

Mean-field limit of systems of topologically interacting particles

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joint works with Adrien Blanchet
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1. Topological interactions
2. Smooth rank-based dynamics
3. Nearest neighbor
4. Conclusion

1. Topological interactions

System with locally interacting agents

Emergence of spatio-temporal coordination

Patterns, structures, correlations, synchronization

No leader



Mean-field (metric) interaction

Particle interact with all particles within a certain distance

Examples: Alignment (Vicsek ...)

Consensus (Cucker-Smale, Motsch-Tadmor, ...)

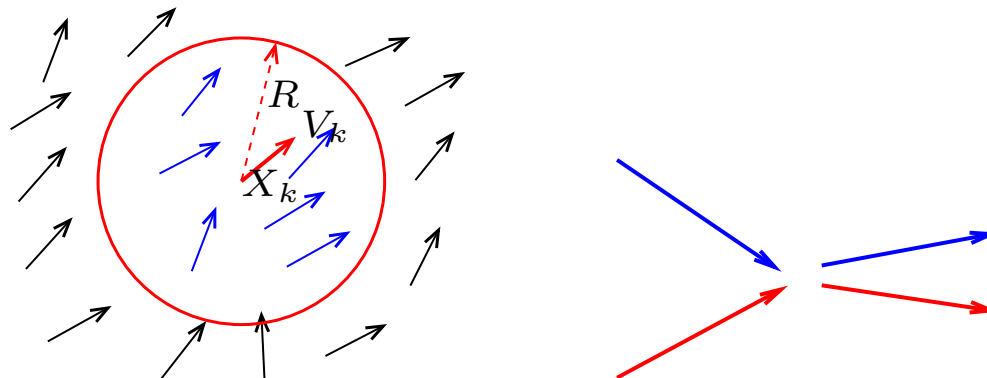
Attraction-repulsion (Bertozzi, Carrillo, ...)

Binary (ternary, ...)

Particle interacts with a partner at contact

Examples: Hard-sphere collisions (Boltzmann ...)

Alignment (Bertin, Droz & Grégoire, ...)

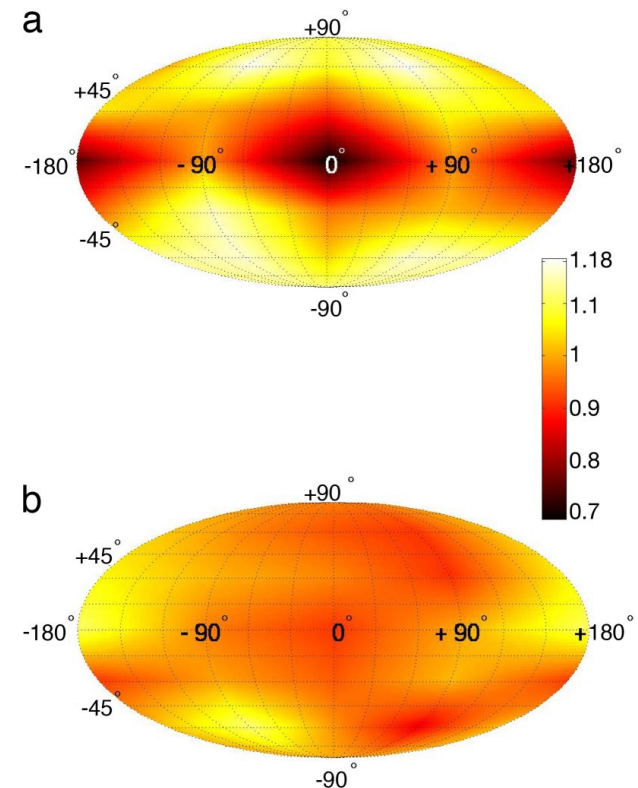
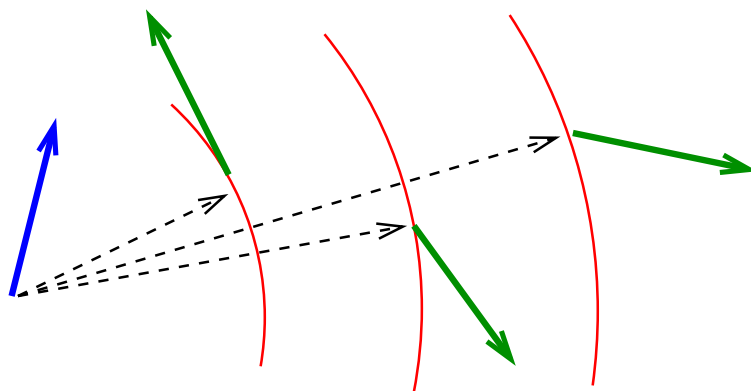


