical terms:** The sign of $\gamma - \gamma_0$ cannot change arbitrarily often near any open part of ∂D . And, either a jump from γ_0 or a local strict increase or decrease from γ_0 near ∂D .

Illustration of numerical implementation

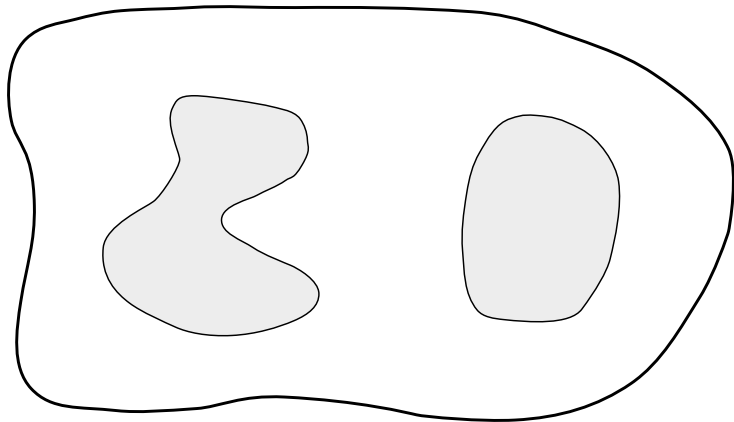


Illustration of numerical implementation

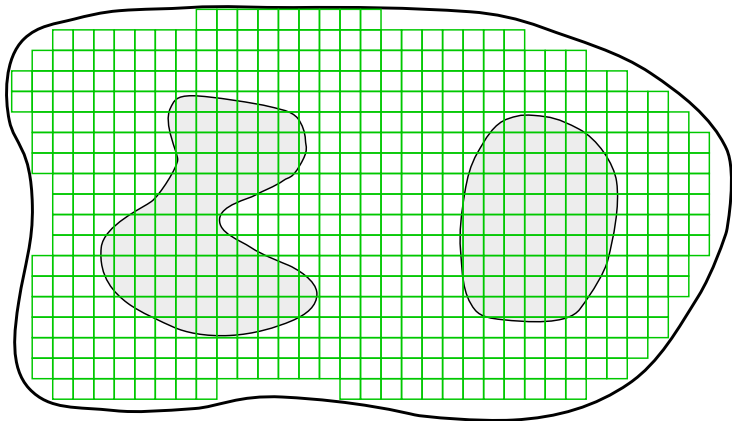


Illustration of numerical implementation

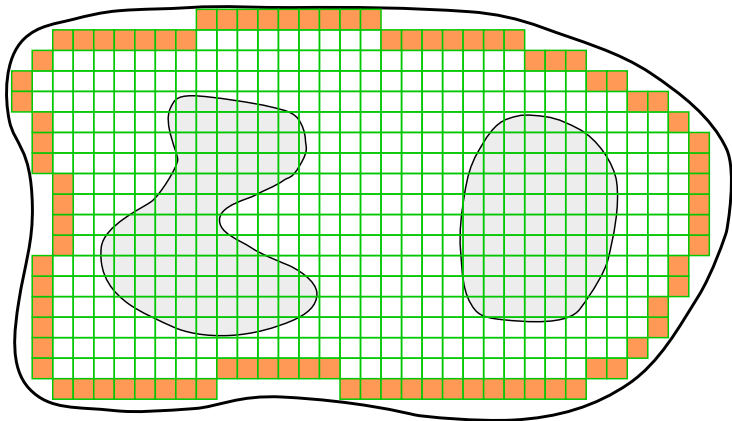


Illustration of numerical implementation

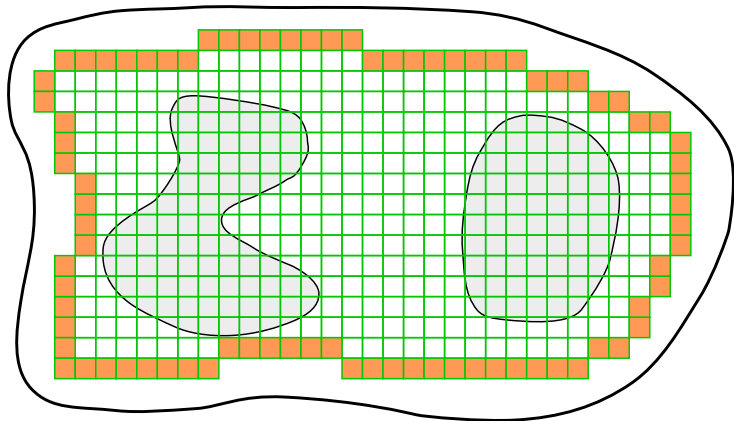


Illustration of numerical implementation

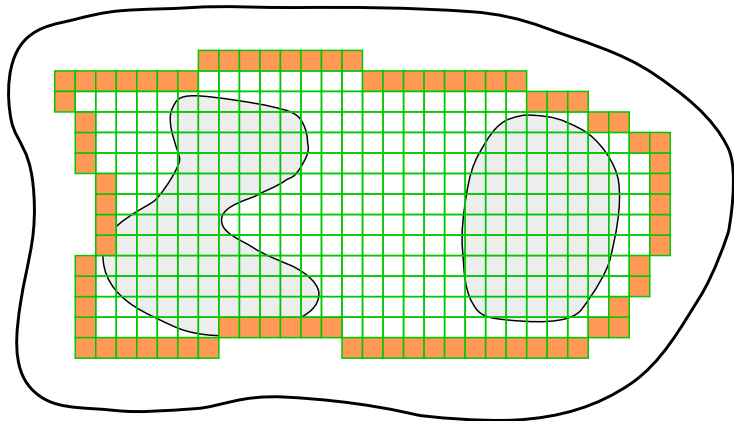


Illustration of numerical implementation

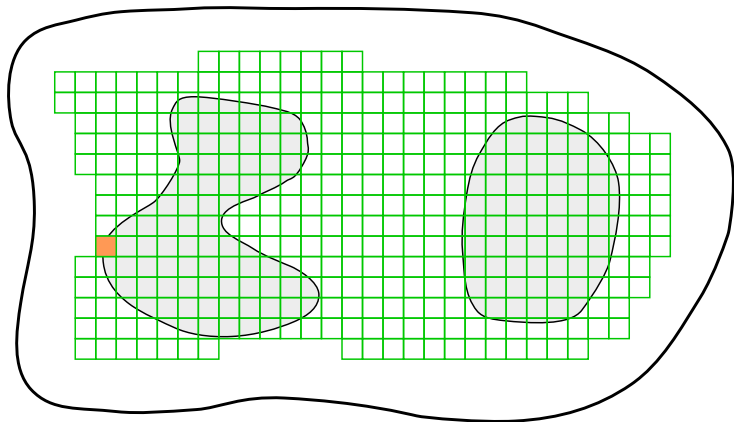


Illustration of numerical implementation

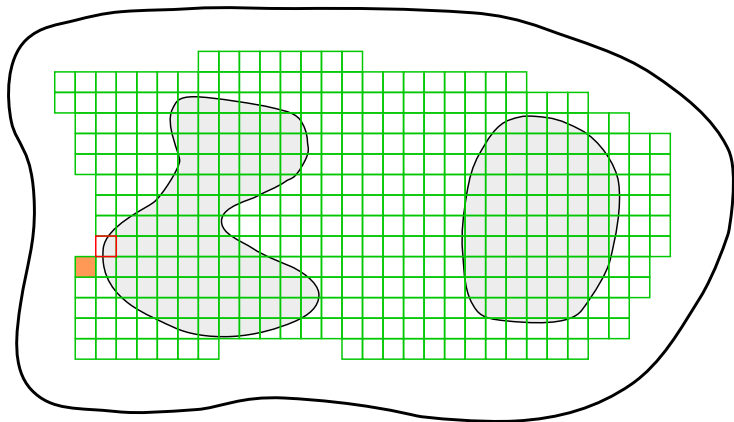


Illustration of numerical implementation

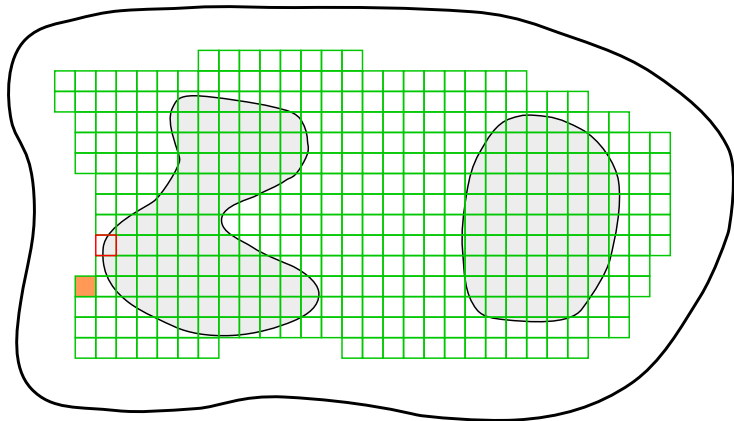
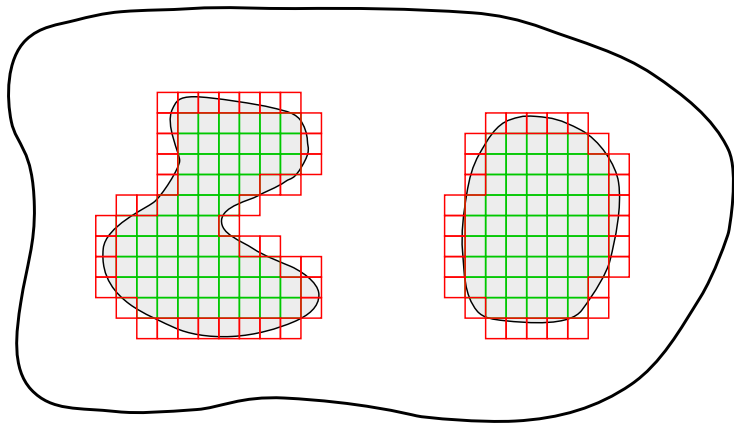


Illustration of numerical implementation



Proof of $D \subseteq C \Rightarrow \Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C$

- Define the ϵ -truncation of γ , with $\epsilon > 0$, as

