# Hydrodynamics of the $N\mbox{-}{\sf BBM}$ process

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Illustration by Eric Brunet

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Brunet and Derrida N branching particles in  $\mathbb{R}$  with selection: discrete time.

Particle at x dies and creates random offsprings around x.

Select the rightmost  $\boldsymbol{N}$  particles.

iterate

Pascal Maillard studied the N-BBM process.

N particles move as independent Brownian motions in  $\mathbb{R}$ ,

each particle, at rate 1, creates a new particle at its current position.

At each branching time, the left-most particle is removed.

The number N of particles is then conserved.

Brunet Derrida (1997) Shift in the velocity of a front due to a cutoff PRE

Brunet, Derrida, Mueller, Munier (2006). Noisy traveling waves: effect of selection on genealogies. EPL + (06) + (07)

Bérard, Gouéré (2010) Brunet-Derrida behavior of branching-selection particle systems on the line CMP.

Bérard, Maillard (2014) The limiting process of N-BRW with polynomial tails EJP.

Durrett, Remenik (2011) Brunet-Derrida particle systems, free boundary problems and Wiener-Hopf equations AOP.

Derrida, Shi (2017) Large deviations for the BBM in presence of selection or coalescence Preprint.

Julien Berestycki, Brunet, Derrida (2017) Exact solution and precise asymptotics of a Fisher-KPP type front ArXiv

#### Hydrodynamics

Density  $\rho$  with left boundary  $L_0 = \arg \max_a \int_a^\infty \rho(x) dx > -\infty$ 

Time zero: iid continuous random variables with density  $\rho$ .

 $X_t :=$  set of positions of N-BBM particles at time t.

# Theorem 1. [Existence]

For every  $t \ge 0$ , there is a density function  $\psi(\cdot, t) : \mathbb{R} \to \mathbb{R}^+$  such that,

$$\lim_{N \to \infty} \frac{|X_t \cap [a, \infty)|}{N} = \int_a^\infty \psi(r, t) dr, \quad \text{a.s. and in } L^1.$$

for any  $a \in \mathbb{R}$ .

#### Free boundary problem.

Density  $\rho$  with left boundary  $L_0 = \arg \max_a \int_a^\infty \rho(x) dx > -\infty$ 

Find  $((u(r,t), L_t) : r \in \mathbb{R}, t \in [0,T])$  such that:  $u_t = \frac{1}{2}u_{rr} + u, \quad \text{in } (L_t, +\infty);$   $u(r,0) = \rho(r);$  $u(L_t,t) = 0, \quad \int_{L_t}^{\infty} u(r,t)dr = 1.$ 

If one finds a continuous function  $L_t$  such that

$$e^t P(L_s \le B_s^{\rho}, 0 \le s \le t) = 1, \qquad t \ge 0.$$

where  $B_s^{\rho}$  is BM with random initial position  $B_0^{\rho} \sim \rho$ , then

$$\int \varphi(r)u(r,t)dr = e^t E(\varphi(B_t^{\rho}) \mathbf{1}\{L_s \le B_s^{\rho}, 0 \le s \le t\})$$

**Theorem 2.** If  $L_t$  is a continuous function such that

 $((u(r,t), L_t) : t \in [0,T]$ 

is a solution of the free boundary problem, then the hydrodynamic limit  $\psi$  coincides with  $u{:}$ 

$$\psi(\cdot, t) = u(\cdot, t), \quad t \in [0, T].$$
(1)

Lee (2017) proved that if  $\rho \in C_c^2([L_0,\infty))$  and  $\rho'_{L_0} = 2$  then there exist T > 0 and a solution (u, L) of the free boundary problem with the following properties:

- $\{L_t : t \in [0,T]\}$  is in  $C^1[0,T]$ ,  $L_{t=0} = L_0$
- $u \in C(D_{L,T}) \cap C^{2,1}(D_{L,T})$ , where  $D_{L,T} = \{(r,t) : L_t < r, 0 < t < T\}$ .

#### **General strategy**

We use a kind of Trotter-Kato approximation as upper and lower bounds.

Durrett and Remenik upperbound for the Brunet-Derrida model. Leftmost particle motion is *increasing*: natural lower bounds.

