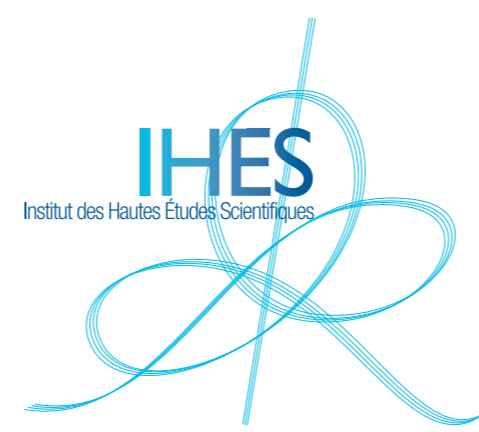




Discovering the torsion bigravity world

Vasilisa Nikiforova

Institut des Hautes Études Scientifiques



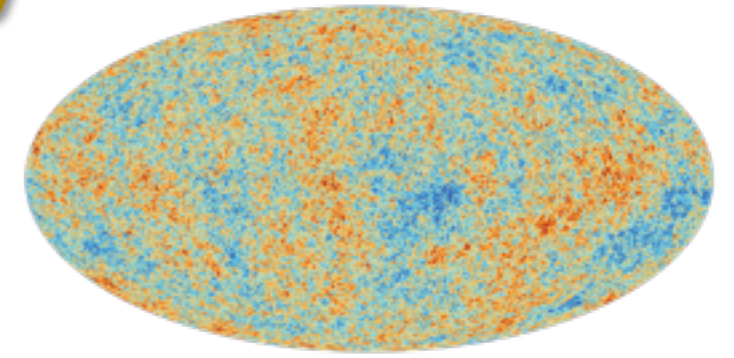
**Tours
2022**



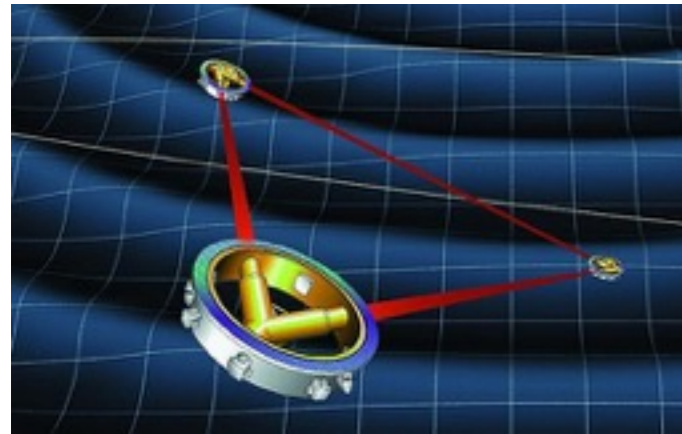
Golden Era of gravity



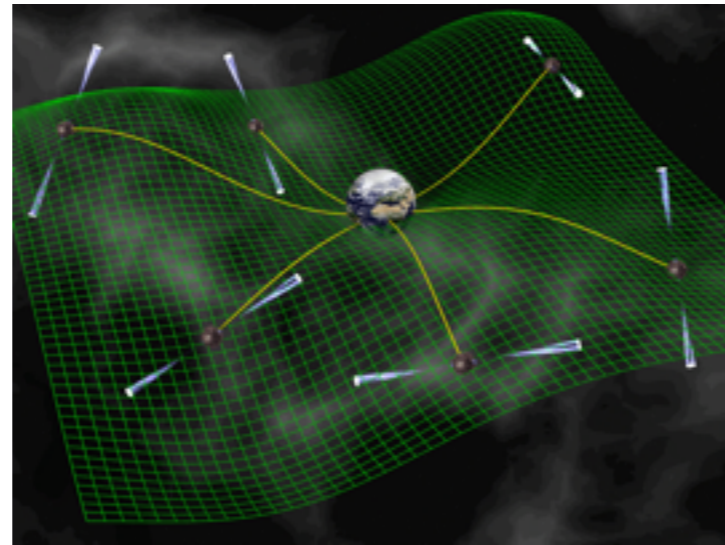
LIGO-Virgo



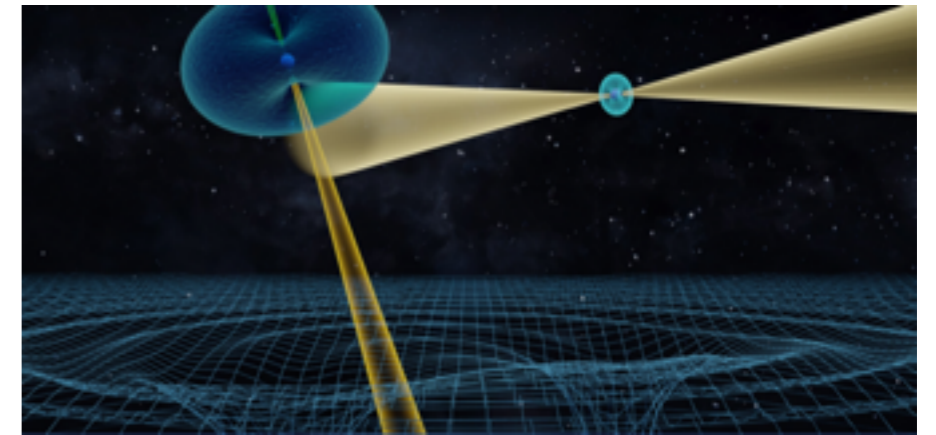
CMB



LISA



pulsar timing array



binary pulsars



Microscope

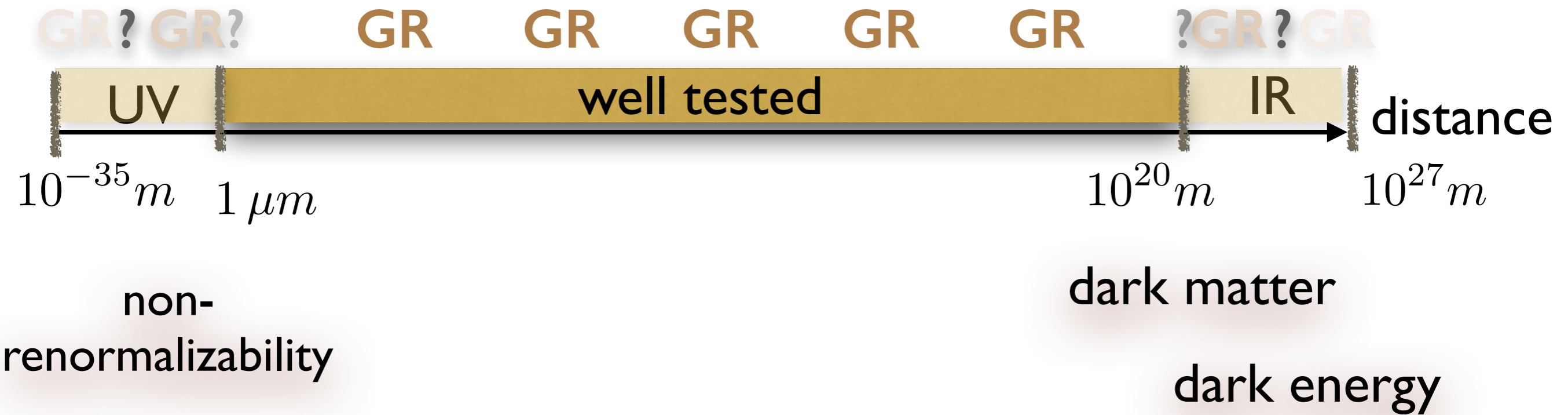


SgrA*