Particles interact with other particles
according to their rank
[Ballerini et al, PNAS 2008]

Interaction ruling animal collective behavior depends on topological rather than metric distance: Evidence from a field study

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Kinetic models

Largely investigated for mean-field metric
or binary interactions

Kinetic models for topological interactions are recent:

J. Haskovec [Physica D 2013], Benedetto et al [arxiv 2021]
for Cucker-Smale type interactions

Y. Brenier
for a competition model

Our work: Boltzmann approach

Adapting 'Choose the Leader' model from
[Carlen, D., Wennberg, M3AS 2013],
[Carlen, Chatelin, D., Wennberg, Physica D 2013]

2. Smooth rank-based dynamics

A. Blanchet & P. D., J. Stat. Phys **163** (2016) 41-60

P. D. & M. Pulvirenti, Ann. Appl. Prob. **29** (2019) 2594-2612

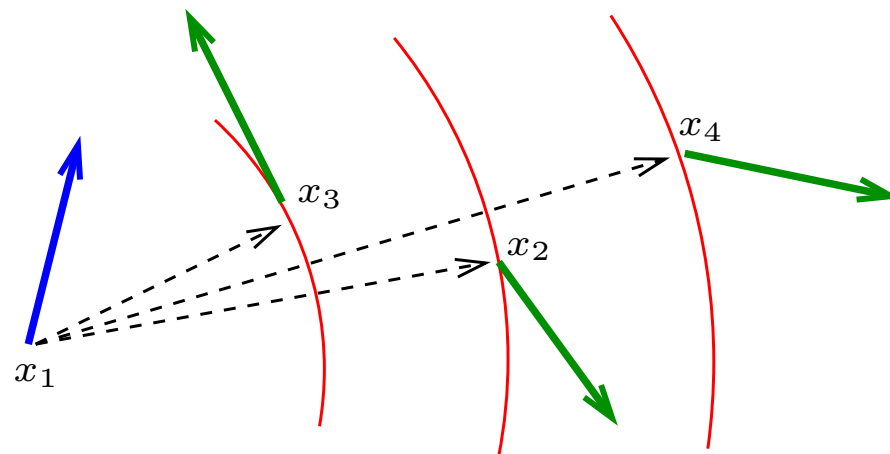
N particles $(x_i(t), v_i(t))_{i=1,\dots,N}$, $(x_i, v_i) \in \mathbb{R}^{2n}$

At given time t , Rank $R^N(i, j)$ of j w.r.t. i :

Sort $(|x_k - x_i|)_{k=1,\dots,n, k \neq i}$ by increasing order

$R^N(i, j)$ is rank of $|x_j - x_i|$ in the list

$R^N(i, j) \in \{1, \dots, N - 1\}$



$$R(1, 3) = 1 \quad R(1, 2) = 2 \quad R(1, 4) = 3$$

Normalized rank: $r^N(i, j) = \frac{R^N(i, j)}{N-1} \in \bigcup_{k=1}^N \left\{ \frac{k}{N-1} \right\}$

Given: $K : [0, 1] \rightarrow [0, \infty)$ s.t. $\int_0^1 K(r) dr = 1$

define $K^N(r) = \frac{K(r)}{\sum_{k=1}^{N-1} K(\frac{k}{N-1})}$

Interaction probabilities:

π_{ij}^N probability of i interacting with j : $\pi_{ij}^N = K^N(r^N(i, j))$

Note: $\sum_{\substack{j=1 \\ j \neq i}}^N \pi_{ij}^N = \sum_{\substack{j=1 \\ j \neq i}}^N K^N(r^N(i, j)) = 1$