Upper and lower bounds method was used in several papers:

- De Masi, F and Presutti (2015) Symmetric simple exclusion process with free boundaries. PTRF

• Carinci, De Masi, Giardinà, and Presutti (2016) Free boundary problems in PDEs and particle systems. SpringerBriefs in Mathematical Physics.

We introduce labelled versions of the processes and a coupling of trajectories to prove the lowerbound. **Ranked BBM, a tool** Let  $(Z_0^1, \ldots, Z_0^N)$  BBM initial positions.  $B_0^{i,1} = Z_0^i$ , iid with density  $\rho$ .  $N_t^i$ : is the size of the *i*th BBM family.

 $B_t^{i,j}:$  is the j-th member of the i-th family at time  $t,\ 1\leq j\leq N_t^i.$  birth-time order.

BBM: 
$$Z_t = \{B_t^{i,j} : 1 \le j \le N_t^i, 1 \le i \le N\}$$

 $B^{i,j}_{[0,t]}$  trajectory coincides with ancestors before birth.

(i, j) is the rank of the *j*th particle of *i*-family

**N-BBM as subset of BBM** 

Let  $X_0 = Z_0$ ,  $\tau_0 = 0$ 

 $\tau_n$  branching times of BBM.

$$X_t := \{ B_t^{i,j} : B_{\tau_n}^{i,j} \ge L_{\tau_n}, \text{ for all } \tau_n \le t \}$$

 $L_{\tau_n} :=$  defined iteratively such that  $|X_t| = N$  for all t

 $X_t$  has the law of N-BBM.

# Stochastic barriers.

Fix  $\delta>0$ 

 $X_0^{\delta,\pm} = Z_0.$ 

The upper barrier. Post-selection at time  $k\delta$ .

$$\begin{split} X^{\delta,+}_{k\delta} &:= N \text{ rightmost } \{B^{i,j}_{k\delta} : B^{i,j}_{(k-1)\delta} \in X^{\delta,+}_{(k-1)\delta} \} \\ L^{\delta,+}_{k\delta} &:= \min X^{\delta,+}_{k\delta} \end{split}$$

The number of particles in  $X_{k\delta}^{\delta,+}$  is exactly N for all k.

The lower barrier.

Pre selection at time  $(k-1)\delta$ .

Select maximal number of rightmost particles at time  $(k-1)\delta$  keeping no more than N particles at time  $k\delta$ .

$$\begin{split} L^{\delta,-}_{(k-1)\delta} &:= \text{cutting point at time } (k-1)\delta \\ X^{\delta,-}_{k\delta} &:= \{B^{i,j}_{k\delta} : B^{i,j}_{(k-1)\delta} \in X^{\delta,-}_{(k-1)\delta} \cap [L^{\delta,-}_{(k-1)\delta},\infty)\} \end{split}$$

Only entire families of particles at time  $(k-1)\delta$  are kept at time  $k\delta$ .

The number of particles in  $X_{k\delta}^{\delta,-}$  is N - O(1).









#### Mass transport partial order

 $X \preccurlyeq Y \quad \text{if and only if} \quad |X \cap [a,\infty)| \leq |Y \cap [a,\infty)| \quad \forall a \in \mathbb{R}.$ 

**Proposition 3.** Coupling  $((\hat{X}_{k\delta}^{\delta,-}, \hat{X}_{k\delta}, \hat{X}_{k\delta}^{\delta,+}) : k \ge 0)$  such that

$$\hat{X}_{k\delta}^{\delta,-} \preccurlyeq \hat{X}_{k\delta} \preccurlyeq \hat{X}_{k\delta}^{\delta,+}, \quad k \ge 0.$$

 $\hat{X}_t^{\delta,-}$  is a subset of  $\hat{Z}_t$ , a BBM with the same law as  $Z_t$ .

**Deterministic barriers.**  $u \in L^1(\mathbb{R}, \mathbb{R}_+)$ .

Gaussian kernel: 
$$G_t u(a) := \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi t}} e^{-(a-r)^2/2t} u(r) dr.$$

$$e^t G_t \rho$$
 solves  $u_t = \frac{1}{2}u_{rr} + u$  with initial  $\rho$ .

Cut operator  $C_m$  is defined by

$$C_m u(a) := u(a) \mathbf{1} \bigg\{ \int_a^\infty u(r) dr \le m \bigg\},$$

so that  $C_m u$  has total mass  $m \wedge ||u||_1$ .