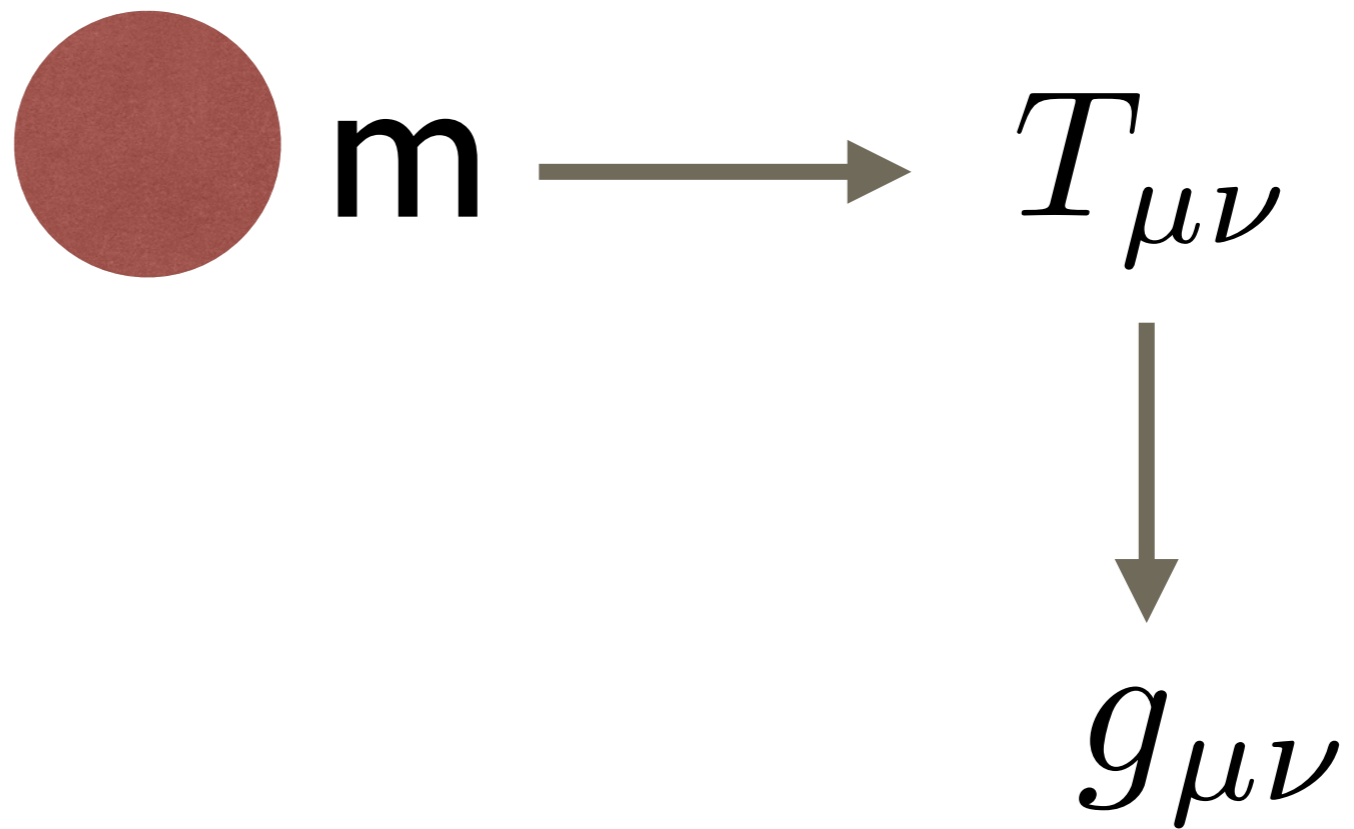


EHT

Puzzles of gravity

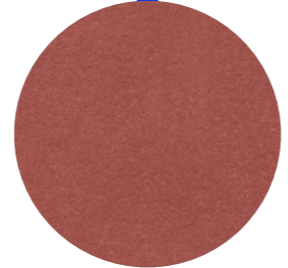


Sources of gravity



Sources of gravity

quantum spin density $S_{\lambda\mu\nu} \equiv \bar{\psi}\gamma^{[\lambda\mu\nu]}\psi \longrightarrow ? T^{\lambda}_{\mu\nu}$



m



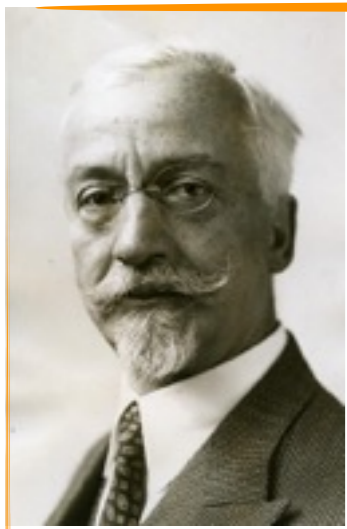
$T_{\mu\nu}$



$g_{\mu\nu}$

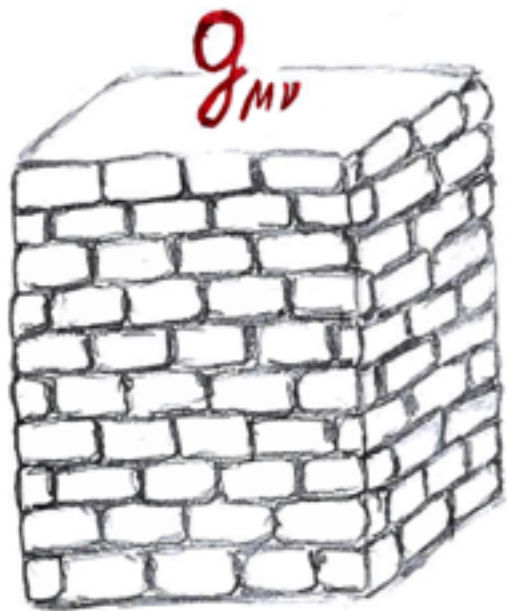
mass and spin are characteristics of a particle linked to space-time symmetries

Quantum spin as a source of gravity ?