Free flights $\dot{x}_i = v_i$, $\dot{v}_i = 0$

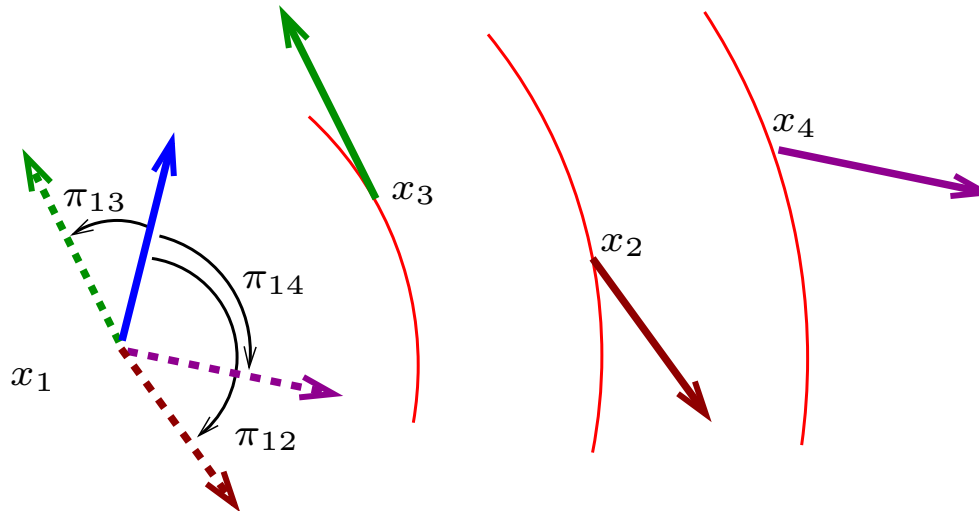
Collisions: at Poisson random times with rate N :

Pick $i \in \{1, \dots, N\}$ with uniform probability $1/N$

Pick $j \in \{1, \dots, N\}$, $j \neq i$ with probability π_{ij}^N

Perform: (x_i, x_j) remains unchanged

(v_i, v_j) changed into (v_j, v_i)



Properties

Rank is a function of positions $r^N(i, j) = r^N(i, j)(x_1, \dots, x_N)$
and so is π_{ij}^N : $\pi_{ij}^N = K^N(r^N(i, j)(x_1, \dots, x_N))$

Rank is permutation-invariant

$$r^N(\sigma(i) \sigma(j))(x_{\sigma(1)}, \dots, x_{\sigma(N)}) = r^N(i, j)(x_1, \dots, x_N)$$

Goal

Derive a Boltzmann operator
under Propagation of Chaos assumption

Note: BBGKY does not work in general
because all particles interact with all particles
but works if K is real analytic [P.D., Pulvirenti]

Previous work using similar ideas:

[D. & Ringhofer, SIAM Appl. Math. 2007] (min interaction)

Notations: $\vec{x} = (x_1, \dots, x_N)$, $\vec{v} = (v_1, \dots, v_N)$
 $Z_j = (x_j, v_j)$, $\vec{Z} = (Z_1, \dots, Z_N)$, $d\vec{Z} = dZ_1 \dots dZ_N$

$f^N(\vec{Z}, t)$: N -particle distribution function

Find eq. for f^N (master equation)

Follow strategy of [Carlen, D., Wennberg, M3AS 2013]

Take $\Phi^N(\vec{Z})$ a test function

Drop drift term for simplicity

$$\frac{d}{dt} \int f^N \Phi^N d\vec{Z} =$$

$$N \int \left[\frac{1}{N} \sum_{i \neq j} \pi_{ij}^N(\vec{x}) \Phi^N(Z_1, \dots, x_i, v_j, \dots, x_j, v_j, \dots, Z_N) - \Phi^N(\vec{Z}) \right] f^N(\vec{Z}, t) d\vec{Z}$$

If $f^N|_{t=0}$ is permutation invariant

then $f^N(t)$ is permutation invariant

Marginal: $f_N^k(Z_1, \dots, Z_k, t) = \int f^N(\vec{Z}, t) dZ_{k+1} \dots dZ_N$

To compute eq. for f_N^1 , use $\Phi^N(\vec{Z}) = \phi(Z_1)$

Propagation of chaos

Assume $f^N(\vec{Z}, t) = \prod_{\ell=1}^N f_N^1(Z_\ell, t) + \text{negligible terms as } N \rightarrow \infty$

Define $\rho_N^1(x, t) = \int f_N^1(x, v, t) dv$

$$\begin{aligned} \frac{d}{dt} \int f_1^N(Z_1) \phi(Z_1) dZ_1 &= \frac{1}{S^N(K)} \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f_1^N(Z_1) f_1^N(Z_2) \\ &\quad \times \left(\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell \right) dZ_1 dZ_2 \end{aligned}$$

where

$$S^N(K) = \frac{1}{N-1} \sum_{k=1}^{N-1} K\left(\frac{k}{N-1}\right) \approx \int_0^1 K(r) dr = 1$$