$$\gamma_\epsilon = \begin{cases} \epsilon\gamma_0 & \text{in } D_0, \\ \epsilon^{-1}\gamma_0 & \text{in } D_\infty, \\ \gamma_- & \text{in } D_-, \\ \gamma_+ & \text{in } D_+, \\ \gamma_0 & \text{in } \Omega \setminus D. \end{cases} \quad \gamma_{\alpha,C} = \begin{cases} \alpha\gamma_0 & \text{in } C, \\ \gamma_0 & \text{in } \Omega \setminus C. \end{cases}$$

- Let $0 < \epsilon_0 < 1$ be small enough that $\epsilon_0\gamma_0 \leq \gamma$ in D_- and $\epsilon_0^{-1}\gamma_0 \geq \gamma$ in D_+ .
- Assume $D \subseteq C$ and $0 < \epsilon \leq \epsilon_0$ then $\gamma_{\epsilon,C} \leq \gamma_\epsilon \leq \gamma_{\epsilon^{-1},C}$. By monotonicity:

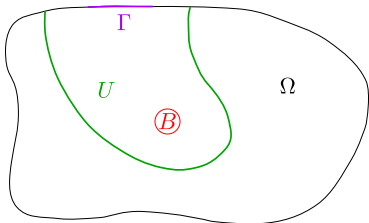
$$\Lambda(\gamma_{\epsilon,C}) \geq \Lambda(\gamma_\epsilon) \geq \Lambda(\gamma_{\epsilon^{-1},C}).$$

Letting $\epsilon \rightarrow 0$ gives $\Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C$.

On the unique continuation principle

For $U \subseteq \bar{\Omega}$ relatively open and connected, we say that $\varsigma \in L_+^\infty(\Omega)$ satisfies the weak **unique continuation principle** (UCP) in U for the conductivity equation if:

- If $\nabla \cdot (\varsigma \nabla v) = 0$ in U° and $v \equiv 0$ in an open set $B \subseteq U$, then $v \equiv 0$ in U .
- If $\nabla \cdot (\varsigma \nabla v) = 0$ in U° with vanishing Cauchy data on $\partial U \cap \Gamma$, then $v \equiv 0$ in U .



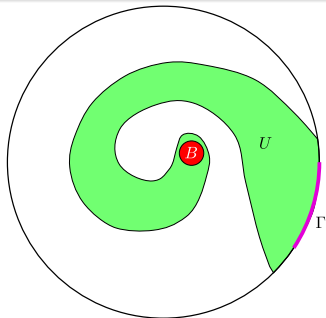
- This is e.g. satisfied for:
 - $d = 2$: L_+^∞
 - $d \geq 3$: $L_+^\infty \cap \text{Lipschitz}$
 - $d \geq 2$: Piecewise analytic (allows discontinuities)

Localised potentials

Lemma 4 (Gebauer)

Let $U \subset \bar{\Omega}$ be a relatively open connected set that intersects Γ . Let $B \subset U$ be an open set and assume $\varsigma \in L_+^\infty(\Omega)$ satisfies the UCP in U . Then there are sequences $(f_i) \subset L_\diamond^2(\Gamma)$ and $(u_i) \subset H_\diamond^1(\Omega)$ with $u_i = u_{f_i}^\varsigma$ such that

$$\lim_{i \rightarrow \infty} \int_B |\nabla u_i|^2 dx = \infty \quad \text{and} \quad \lim_{i \rightarrow \infty} \int_{\Omega \setminus U} |\nabla u_i|^2 dx = 0.$$