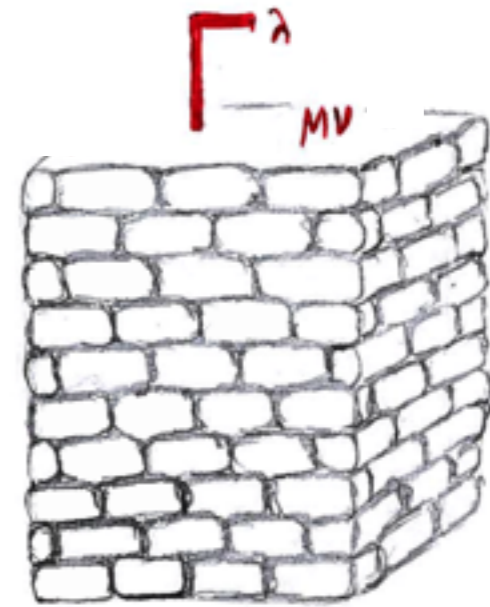


Metric and (torsionfull) connection

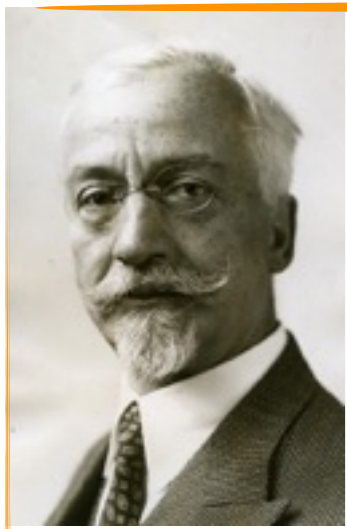
Elie Cartan



metric

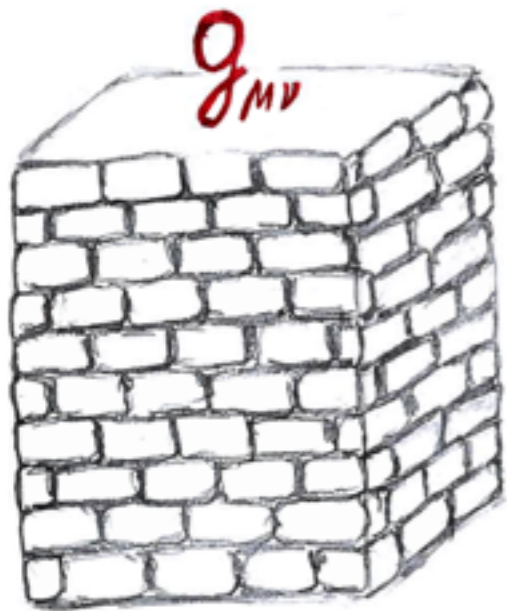


connection

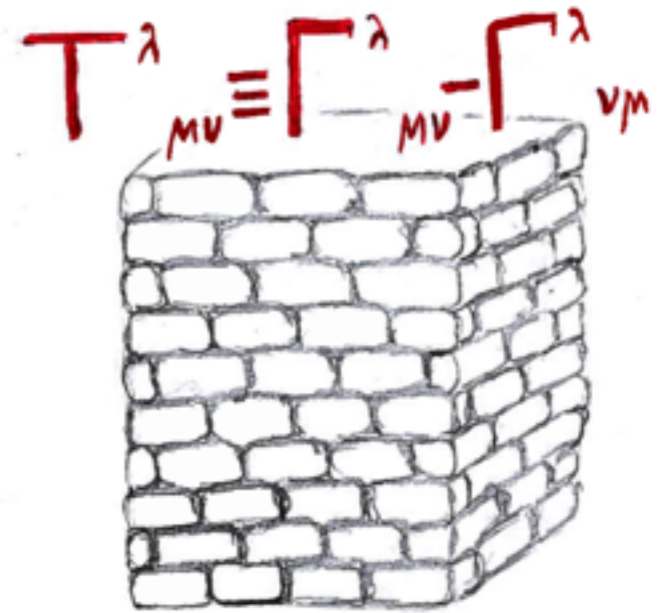


Metric and (torsionfull) connection

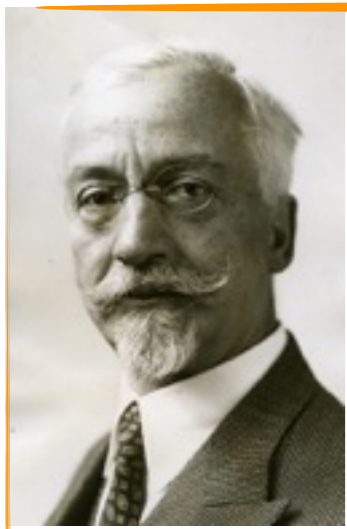
Elie Cartan



metric

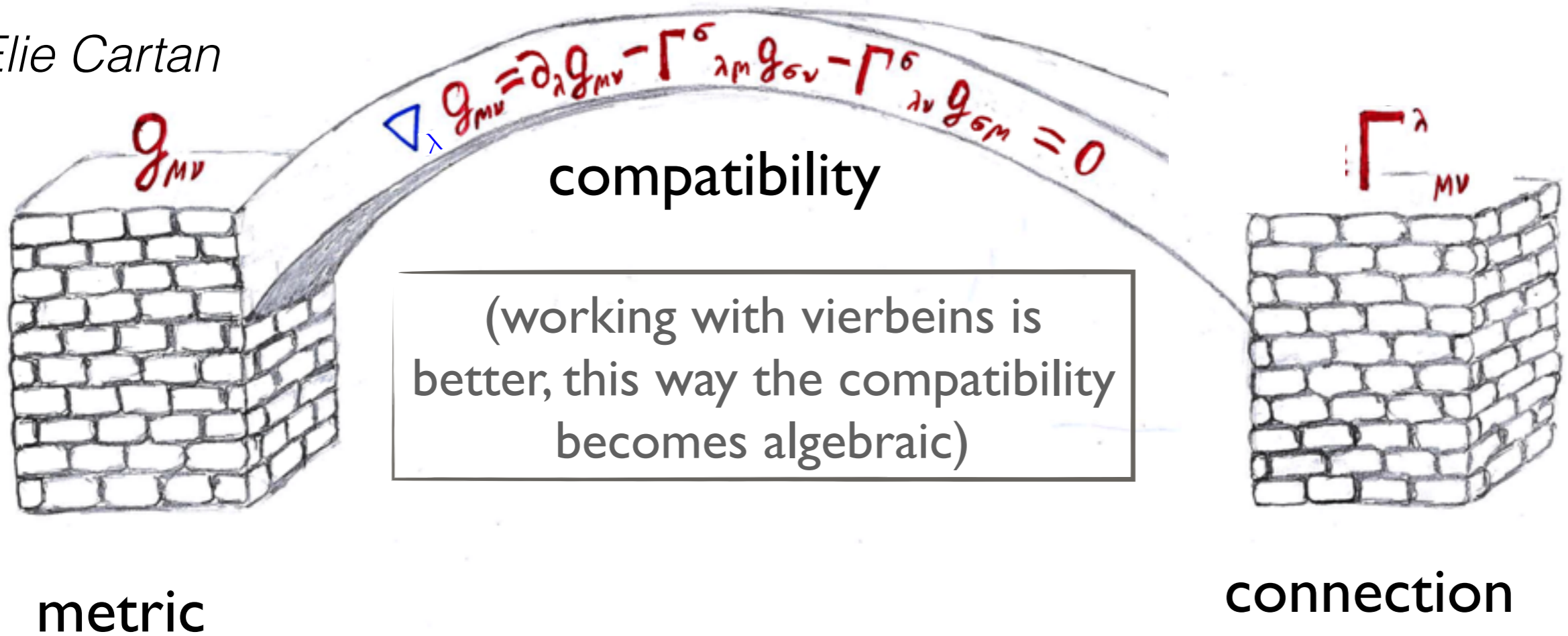


connection



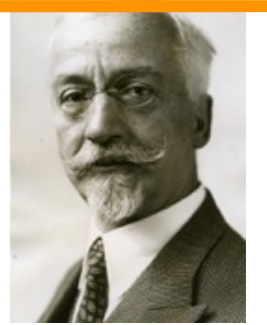
Metric and (torsionfull) connection

Elie Cartan

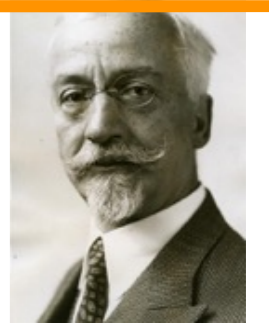




Einstein-Cartan theory

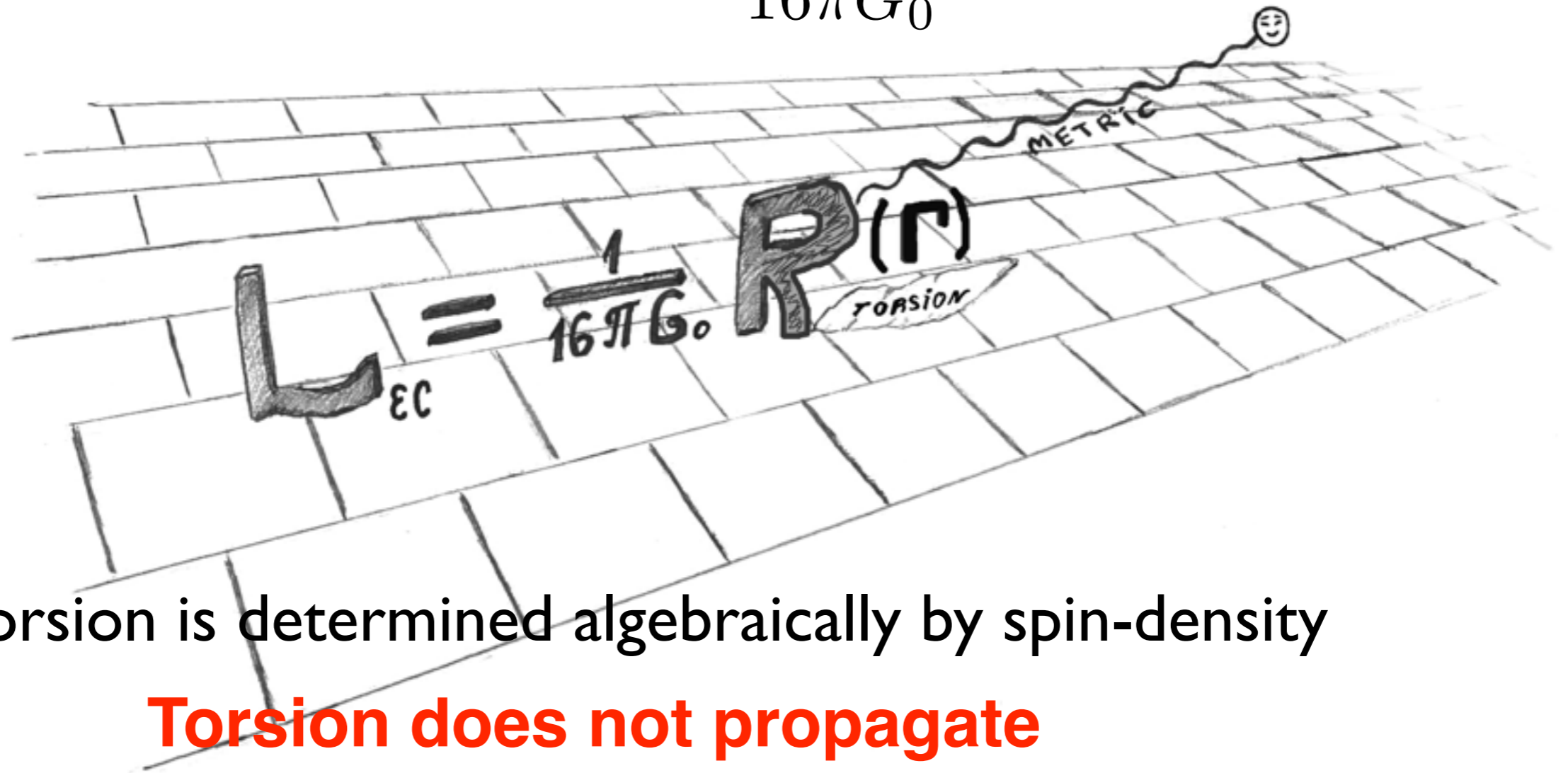


$$L_{\text{Einstein-Cartan}} = \frac{1}{16\pi G_0} \mathcal{R}(\Gamma)$$



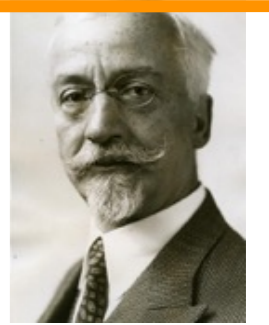