For  $\delta > 0$  and  $k \in \mathbb{N}$ , define the upper and lower barriers:

$$\begin{split} S_0^{\delta,\pm}\rho &:= \rho \quad \text{Initial condition} \\ S_{k\delta}^{\delta,+}\rho &:= \left(C_1 \left(e^{\delta}G_{\delta}\right)\right)^k \rho \quad (\text{diffuse \& grow}) + \text{cut}; \\ S_{k\delta}^{\delta,-}\rho &:= \left(\left(e^{\delta}G_{\delta}\right)C_{e^{-\delta}}\right)^k \rho \quad \text{cut} + (\text{diffuse \& grow}) \end{split}$$

We have  $\|S_{k\delta}^{\delta,\pm}\rho\|_1 = \|\rho\|_1 = 1$  for all k.

Hydrodynamics of  $\delta$ -barriers

We prove that for fixed  $\delta$ 

the stochastic barriers converge to the macroscopic barriers:

**Theorem 4.** Conditions of Theorem 1 and fixed  $\delta$ :

$$\lim_{N \to \infty} \frac{\left| X_{k\delta}^{\delta, \pm} \cap [r, \infty) \right|}{N} = \int_{a}^{\infty} S_{k\delta}^{\delta, \pm} \rho, \quad \text{a.s. and in } L^{1}.$$

The same is true for the coupling marginals  $\hat{X}_{k\delta}^{\delta,\pm}$ .

#### Convergence of macroscopic barriers

Partial order: Take  $u, v : \mathbb{R} \to \mathbb{R}^+$  and denote

$$u \preccurlyeq v \quad \text{iff} \quad \int_a^\infty u \leq \int_a^\infty v \quad \forall a \in \mathbb{R}.$$

Fix t and take diadic  $\delta = t2^{-n}$ . We prove

- $S_t^{\delta,-}\rho$  is increasing and  $S_t^{\delta,+}\rho$  decreasing in n (diadics).
- $\left\|S_t^{\delta,+}\rho S_t^{\delta,-}\rho\right\|_1 \le c\delta.$
- There exists a continuous function  $\psi$  such that for any t > 0,

$$\lim_{n \to \infty} \|S_t^{\delta, \pm} \rho - \psi(\cdot, t)\|_1 = 0.$$

# Sketch of proof of Theorem 1

By coupling  $\hat{X}_t^{\delta,-} \preccurlyeq \hat{X}_t \preccurlyeq \hat{X}_t^{\delta,+}$ .

Convergences in the sense of the Theorem 1:

 $N \to \infty$ :

The stochastic barriers  $\hat{X}_t^{\delta,\pm}$  converge to the macroscopic barriers  $S_t^{\delta,\pm}.$ 

 $\delta \rightarrow 0$ :

The macroscopic barriers converge to a function  $\psi$ , along diadics  $\delta \rightarrow 0$ .

Corollary: N-BBM  $\hat{X}_t$  converge to  $\psi$  as  $N \to \infty$ .

This is Theorem 1.

#### Sketch of proof of Theorem 2

We show that for continuous  $L_t$ , the solution u of the free boundary problem is in between the barriers:

$$S_{k\delta}^{\delta,-}\rho \preccurlyeq u(\cdot,k\delta) \preccurlyeq S_{k\delta}^{\delta,+}\rho.$$

Here we use the Brownian representation of the solution.

# Proof of Pre-selection inequalities.

Rank order

$$(i,j) \prec (i',j')$$
 if and only if  $B_0^{i,1} < B_0^{i',1}$  or  $i = i'$  and  $j < j'$ . (2)

N rank-selected  $\mathsf{BBM}$ :

$$Y_t := \{ B_t^{i,j} : |\{ B_t^{i',j'} : (i,j) \prec (i',j')\}| < N \},\$$

We have  $X^{\delta,-}_{\delta} \subset Y_{\delta}$ , which in turn implies

 $X_{\delta}^{\delta,-} \preceq Y_{\delta}.$ 



Labeled N-BBM.

$$(X_t^1,\ldots,X_t^N) \in \mathbb{R}^N$$

 $X_t^{\ell}$  is just a labelling of *N*-BBM as function of  $(B_{[0,t]}^{i,j}:i,j)$ : When one of the Brownian particles branches at time *s*, identify  $X_{s-}^n :=$  the branching particle  $X_{s-}^m :=$  the position of the leftmost particle (to be erased) At time *s* put

 $X^m_{\circ} = X^n_{\circ}$ 

 $X_t^m$  will follow the newborn Brownian particle until next branching.