Simultaneous localization of power densities

Lemma 5

For $\varsigma \in L_+^\infty(\Omega)$ and $(f_i) \subset L_\diamond^2(\Gamma)$, suppose that $u_i = u_{f_i}^\varsigma$ satisfies

$$\lim_{i \rightarrow \infty} \int_B |\nabla u_i|^2 dx = \infty \quad \text{and} \quad \lim_{i \rightarrow \infty} \int_{\Omega \setminus U} |\nabla u_i|^2 dx = 0.$$

If $C = C_0 \cup C_\infty$ with $C \subset \Omega \setminus \bar{U}$ and $\hat{u}_i = u_{f_i}^\sigma$ with $\sigma = \sigma(\varsigma, C_0, C_\infty)$, then it also holds

$$\lim_{i \rightarrow \infty} \int_B |\nabla \hat{u}_i|^2 dx = \infty \quad \text{and} \quad \lim_{i \rightarrow \infty} \int_{\Omega \setminus U} |\nabla \hat{u}_i|^2 dx = 0.$$

If $\text{supp}(\varsigma_1 - \varsigma_2) \subset \Omega \setminus \bar{U}$ the localisation for ς_1 is transferred to ς_2 (Harrach–Ullrich).

Improved monotonicity principles

Lemma 6

Let $C = C_0 \cup C_\infty$, $\varsigma, \varsigma_1, \varsigma_2 \in L_+^\infty(\Omega)$, and $f \in L_\diamond^2(\Gamma)$.

- *Different background conductivity ($\sigma_1 = \sigma(\varsigma_1, C_0, C_\infty)$ and $\sigma_2 = \sigma(\varsigma_2, C_0, C_\infty)$):*

$$\int_{\Omega \setminus C} \frac{\varsigma_2}{\varsigma_1} (\varsigma_1 - \varsigma_2) |\nabla u_2|^2 dx \leq \langle (\Lambda_2 - \Lambda_1)f, f \rangle_{L^2(\Gamma)} \leq \int_{\Omega \setminus C} (\varsigma_1 - \varsigma_2) |\nabla u_2|^2 dx.$$

- *With and without perfectly conducting inclusions ($\sigma_1 = \sigma(\varsigma, C_0, C_\infty)$ and $\sigma_2 = \sigma(\varsigma, C_0, \emptyset)$):*

$$\int_{C_\infty} \varsigma |\nabla u_2|^2 dx \leq \langle (\Lambda_2 - \Lambda_1)f, f \rangle_{L^2(\Gamma)} \leq K \int_{C_\infty} |\nabla u_2|^2 dx,$$

where $K > 0$ is independent of f .

- *With and without perfectly insulating inclusions ($\sigma_1 = \sigma(\varsigma, \emptyset, C_\infty)$ and $\sigma_2 = \sigma(\varsigma, C_0, C_\infty)$):*

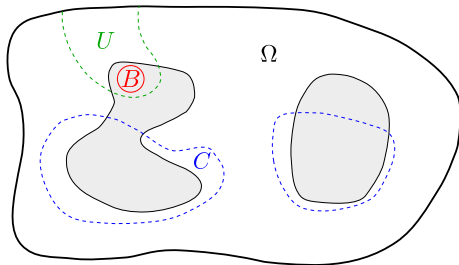
$$\int_{C_0} \varsigma |\nabla u_1|^2 dx \leq \langle (\Lambda_2 - \Lambda_1)f, f \rangle_{L^2(\Gamma)} \leq \int_{C_0} \varsigma |\nabla u_2|^2 dx.$$

Proof of $D \not\subseteq C \Rightarrow \neg(\Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C)$

- Assume $D \not\subseteq C$, i.e. $D \setminus C$ contains an open ball B that can be connected to Γ via a relatively open connected set $U \subset \bar{\Omega}$.
- We may assume that $\bar{U} \cap C = \emptyset$ and either of the following four options holds:

$$\begin{array}{ll} \text{(a): } \bar{U} \cap (D \setminus D_+) = \emptyset, & \text{(b): } \bar{U} \cap (D \setminus D_-) = \emptyset, \\ \text{(c): } \bar{U} \cap (D \setminus D_\infty) = \emptyset, & \text{(d): } \bar{U} \cap (D \setminus D_0) = \emptyset. \end{array}$$

- In the following we consider case (d).