Einstein-Cartan theory

$$L_{\text{Einstein-Cartan}} = \frac{1}{16\pi G_0} \mathcal{R}(\Gamma)$$



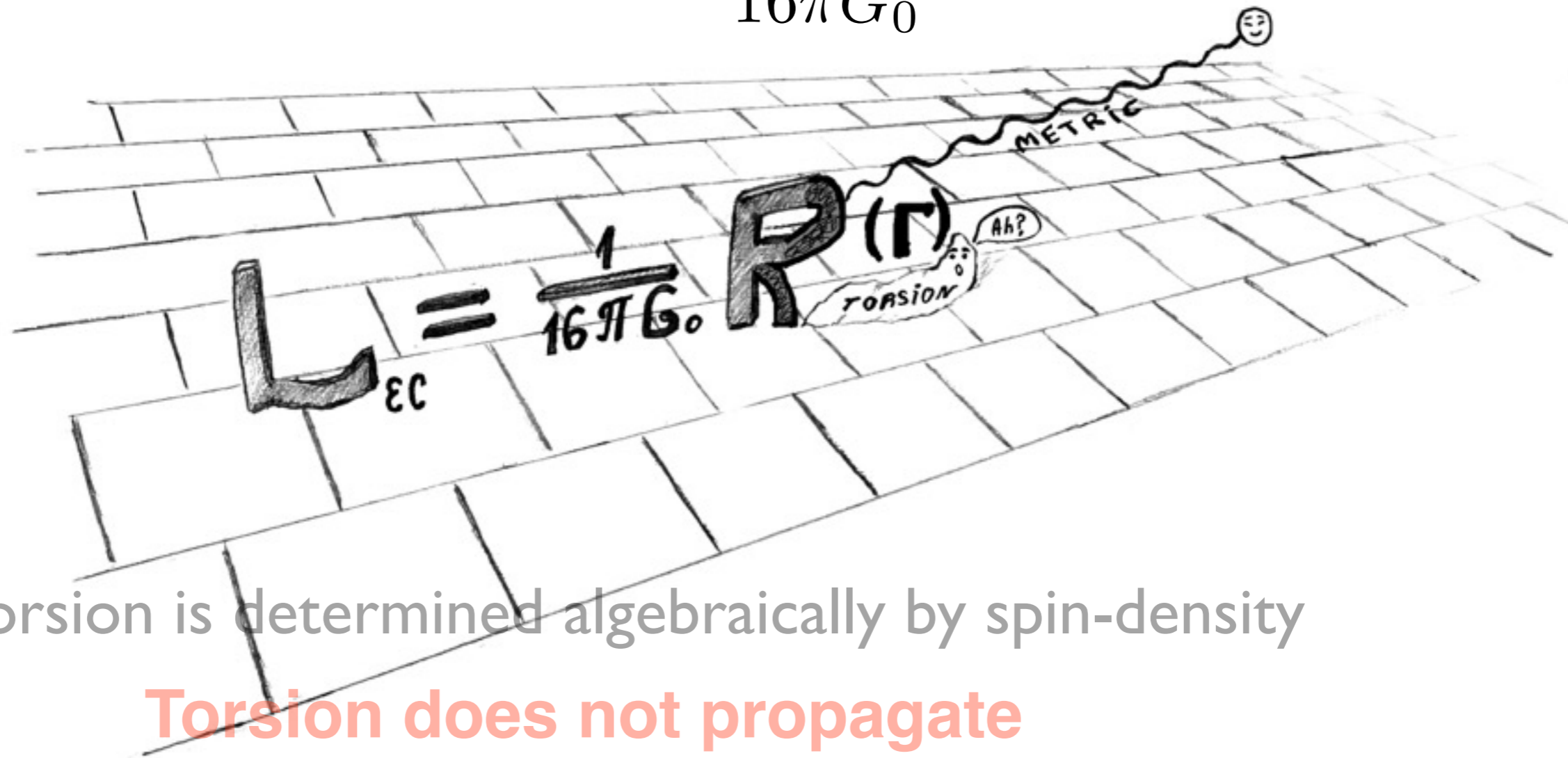
torsion is determined algebraically by spin-density

Torsion does not propagate



Einstein-Cartan theory

$$L_{\text{Einstein-Cartan}} = \frac{1}{16\pi G_0} \mathcal{R}(\Gamma)$$



Torsion does not propagate

Could there be a theory of gravity containing dynamical torsion ?

Sezgin-van Nieuwenhuizen'80,
Hayashi-Shirafuji'81

Torsion bigravity

Sezgin-van Nieuwenhuizen'80,
Hayashi-Shirafuji'81

Nair
Randjbar-Daemi
Rubakov'09

Nikiforova
Randjbar-Daemi
Rubakov'09

Damour
Nikiforova'19



Torsion bigravity

Sezgin-van Nieuwenhuizen'80,