Need to estimate the behavior as $N \rightarrow \infty$ of:

$$\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell$$

$$\int K(r^N(1, 2)(\vec{x})) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell = \mathbb{E} \left\{ K(r^N(1, 2)(x_1, x_2, x_3, \dots, x_N)) \right\}$$

when $x_3, \dots, x_N$ are iid random variables following the law ρ_N^1

$$= \sum_{j=0}^{N-2} K\left(\frac{j+1}{N-1}\right) \binom{N-2}{j} \text{Prob}\left(\text{Card}\{k \in \{3, \dots, N\} \text{ s.t. } |x_k - x_1| \leq |x_2 - x_1|\} = j\right)$$

$$= \sum_{j=0}^{N-2} K\left(\frac{j+1}{N-1}\right) \binom{N-2}{j} m^j (1-m)^{(N-2)-j}$$

$$\text{where } m = M_{\rho_N^1}(x_1, |x_2 - x_1|) = \int_{|x - x_1| \leq |x_2 - x_1|} \rho_N^1(x) dx$$

$\approx$ Bernstein polynomial approximation of $K(m)$

$$= K(m) + \mathcal{O}\left(\frac{1}{N}\right)$$

In the limit $N \rightarrow \infty$, if $f_N^1(Z_1) \rightarrow f(Z_1)$, then

$$\begin{aligned} & \frac{d}{dt} \int f(Z_1, t) \phi(Z_1) dZ_1 = \\ &= \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(Z_1) f(Z_2) K(M_\rho(x_1, |x_2 - x_1|)) dZ_1 dZ_2 \end{aligned}$$

In strong form (adding back the drift term):

$$\partial_t f + v \cdot \nabla_x f = Q(f)$$

$$Q(f)(x, v) = \rho(x) \int f(x', v) K(M_\rho(x, |x' - x|)) dx' - f(x, v)$$

$$\text{with } M_\rho(x, s) = \int_{|x' - x| \leq s} \rho(x') dx'$$

$$\begin{aligned}\int Q(f) dv &= \rho(x) \left(\int \rho(x') K(M_\rho(x, |x' - x|)) dx' - 1 \right) \\ &= \rho(x) \left(\int K(m) dm - 1 \right) = 0\end{aligned}$$

So, the continuity eq. is satisfied

$$\partial_t \rho + \nabla_x \cdot j = 0, \quad j(x, t) = \int_{\mathbb{R}^n} f(x, v, t) v dv$$

3. Nearest neighbor

A. Blanchet & P. D., J. Stat. Phys. **169** (2019), 929-950

At each collision event:

A particle is selected randomly with uniform probability

When selected, particle i follows its nearest neighbor

i.e. the probability π_{ij}^N of i interacting with j is

$$\pi_{ij}^N = \delta(R^N(i, j) - 1)$$

here: rank is un-normalized: $R^N(i, j) \in \{1, \dots, N - 1\}$, so:

$\pi_{ij}^N = 0$ except if $j =$ nearest neighbor in which case $\pi_{ij}^N = 1$

Rate of collision events: $\lambda(N) N$ with $\lambda(N) \rightarrow \infty$ TBD

Master eq. unchanged: - multiply by $\lambda(N)$

- replace $K^N(r^N(i, j))$ by $\delta(R^N(i, j) - 1)$

$$\begin{aligned} \frac{d}{dt} \int f_1^N(Z_1) \phi(Z_1) dZ_1 &= \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f_1^N(Z_1) f_1^N(Z_2) \\ &\times \left(\lambda(N) (N-1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell \right) dZ_1 dZ_2 \end{aligned}$$

Need to estimate the behavior as $N \rightarrow \infty$ of:

$$\lambda(N) (N-1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho_N^1(x_\ell) dx_\ell$$

Drop scripts N and 1 for clarity

e.g. f_1^N replaced by f ; ρ_1^N replaced by ρ

$$\int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho(x_\ell) dx_\ell = \mathbb{E} \left\{ \delta(R^N(1, 2)(\vec{x}) - 1) \right\}$$

when $x_3, \dots, x_N$ are iid random variables following the law ρ

$$= \text{Prob} \left(\text{Card} \{ k \in \{3, \dots, N\} \text{ s.t. } |x_k - x_1| \leq |x_2 - x_1| \} = 0 \right)$$

$$= (1 - m)^{(N-2)}$$

$$\text{where } m = M_\rho(x_1, |x_2 - x_1|) = \int_{|x - x_1| \leq |x_2 - x_1|} \rho(x) dx$$

$$\begin{aligned} \text{So: } \lambda(N) (N - 1) \int \delta(R^N(1, 2)(\vec{x}) - 1) \prod_{\ell=3}^N \rho(x_\ell) dx_\ell &= \\ &= \lambda(N) (N - 1) (1 - M_\rho(x_1, |x_2 - x_1|))^{(N-2)} \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} \int f(Z_1) \phi(Z_1) dZ_1 &= \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) f(x_2, v_2) \\ &\quad \times \lambda(N) (N-1) (1 - M_\rho(x_1, |x_2 - x_1|))^{(N-2)} dx_1 dx_2 dv_1 dv_2 \end{aligned}$$