Proof of $D \not\subseteq C \Rightarrow \neg(\Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C)$

- Recall the definition of γ , and now introduce also some new auxiliary conductivities:

$$\gamma := \begin{cases} 0 & \text{in } D_0 \\ \infty & \text{in } D_\infty \\ \gamma_- & \text{in } D_- \\ \gamma_+ & \text{in } D_+ \\ \gamma_0 & \text{in } \Omega \setminus D \end{cases} \quad \gamma_1 := \begin{cases} \gamma_- & \text{in } D_- \\ \gamma_+ & \text{in } D_+ \\ \gamma_0 & \text{in } \Omega \setminus (D_- \cup D_+) \end{cases}$$

and $\gamma_2 := \sigma(\gamma_1, \emptyset, D_\infty)$ and $\gamma_C := \sigma(\gamma_0, C, \emptyset)$.

- We will now estimate each of the following terms:

$$\Lambda_C^\emptyset - \Lambda(\gamma) = [\Lambda_C^\emptyset - \Lambda(\gamma_0)] + [\Lambda(\gamma_0) - \Lambda(\gamma_1)] + [\Lambda(\gamma_1) - \Lambda(\gamma_2)] + [\Lambda(\gamma_2) - \Lambda(\gamma)]$$

- Let (f_i) simultaneously localize potentials $u_{0,i}$, $u_{1,i}$, $u_{2,i}$, and $u_{C,i}$ in B along the set U .

Proof of $D \not\subseteq C \Rightarrow \neg(\Lambda_C^\emptyset \geq \Lambda(\gamma) \geq \Lambda_\emptyset^C)$

- Using the improved monotonicity inequalities:

$$\langle (\Lambda_C^\emptyset - \Lambda(\gamma_0))f_i, f_i \rangle_{L^2(\Gamma)} \leq \sup_C(\gamma_0) \int_C |\nabla u_{C,i}|^2 dx \rightarrow 0$$

$$\langle (\Lambda(\gamma_0) - \Lambda(\gamma_1))f_i, f_i \rangle_{L^2(\Gamma)} \leq \sup_{D_- \cup D_+}(\gamma_1 - \gamma_0) \int_{D_- \cup D_+} |\nabla u_{0,i}|^2 dx \rightarrow 0$$

$$\langle (\Lambda(\gamma_1) - \Lambda(\gamma_2))f_i, f_i \rangle_{L^2(\Gamma)} \leq K \int_{D_\infty} |\nabla u_{1,i}|^2 dx \rightarrow 0$$

$$\langle (\Lambda(\gamma_2) - \Lambda(\gamma))f_i, f_i \rangle_{L^2(\Gamma)} \leq -\inf_{D_0}(\gamma_0) \int_{D_0} |\nabla u_{2,i}|^2 dx \rightarrow -\infty$$

- In total this gives $\lim_{i \rightarrow \infty} \langle (\Lambda_C^\emptyset - \Lambda(\gamma))f_i, f_i \rangle_{L^2(\Gamma)} = -\infty$, i.e. $\Lambda_C^\emptyset \not\geq \Lambda(\gamma)$.

Degenerate and singular inclusions

Extension of the method with Hyvönen:

- A nonnegative function w on \mathbb{R}^d is called an A_2 Muckenhoupt weight, if w and $1/w$ are locally integrable and satisfy

$$\exists C > 0, \forall B \text{ open ball in } \mathbb{R}^d : \left(\int_B w \, dx \right) \left(\int_B \frac{1}{w} \, dx \right) \leq C.$$

- If Σ is a Lipschitz hypersurface, then w can locally behave as

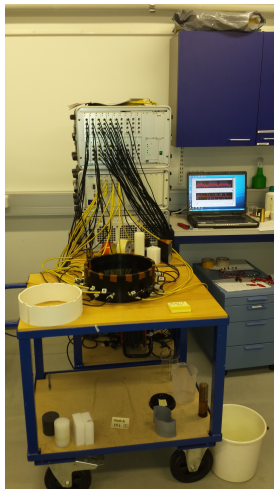
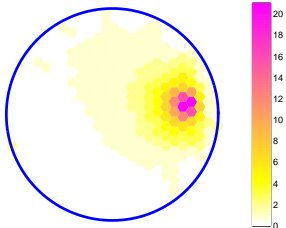
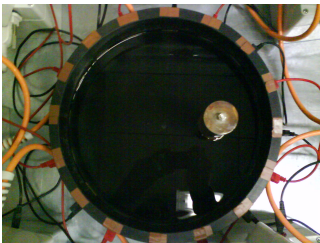
$$\text{dist}(\cdot, \Sigma)^s, \quad s \in (-1, 1).$$

Or near a point x_0 as

$$\text{dist}(\cdot, x_0)^s, \quad s \in (-d, d).$$

- We can allow γ to be the restriction of an A_2 weight in the interior of D , and still recover D with the Monotonicity Method.