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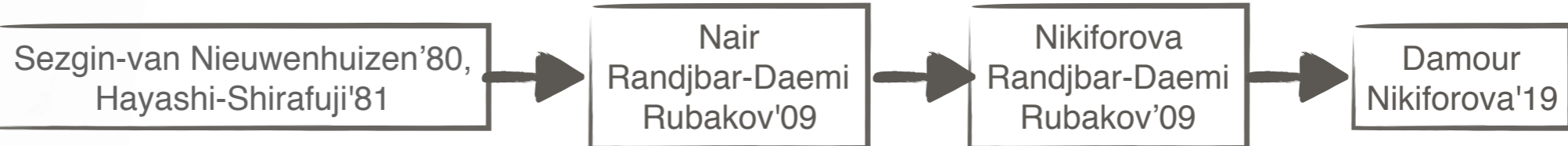
Damour
Nikiforova'19

L_{GR} +

$$L = \frac{1}{16\pi G_0} R(g)$$



Torsion bigravity

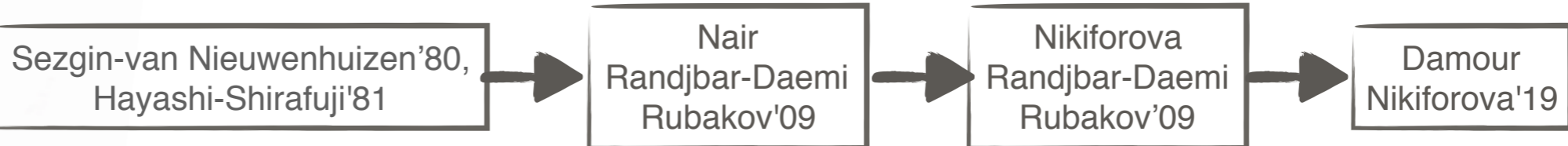


$$L_{\text{GR}} + L_{\text{Einstein-Cartan}}$$

$$L = \frac{1}{16\pi G_0} R(g) + \frac{1}{16\pi G_0} \mathcal{R}(\Gamma)$$



Torsion bigravity

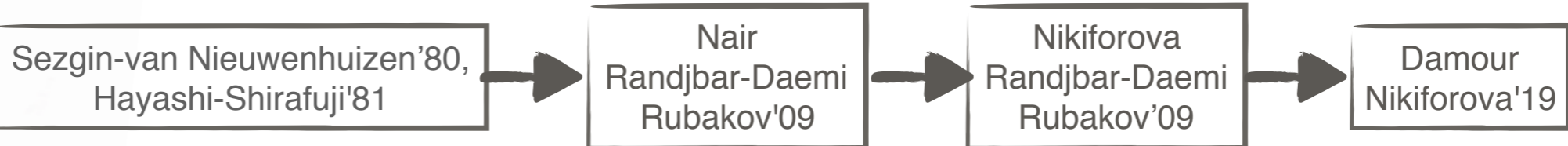


L_{GR} + $L_{\text{Einstein-Cartan}}$ + quadratic terms

$$L = \frac{1}{16\pi G_0} R(g) + \frac{1}{16\pi G_0} \mathcal{R}(\Gamma) + \frac{1}{16\pi G_0} \left(\mathcal{R}_{(\mu\nu)}(\Gamma) \mathcal{R}^{(\mu\nu)}(\Gamma) - \frac{1}{3} \mathcal{R}(\Gamma)^2 \right) + \mathcal{R}_{[\mu\nu]}(\Gamma) \mathcal{R}^{[\mu\nu]}(\Gamma)$$