Use polar coordinates $x_2 = x_1 + r\omega$, $r \in [0, \infty)$, $\omega \in \mathbb{S}^{n-1}$

$$\begin{aligned} \frac{d}{dt} \int f(Z_1) \phi(Z_1) dZ_1 &= \\ &= \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) F_f^N(x_1, v_2) dx_1 dv_1 dv_2 \end{aligned}$$

$$F_f^N(x_1, v_2) = \lambda(N) (N-1) \int f(x_1 + r\omega, v_2) (1 - M_\rho(x_1, r))^{(N-2)} r^{n-1} dr d\omega$$

Change of variables $m = M_\rho(x_1, r) \Leftrightarrow r = R_\rho(x_1, m)$, $m \in [0, 1]$

$$F_f^N(x_1, v_2) = \lambda(N) (N - 1) \int_0^1 G_f(x_1, v_2, m) (1 - m)^{N-2} dm$$

$$\text{with } G_f(x_1, v_2, m) = \frac{\int f(x_1 + R_\rho(x_1, m)\omega, v_2) d\omega}{\int \rho(x_1 + R_\rho(x_1, m)\omega) d\omega}$$

As $m \rightarrow 0$

$$G_f(x_1, v_2, m) = \frac{f(x_1, v_2)}{\rho(x_1)} + m^{\frac{2}{n}} \frac{n^{\frac{2}{n}-1}}{2} \frac{1}{\rho^{\frac{2}{n}+1}(x_1)} D(\rho, f)(x_1, v_2) + o(m^{\frac{2}{n}})$$

with

$$D(\rho, f)(x_1, v_2) = \Delta_x f(x_1, v_2) - \frac{f(x_1, v_2)}{\rho(x_1)} \Delta_x \rho(x_1)$$

Take:

$$\frac{1}{\lambda(N)} = (N-1) \frac{n^{\frac{2}{n}-1}}{2} \int_0^1 m^{\frac{2}{n}} (1-m)^{N-2} dm$$

$\lambda(N) \rightarrow \infty$ as $N \rightarrow \infty$. For instance: $\lambda(N) = 2N$ if $n = 2$

larger frequency than in smooth case

Then, as $N \rightarrow \infty$:

$$F_f^N(x_1, v_2) - \lambda(N) \frac{f(x_1, v_2)}{\rho(x_1)} \longrightarrow \frac{1}{\rho^{\frac{2}{n}+1}(x_1)} D(\rho, f)(x_1, v_2)$$

Concentration of measure at $m = 0$

In the limit $N \rightarrow \infty$, if $f_N^1(Z_1) \rightarrow f(Z_1)$, then

$$\begin{aligned} & \frac{d}{dt} \int f(x_1, v_1, t') \phi(x_1, v_1) dx_1 dv_1 = \\ &= \int (\phi(x_1, v_2) - \phi(x_1, v_1)) f(x_1, v_1) \frac{1}{\rho^{\frac{2}{n}+1}(x_1)} D(\rho, f)(x_1, v_2) dx_1 dx_2 dv_2 \\ &= \int \phi(x_1, v_1) \frac{1}{\rho^{\frac{2}{n}}(x_1)} \left[\Delta_x f(x_1, v_1) - \frac{f(x_1, v_1)}{\rho(x_1)} \Delta_x \rho(x_1) \right] dx_1 dv_1 \end{aligned}$$

In strong form (adding back the drift term):

$$\begin{aligned} \partial_t f + v \cdot \nabla_x f &= Q(f) \\ Q(f)(x, v) &= \frac{1}{\rho(x)^{\frac{2}{n}}} \left(\Delta_x f(x, v) - \frac{f(x, v)}{\rho(x)} \Delta_x \rho(x) \right) \end{aligned}$$

$$\int Q(f) dv = 0 \quad (\text{obvious})$$

So, the continuity eq. is satisfied

$$\partial_t \rho + \nabla_x \cdot j = 0, \quad j(x, t) = \int_{\mathbb{R}^n} f(x, v, t) v dv$$

4. Conclusion

Smooth rank-based topological interaction

- Derivation of spatially non-local Boltzmann model

- Highly nonlinear new collision operator

- Proof of propagation of chaos and convergence

Nearest-neighbor interaction

- gives rise to spatial nonlinear diffusion

- Still preserves continuity eq.

Perspectives

- Well-posedness of the resulting eqs.

- Proof of propagation of chaos in NN interaction

- Hydrodynamic limits

- Extension to more complex interaction rules