Examples - real data



Kuopio impedance tomograph

Lipschitz cracks

Definition 7

A collection of cracks χ lies in the class \mathcal{X} if for some $N \in \mathbb{N}_0$

$$\chi = \bigcup_{i=1}^N \sigma_i,$$

where the $\sigma_i \subset \Omega$ are $(d-1)$ dimensional orientable Lipschitz surfaces with non-empty Lipschitz boundary $\partial\sigma_i$, and with

$$\text{dist}(\sigma_i, \sigma_j) > 0 \text{ for } i \neq j \text{ and } \text{dist}(\sigma_i, \partial\Omega) > 0 \text{ for all } i.$$

We refer to $D \in \mathcal{X}$ as “a (D_0, D_∞) collection of cracks” if:

- $D = D_0 \cup D_\infty$ for $D_0, D_\infty \in \mathcal{X}$,
- $\text{dist}(D_0, D_\infty) > 0$,
- each crack in D_0 is perfectly insulating,
- and each crack in D_∞ is perfectly conducting.

Conductivity problem with cracks

Conductivity problem where D is a (D_0, D_∞) collection of cracks.

$$\begin{aligned} -\nabla \cdot (\gamma_0 \nabla u) &= 0 \quad \text{in } \Omega \setminus \bar{D}, \\ \gamma_0 \frac{\partial u}{\partial \nu} &= \begin{cases} f & \text{on } \Gamma, \\ 0 & \text{on } \partial\Omega \setminus \bar{\Gamma}, \end{cases} \\ \gamma_0 \frac{\partial u}{\partial n} &= 0 \quad \text{on } D_0, \end{aligned}$$

u is locally constant on D_∞ ,

$$\int_{D_i} \left[\gamma_0 \frac{\partial u}{\partial n} \right] dS = 0 \quad \text{for each component } D_i \text{ of } D_\infty.$$

Conductivity problem with cracks

- Weak problem for electric potential $u = u_{D_0, f}^{D_\infty}$:

$$\int_{\Omega} \gamma_0 \nabla u \cdot \nabla v \, dx = \int_{\Gamma} f v|_{\Gamma} \, dS, \quad \forall v \in \mathcal{H}_{D_0}^{D_\infty}.$$

where

$$\mathcal{H}_{D_0}^{D_\infty} = \{v \in H_{\diamond}^1(\Omega \setminus \overline{D_0}) \mid v \text{ is locally constant on } D_\infty\}.$$

Note that ∇u extends to an L^2 -function in all of Ω .

- The local ND map is denoted $\Lambda_{D_0}^{D_\infty}$.

Conductivity problem with cracks

A different way to understand the crack problems:

$$\mathcal{H}_{\emptyset}^{D\infty} \subseteq \mathcal{H}_{D_0}^{D\infty} \subseteq \mathcal{H}_{D_0}^{\emptyset}.$$

And in the inner product $\langle u, v \rangle_* = \int_{\Omega} \gamma_0 \nabla u \cdot \nabla v \, dx$:

- $u_{\emptyset, f}^{D\infty}$ is the orthogonal projection of $u_{D_0, f}^{D\infty}$ onto $\mathcal{H}_{\emptyset}^{D\infty}$.
- $u_{D_0, f}^{D\infty}$ is the orthogonal projection of $u_{D_0, f}^{\emptyset}$ onto $\mathcal{H}_{D_0}^{D\infty}$.

Monotonicity reconstruction of cracks

Theorem 8 (Garde–Vogelius)

Let D be a (D_0, D_∞) collection of cracks. Given any $C \in \mathcal{A}$, then

$$D \subset C \quad \text{if and only if} \quad \Lambda_C^\emptyset \geq \Lambda_{D_0}^{D_\infty} \geq \Lambda_\emptyset^C.$$

Theorem 9 (Garde–Vogelius)

Let $D \in \mathcal{X}$.