Labeled rank-selected N-BBM.

$$((Y_t^1, \sigma_t^1), \dots, (Y_t^N, \sigma_t^N)) \in (\mathbb{R} \times \mathbb{N}^2)^N$$

 $Y_t^{\ell}$  is a labelling of the rank-selected N-BBM  $Y_t$ .

 $\sigma_t^\ell$  tracks the rank of the  $Y^\ell\text{-particles}$  in the Y-tree.

When one of the Brownian particles branches at time s, identify

$$Y_{s-}^n :=$$
 the branching particle,  $\sigma_{s-}^n = (i,j)$ 

$$Y_{s-}^h :=$$
 lowest ranked Y-particle (to be erased)

At time s put

 $Y^h_s=Y^n_{s-}$  and this particle will follow now the newborn Brownian particle  $\sigma^h_s=(i,M^i_{s-}+1)$  (youngest new element of the  $Y^i$  branching family)

Coupling.  $(X_t^1, \ldots, X_t^N), ((Y_t^1, \sigma_t^1) \ldots (Y_t^N, \sigma_t^N))$ Between branchings  $X_t^{\ell} - Y_t^{\ell}$  and  $\sigma_t^{\ell}$  are constant. *s* branching time for *X* process.

 $X_{s-}^n$  and  $Y_{s-}^n$  branching particles.

 $\sigma_{s-}^n=(i,j)$  rank of Y-branching particle

 $X_{s-}^m :=$  leftmost X-particle (to be erased).

 $Y_{s-}^h :=$  lowest-rank *Y*-particle (to be erased).

At time s put

 $X_{s}^{m} = X_{s-}^{n}, Y_{s}^{h} = Y_{s-}^{m}, Y_{s}^{m} = Y_{s-}^{n}$ 

 $X^m_s$  and  $Y^m_s$  will follow now the (same) newborn Brownian particle  $\sigma^h_s=\sigma^n_{s-},\,\sigma^m_s=(i,M^i_{s-}+1)$  (youngest new element of the branching family)



Relative positions of particles at branching time s.



Coupling between  $\underline{x}(t)$  and  $(\underline{y}(t), \underline{\sigma}(t))$ . When n = m only the *h*-th *Y*-particle jumps to  $Y_{s-}^n$ . When n = h only the *m*-th *X*-particle jumps to  $X_{s-}^h$ . Perform two cases simultaneously. The coupling satisfies

$$Y_t^\ell \le X_t^\ell, \quad \text{for all } t, \ell.$$

Hence

$$\hat{X}^{\delta,-}_{\delta} \preceq \hat{Y}_t \preceq X_t$$
, a.s.

 $M_t^i :=$  size of  $Y_0^i$  family at time t.





#### The post-selection process

N-BBM  $X_t$  is a subset of the BBM  $Z_t$ .

 $X_{\delta}^{\delta,+} = N$  right-most Z-particles at time  $\delta$ . Hence,

$$X_{\delta} \preceq X_{\delta}^{\delta,+}$$

#### **Domination** We have proven the dominations

$$\hat{X}_{k\delta}^{\delta,-} \preceq X_{k\delta} \preceq \hat{X}_{k\delta}^{\delta,+}.$$

for k = 1. Iterate to obtain the same for all k.

Construct the coupling for each time interval and then the Brownian tree  $\hat{B}$  containing  $\hat{Y}_{k\delta} \supset \hat{X}_{k\delta}^{\delta,-}$ .

Similarly construct Brownian tree containing  $\hat{X}_{k\delta}^{\delta,+}$ .

Construct new BBM process  $\hat{B}_{[0,t]}$  by

Attaching independent BBM to loose branches of  $\hat{Y}_t$ .