- Given any $\chi \in \mathcal{X}$, then

$$\chi \subseteq D \quad \text{if and only if} \quad \Lambda_D^\emptyset \geq \Lambda_\chi^\emptyset.$$

- Given any $\chi \in \mathcal{X}$, then

$$\chi \subseteq D \quad \text{if and only if} \quad \Lambda_\emptyset^\chi \geq \Lambda_\emptyset^D.$$

Monotonicity reconstruction of cracks

More complicated setting:

- No open set inside inclusion to localise in.
- Usual monotonicity inequalities become trivial when collapsing inclusions to zero volume.
- Less general uniqueness results; we now assume $\gamma_0 \in C^2(\overline{\Omega})$ and positive.

Monotonicity reconstruction of cracks

The key idea for the “difficult direction” of the proof, is to construct:

- Sequence of Neumann boundary values (f_n) ,
- Sequence of potentials $\tilde{u}_n = u_{\Sigma_0, f_n}^{\Sigma_\infty}$,
- Sequence of potentials $\hat{u}_n = u_{\emptyset, f_n}^{\Sigma_\infty}$,

such that \tilde{u}_n and \hat{u}_n localise (blow up) in a set intersecting Σ_∞ , and *also* the difference

$$\tilde{u}_n - \hat{u}_n$$

localise in the same way.

Localised potentials

We will need the constructive version of localised potentials.

Lemma 10 (Gebauer '08)

Let H , K_1 , and K_2 be Hilbert spaces, let $A_j \in \mathcal{L}(K_j, H)$ for $j = 1, 2$, and assume that A_2^* is injective. Assume that there exists $y_0 \in R(A_1)$ such that $y_0 \notin R(A_2)$. For $n \in \mathbb{N}$ we define

$$\xi_n = (A_2 A_2^* + \frac{1}{n} I)^{-1} y_0$$

and

$$x_n = \frac{\xi_n}{\|A_2^* \xi_n\|_{K_2}^{3/2}}.$$

Then

$$\lim_{n \rightarrow \infty} \|A_1^* x_n\|_{K_1} = \infty \quad \text{and} \quad \lim_{n \rightarrow \infty} \|A_2^* x_n\|_{K_2} = 0.$$

An auxiliary operator

- Let $V \in \mathcal{A}$.
- Let Σ be a $(\Sigma_0, \Sigma_\infty)$ collection of cracks.

For $F \in L^2(V)^d$ we define $w = w_{\Sigma_0, F}^{\Sigma_\infty} \in \mathcal{H}_{\Sigma_0}^{\Sigma_\infty}$ as the unique solution of:

$$\int_{\Omega} \gamma_0 \nabla w \cdot \nabla v \, dx = \int_V F \cdot \nabla v \, dx, \quad \forall v \in \mathcal{H}_{\Sigma_0}^{\Sigma_\infty}.$$

We define $L_{\Sigma_0}^{\Sigma_\infty}(V) : L^2(V)^d \rightarrow L_\diamond^2(\Gamma)$ as

$$L_{\Sigma_0}^{\Sigma_\infty}(V)F = w_{\Sigma_0, F}^{\Sigma_\infty}|_\Gamma.$$

Then

$$(L_{\Sigma_0}^{\Sigma_\infty}(V))^* f = \nabla u_{\Sigma_0, f}^{\Sigma_\infty}|_V, \quad f \in L_\diamond^2(\Gamma).$$

Lemma 11

Let $V \in \mathcal{A}$ and let Σ be a $(\Sigma_0, \Sigma_\infty)$ collection of cracks.

- If $\Sigma_0 \in V$ then $R(L_{\emptyset}^{\Sigma_\infty}(V)) = R(L_{\Sigma_0}^{\Sigma_\infty}(V))$.
- If $\Sigma_\infty \in V$ then $R(L_{\Sigma_0}^{\emptyset}(V)) = R(L_{\Sigma_0}^{\Sigma_\infty}(V))$.

The proof becomes much more complicated due to different function spaces involved.

Lemma 12

Let Σ be a $(\Sigma_0, \Sigma_\infty)$ collection of cracks. Assume that $\Sigma_0 \in V$ and $\Sigma_\infty \in W$ for $V, W \in \mathcal{A}$ with $\text{dist}(V, W) > 0$.

- If $\Sigma_0 \neq \emptyset$ then there exists a sequence (f_n) in $L^2_\diamond(\Gamma)$ such that

$$\lim_{n \rightarrow \infty} \langle (\Lambda_W^\emptyset - \Lambda_\emptyset^\emptyset) f_n, f_n \rangle = 0,$$

$$\lim_{n \rightarrow \infty} \langle (\Lambda_\emptyset^\emptyset - \Lambda_\emptyset^W) f_n, f_n \rangle = 0,$$

$$\lim_{n \rightarrow \infty} \langle (\Lambda_{\Sigma_0}^{\Sigma_\infty} - \Lambda_\emptyset^{\Sigma_\infty}) f_n, f_n \rangle = \infty.$$

- Analogous result for Σ_∞ .