**Proposition 5.**  $\hat{B}_{[0,t]}$  has the same law as the BBM  $B_{[0,t]}$  and

$$\hat{Y}_t := \{ \hat{B}_t^{i,j} : \left| \{ \hat{B}_t^{i',j'} : (i,j) \prec (i',j') \} \right| < N \}$$

is the rank selected process associated to  $\hat{B}_{[0,t]}$ .

The rightmost families with up to N total particles coincide

$$\hat{N}_t^i = M_t^i \quad \text{if} \quad \sum_j N_t^j \mathbf{1}\{B_0^{j,1} \ge B_0^{i,1}\} \le N$$

# Hydrodynamic limit for the barriers

# Macroscopic left boundaries

For  $\delta>0$  and  $\ell\leq k$  denote

$$L_{\ell\delta}^{\delta,+} := \sup_{r} \left\{ \int_{-\infty}^{r} S_{\ell\delta}^{\delta,+} \rho(r') dr' = 0 \right\};$$
  
$$L_{\ell\delta}^{\delta,-} := \sup_{r} \left\{ \int_{-\infty}^{r} S_{\ell\delta}^{\delta,-} \rho(r') dr' < 1 - e^{-\delta} \right\}.$$
(1)

Brownian representation of macroscopic barriers:

 $B_{[0,t]} = (B_s : s \in [0,t])$  Brownian motion with

 $B_0$ , random variable with density  $\rho$ .

**Lemma 6.** For test function  $\varphi \in L^{\infty}(\mathbb{R})$  and t > 0,

$$\int \varphi S_{k\delta}^{\delta,+} \rho = e^{k\delta} E[\varphi(B_{k\delta}) \mathbf{1} \{ B_{\ell\delta} > L_{\ell\delta}^{\delta,+} : 1 \le \ell \le k \}].$$
$$\int \varphi S_{k\delta}^{\delta,-} \rho = e^{k\delta} E[\varphi(B_{k\delta}) \mathbf{1} \{ B_{\ell\delta} > L_{\ell\delta}^{\delta,-} : 0 \le \ell \le k-1 \}].$$

Generic LLN over trajectories of BBM

Let  $B_0^{i,1}$  iid with density  $\rho$ .

 $N_t^i$  size at time t of the i-th BBM family.  $EN_t^i = e^t$ .

**Proposition 7.** Let g be bounded. Then

$$\mu_t^N g := \frac{1}{N} \sum_{i=1}^N \sum_{j=1}^{N_t^i} g(B_{[0,t]}^{i,j}) \underset{N \to \infty}{\longrightarrow} e^t Eg(B_{[0,t]}), \quad \text{a.s. and in } L^1.$$
 (0)

a.s. and in  $L^1$ .

Proof. By the many-to-one Lemma we have

$$E\mu_t^N g = EN_t Eg(B_{[0,t]}) = e^t Eg(B_{[0,t]}), \tag{1}$$

The variance of  $\mu_t^N g$  is order 1/N, by family independence.

Corollary 8 (Hydrodynamics of the BBM).

$$\lim_{N \to \infty} \frac{1}{N} \sum_{i=1}^{N} \sum_{j=1}^{N_t^i} \varphi(B_t^{i,j}) = e^t E \varphi(B_t) \quad \text{a.s. and in } L^1.$$
$$= e^t \int \varphi(r) G_t \rho(r) dr, \qquad (2)$$

#### Proof of Hydrodynamics for barriers

**Proof of Theorem 4** BBM representation of stochastic barriers:

$$\begin{aligned} \pi_{k\delta}^{N,\delta,+}\varphi &= \frac{1}{N}\sum_{i=1}^{N}\sum_{j=1}^{N_{k\delta}^{i}}\varphi(B_{k\delta}^{i,j})\mathbf{1}\{B_{\ell\delta}^{i,j} \geq \underline{L}_{\ell\delta}^{N,\delta,+}: 1 \leq \ell \leq k\}\\ \pi_{k\delta}^{N,\delta,-}\varphi &= \frac{1}{N}\sum_{i=1}^{N}\sum_{j=1}^{N_{k\delta}^{i}}\varphi(B_{k\delta}^{i,j})\mathbf{1}\{B_{\ell\delta}^{i,j} \geq \underline{L}_{\ell\delta}^{N,\delta,-}: 0 \leq \ell \leq k-1\}. \end{aligned}$$

We show that as  $N \to \infty$ 

 $L_{\ell\delta}^{N,\delta,\pm}$  can be replaced by  $L_{\ell\delta}^{\delta,\pm}$ , and get the result by the generic LLN. Use the fact that the random left boundaries are exact quantiles of 1.

# **Proof of Theorem 2** The limit function $\psi$ is the solution of the free boundary problem.

The local solution of the free boundary problem is in between the barriers:

**Theorem 9.** Let  $t \in (0, T]$ ,  $\delta \in \{2^{-n}t, n \in \mathbb{N}\}$ . Then

$$S_t^{\delta,-}\rho \preccurlyeq u(\cdot,t) \preccurlyeq S_t^{\delta,+}\rho, \qquad t=k\delta$$

The upperbound is immediate. The lower bound reduces to show the following stochastic order between conditioned probability measures:

$$P_{u_0}(B_t \ge r | \tau^L \le \delta) \le P_{u_1}(B_t \ge r | \tau^L > \delta)$$
(3)

where  $u_1 = C_{e^{-\delta}} u$ ,  $u_0 = u - u_1$ ,

$$P_{u_i}(B_t \in A) := \frac{1}{\|u_i\|_1} \int u_i(x) P_x(B_t \in A) dx.$$
(4)

and  $\tau^L$  is the hitting time of the boundary.

**Stationary** *N***-BBM**  $X_t$  be *N*-BBM. Process as seen from leftmost particle:

$$X'_t := \{x - \min X_t : x \in X_t\}$$

Durrett and Remenik for a related Brunet-Derrida process proved:

**Theorem 10.** N-BBM as seen from leftmost particle is Harris recurrent.

 $\nu_N$  unique invariant measure. Under  $\nu_N$  asymptotic speed

$$\alpha_N = (N-1)\,\nu_N\big[\min(X\setminus\{0\})\big],\,$$

 $X_t'$  starting with anything converges in distribution to  $u_N$  and

$$\lim_{t \to \infty} \frac{\min X_t}{t} = \alpha_N.$$

 $\alpha_N$  converges to asymptotic speed of the first particle in BBM:

$$\lim_{N \to \infty} \alpha_N = \sqrt{2}.$$
 Berard and Gouéré

**Travelling wave solutions**  $u(r,t) = w(r - \alpha t)$ , where w must satisfy

$$\frac{1}{2}w'' + \alpha w' + w = 0, \quad w(0) = 0, \quad \int_0^\infty w(r)dr = 1.$$

Groisman and Jonckheere (2013): for each speed  $\alpha \geq \sqrt{2}$  there is a solution  $w_{\alpha}$ 

$$w_{\alpha}(r) = \begin{cases} M_{\alpha} r e^{-\alpha r} & \text{if } \alpha = \sqrt{2} \\ M_{\alpha} e^{-\alpha r} \sinh\left(r\sqrt{\alpha^2 - 2}\right) & \text{if } \alpha > \sqrt{2} \end{cases}$$
(4)

where  $M_{\alpha}$  is a normalization constant.

 $w_{\alpha}$  is the unique qsd for Brownian motion with drift  $-\alpha$  and absorption rate 1; see Martínez and San Martín (1994).

*Open problems.* (1) Let  $X_t$  be the *N*-BBM process with initial configuration sampled from the stationary measure  $\nu^N$ .

Show that the empirical distribution of  $X_t$  converges to a measure with density  $w_{\sqrt{2}}(t\sqrt{2}+\cdot)$ , as  $N \to \infty$ . This would be a strong selection principle for N-BBM.

Problem: we do not control the particle-particle correlations in the  $\nu_N$  distributed initial configuration.

If we start with independent particles with distribution  $w_{\sqrt{2}}$ , then Theorem 1 and  $w_{\sqrt{2}}(t\sqrt{2}+\cdot)$  strong solution of FBP imply convergence of the empirical measure to this solution.

(2) "Yaglom limit"? Does  $u(\cdot - L_t, t)$  converges to  $w_{\alpha}$  for some  $\alpha \ge \sqrt{2}$ ? Fix  $\alpha$ , which conditions must satisfy  $\rho$  to converge to  $w_{\alpha}$ ?

(3) Give a simple proof of existence of the solution for the FBP.