Key part of the proof

■ Define $A = L_{\Sigma_0}^{\Sigma_\infty}(V) - L_\emptyset^{\Sigma_\infty}(V)$.

■ Then $A^*f = \nabla(u_{\Sigma_0, f}^{\Sigma_\infty} - u_{\emptyset, f}^{\Sigma_\infty})|_V$.

■ By unique continuation and zero mean conditions on Γ :

$$A^*f = 0 \quad \text{if and only if} \quad u_{\Sigma_0, f}^{\Sigma_\infty} = u_{\emptyset, f}^{\Sigma_\infty}.$$

■ From non-invisibility of cracks for full ND map: $A \neq 0$.

■ There is a $g \in R(A) \setminus \{0\}$ such that also $g \in R(L_\emptyset^{\Sigma_\infty}(V)) = R(L_{\Sigma_0}^{\Sigma_\infty}(V))$ but $g \notin R(L_\emptyset^{\Sigma_\infty}(W))$ (last part from usual proof of localised potentials).

■ Using constructive version of localised potentials, there is a sequence (f_n) such that

$$\begin{aligned} \lim_{n \rightarrow \infty} \|(L_\emptyset^{\Sigma_\infty}(W))^* f_n\|_{L^2(W)^d} &= 0, \\ \lim_{n \rightarrow \infty} \|(L_\emptyset^{\Sigma_\infty}(V))^* f_n\|_{L^2(V)^d} &= \infty, \\ \text{and } \lim_{n \rightarrow \infty} \|A^* f_n\|_{L^2(V)^d} &= \infty. \end{aligned}$$

Difficult direction of main result

- To prove

$$\Lambda_C^\emptyset \geq \Lambda_{D_0}^{D_\infty} \geq \Lambda_\emptyset^C \quad \text{implies} \quad D \subset C,$$

we assume the contrapositive, i.e. $D \not\subset C$.

- We have either of two cases:

- (a): There are $V, W \in \mathcal{A}$ with $\text{dist}(V, W) > 0$ and non-empty $\chi \in \mathcal{X}$, such that

$$\chi \subseteq D_0, \quad \chi \in V, \quad C \subseteq W, \quad \text{and} \quad D_\infty \in W.$$

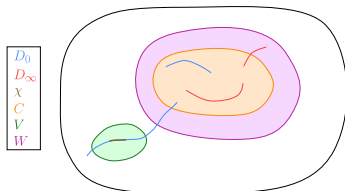
- (b): There are $V, W \in \mathcal{A}$ with $\text{dist}(V, W) > 0$ and non-empty $\chi \in \mathcal{X}$, such that

$$\chi \subseteq D_\infty, \quad \chi \in W, \quad C \subseteq V, \quad \text{and} \quad D_0 \in V.$$

Difficult direction of main result

- Focusing on case (a):

$$\begin{aligned}\Lambda_C^\emptyset - \Lambda_{D_0}^{D_\infty} &= (\Lambda_C^\emptyset - \Lambda_\emptyset^\emptyset) + (\Lambda_\emptyset^\emptyset - \Lambda_\emptyset^{D_\infty}) + (\Lambda_\emptyset^{D_\infty} - \Lambda_{D_0}^{D_\infty}) \\ &\leq (\Lambda_W^\emptyset - \Lambda_\emptyset^\emptyset) + (\Lambda_\emptyset^\emptyset - \Lambda_\emptyset^W) + (\Lambda_\emptyset^{D_\infty} - \Lambda_X^{D_\infty}).\end{aligned}$$



- From our lemma, there is a sequence (f_n) so that

$$\begin{aligned}\lim_{n \rightarrow \infty} \langle (\Lambda_W^\emptyset - \Lambda_\emptyset^\emptyset) f_n, f_n \rangle &= \lim_{n \rightarrow \infty} \langle (\Lambda_\emptyset^\emptyset - \Lambda_\emptyset^W) f_n, f_n \rangle = 0 \\ \lim_{n \rightarrow \infty} \langle (\Lambda_\emptyset^{D_\infty} - \Lambda_X^{D_\infty}) f_n, f_n \rangle &= -\infty.\end{aligned}$$

- Hence $D \not\subset C$ implies $\Lambda_C^\emptyset \not\geq \Lambda_{D_0}^{D_\infty}$.

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