Torsion bigravity



L_{GR} + $L_{\text{Einstein-Cartan}}$ + quadratic terms

$$L = \frac{1}{16\pi G_0(1+\eta)} R(g) + \frac{\eta}{16\pi G_0(1+\eta)} \mathcal{R}(\Gamma) + \frac{\eta}{16\pi G_0 \kappa^2} \left(\mathcal{R}_{(\mu\nu)} \mathcal{R}^{(\mu\nu)} - \frac{1}{3} \mathcal{R}^2 \right) + c_{34} \mathcal{R}_{[\mu\nu]}(\Gamma) \mathcal{R}^{[\mu\nu]}(\Gamma)$$

Explicit field equations

$$\frac{\eta}{(1+\eta)16\pi G_0} \left(\mathcal{R}_{ij}(\Gamma) - \frac{1}{2}\eta_{ij}\mathcal{R}(\Gamma) \right) + \frac{1}{(1+\eta)16\pi G_0} \left(R_{ij}(g) - \frac{1}{2}\eta_{ij}R(g) \right) \\ + \frac{\eta}{16\pi G_0\kappa^2} \left[\mathcal{R}_{ki}(\Gamma)\mathcal{R}_{kj}(\Gamma) + \mathcal{R}_{kl}(\Gamma)\mathcal{R}_{kilj}(\Gamma) - \frac{2}{3}\mathcal{R}(\Gamma)\mathcal{R}_{ij}(\Gamma) - \frac{1}{2}\eta_{ij} \left(\mathcal{R}_{kl}(\Gamma)\mathcal{R}_{kl}(\Gamma) - \frac{1}{3}\mathcal{R}(\Gamma)^2 \right) \right] = T_{ij}$$

2nd order in connection and metric !

stress-energy tensor

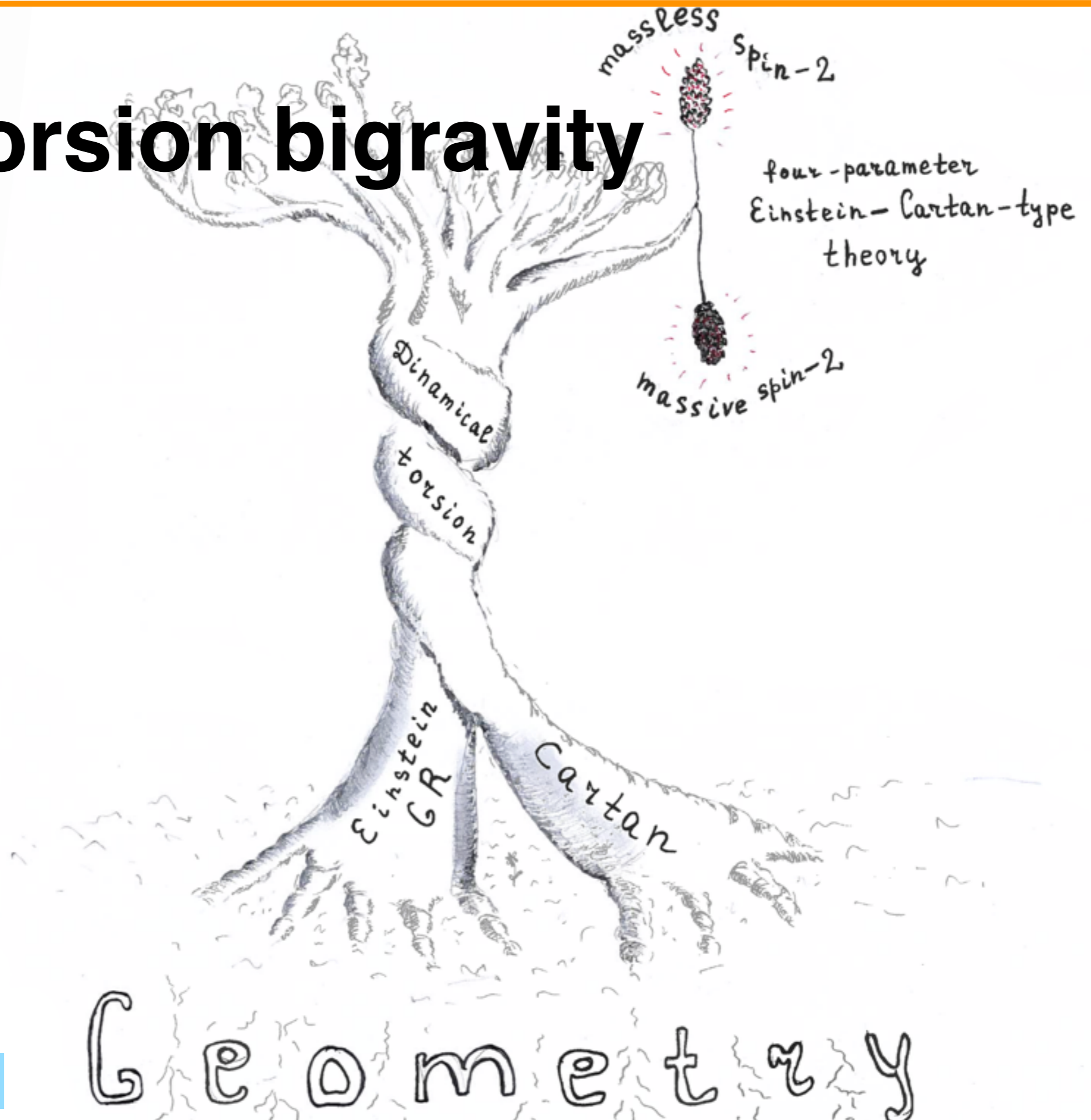
$$\left[\eta_{ik} \left(D_m P_{jm} - \frac{2}{3}D_j P \right) - D_i P_{jk} \right] - \left[\eta_{jk} \left(D_m P_{im} - \frac{2}{3}D_i P \right) - D_j P_{ik} \right] \\ + \frac{\eta}{(1+\eta)16\pi G_0} (K_{ikj} - K_{jki} - K_{ill}\eta_{jk} + K_{jll}\eta_{ik}) \\ + (K_{mkn} - K_{nkm} - K_{mll}\eta_{nk} + K_{nll}\eta_{mk}) \left(\eta_{im}P_{jn} - \eta_{jm}P_{in} - \frac{2}{3}\eta_{im}\eta_{jn}P \right) = S_{ijk}$$

spin density

$$P_{ij} \equiv \frac{\eta}{16\pi G_0\kappa^2} \mathcal{R}_{(ij)}(\Gamma) + c_{34} \mathcal{R}_{[ij]}(\Gamma)$$

$$K_{ijk} = \frac{1}{2} (T_{i[jk]} + T_{j[ki]} - T_{k[ij]})$$

Torsion bigravity



Torsion bigravity

Field content around flat space:

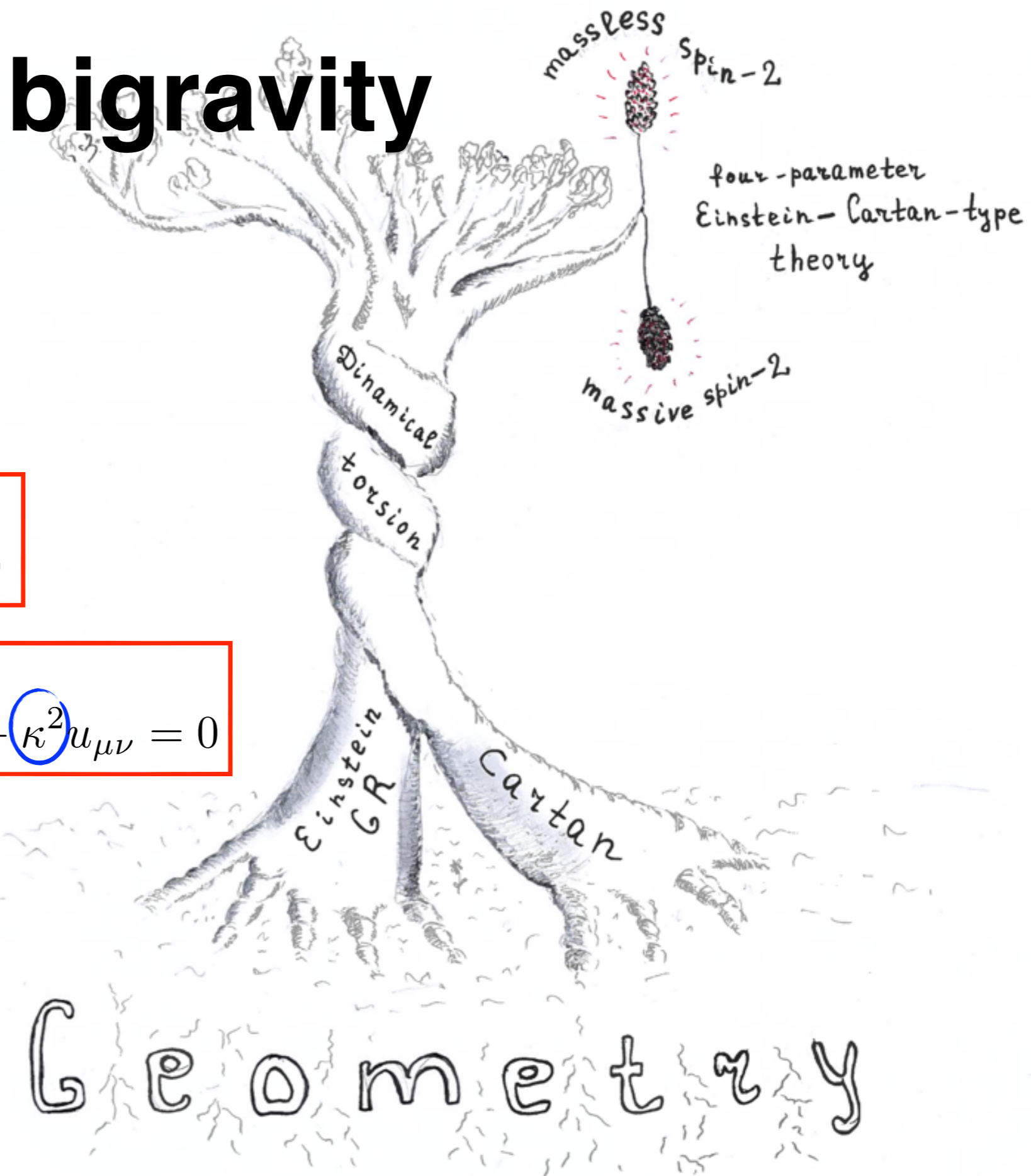
massless spin-2

$$\square \bar{h}_{\mu\nu} + \partial_{\mu\nu} \bar{h} - \partial_{\mu\sigma} \bar{h}_{\nu}^{\sigma} - \partial_{\nu\sigma} \bar{h}_{\mu}^{\sigma} = 0$$

massive spin-2

$$\square u_{\mu\nu} + \partial_{\mu\nu} u - \partial_{\mu\sigma} u_{\nu}^{\sigma} - \partial_{\nu\sigma} u_{\mu}^{\sigma} - \kappa^2 u_{\mu\nu} = 0$$

$$u_{\mu\nu} \equiv \mathcal{R}_{(1)\mu\nu} - \frac{1}{6} g_{\mu\nu} \mathcal{R}_{(1)}$$



Theories with massive spin-2:

Theoretical consistency

bimetric gravity

$$g_{\mu\nu}, f_{\mu\nu}$$

$$L \sim \sqrt{g}R(g) + \sqrt{f}R(f) - V(\sqrt{g^{-1}f})$$

☺ $m = 0, s = 2$ $\alpha \delta g_{\mu\nu} + \beta \delta f_{\mu\nu}$

☺ $m \neq 0, s = 2$ $\bar{\alpha} \delta g_{\mu\nu} + \bar{\beta} \delta f_{\mu\nu}$

ghost-free

(5 dof for the massive spin-2;
generic existence of 5
constraints)

de Rham-Gabadadze-Tolley'10,
Hassan-Rosen'11,
Volkov'14

torsion bigravity

$$g_{\mu\nu}, T^{\lambda}_{\mu\nu}$$

$$L \sim \sqrt{g}R(g) + \sqrt{g}\mathcal{R}(\Gamma) + \sqrt{g} \left(\mathcal{R}_{(\mu\nu)}\mathcal{R}^{(\mu\nu)} - \frac{1}{3}\mathcal{R}^2 \right) + \sqrt{g}\mathcal{R}_{[\mu\nu]}\mathcal{R}^{[\mu\nu]}$$

☺ $m = 0, s = 2$ $\delta g_{\mu\nu} + \frac{1}{\kappa^2}\delta\mathcal{R}_{\mu\nu}$

☺ $m \neq 0, s = 2$ $\delta\mathcal{R}_{\mu\nu} - \frac{1}{6}g_{\mu\nu}\delta\mathcal{R}$

ghost-free

(5 dof for the massive spin-2) around flat space,
Einstein spaces ($R_{\mu\nu} = \Lambda g_{\mu\nu}$),
in nonlinear static spherically symm. solutions

Sezgin-van Nieuwenhuizen'80, Hayashi-Shirafuji'81, Nair-Randjbar-Daemi-Rubakov'09
Nikiforova-Randjbar-Daemi-Rubakov'09, Damour-Nikiforova'19, Nikiforova'20

Theories with massive spin-2:

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torsion bigravity

$$g_{\mu\nu}, T^{\lambda}_{\mu\nu}$$

$$L \sim \sqrt{g}R(g) + \sqrt{g}\mathcal{R}(\Gamma) + \sqrt{g} \left(\mathcal{R}_{(\mu\nu)}\mathcal{R}^{(\mu\nu)} - \frac{1}{3}\mathcal{R}^2 \right) + \sqrt{g}\mathcal{R}_{[\mu\nu]}\mathcal{R}^{[\mu\nu]}$$

☺ $m = 0, s = 2$ $\delta g_{\mu\nu} + \frac{1}{\kappa^2}\delta\mathcal{R}_{\mu\nu}$

☺ $m \neq 0, s = 2$ $\delta\mathcal{R}_{\mu\nu} - \frac{1}{6}g_{\mu\nu}\delta\mathcal{R}$

ghost-free

(5 dof for the massive spin-2; generic existence of 5 constraints)

ghost-free

(5 dof for the massive spin-2) are Einstein spaces ($R_{\mu\nu} = \Lambda g_{\mu\nu}$), in nonlinear static spherically symmetric solutions

future direction



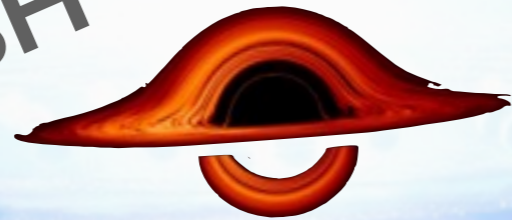
Phenomenology: Torsion bigravity world

OK let us model a world
on a base of torsion bigravity.

And then will see
whether it looks like our real world or not.

Phenomenology: Torsion bigravity world

BH



star



OK let us model a world
on a base of torsion bigravity.

And then will see
whether it looks like our real world or not.

Newtonian
limit

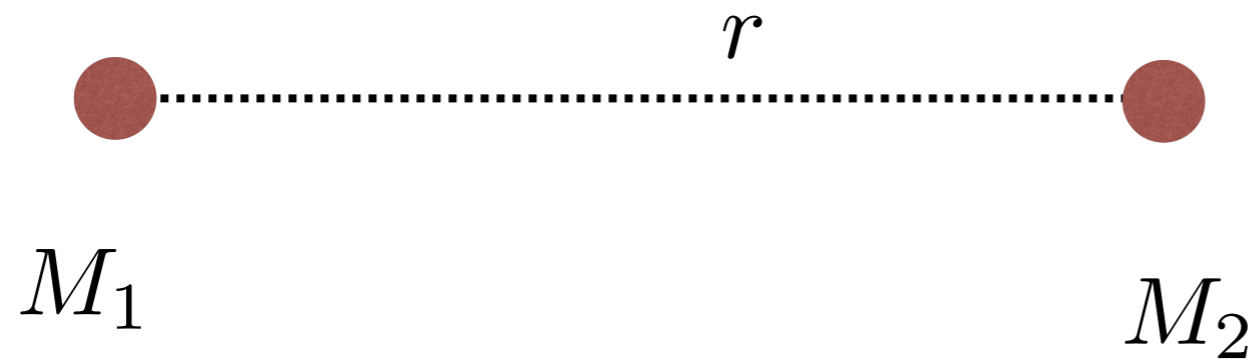


Newtonian limit in torsion bigravity

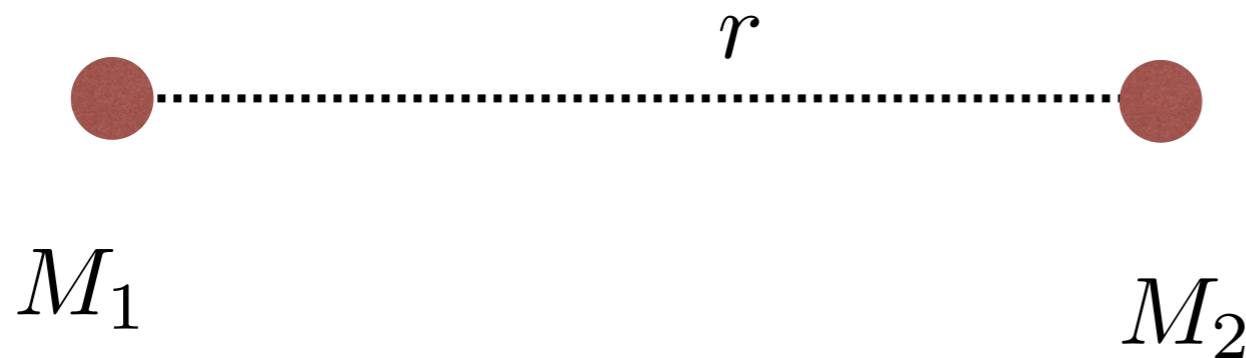
Newtonian
limit



Newtonian limit



Newtonian limit



$$V_{\text{int}}^{\text{Newtonian}} = -G_0 \frac{M_1 M_2}{r} - G_m \frac{M_1 M_2}{r} e^{-\kappa r}$$

massless
spin-2

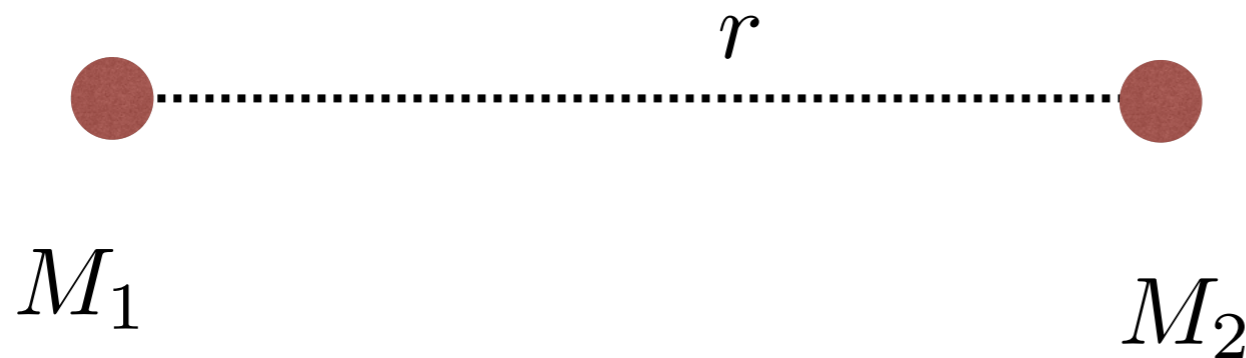
massive
spin-2

Coupling constant :

$$\frac{G_m}{G_0} = \frac{4}{3}\eta$$

$$L_{\text{interaction}} = G_0 T_{00} \left(\frac{-4\pi}{\Delta} \right) T_{00} + G_m T_{00} \left(\frac{-4\pi}{\Delta - \kappa^2} \right) T_{00}$$

Newtonian limit



$$V_{\text{int}}^{\text{Newtonian}} = -G_0 \frac{M_1 M_2}{r} - G_m \frac{M_1 M_2}{r} e^{-\kappa r}$$

$$\frac{G_m}{G_0} = \frac{4}{3}\eta$$

massless
spin-2

massive
spin-2

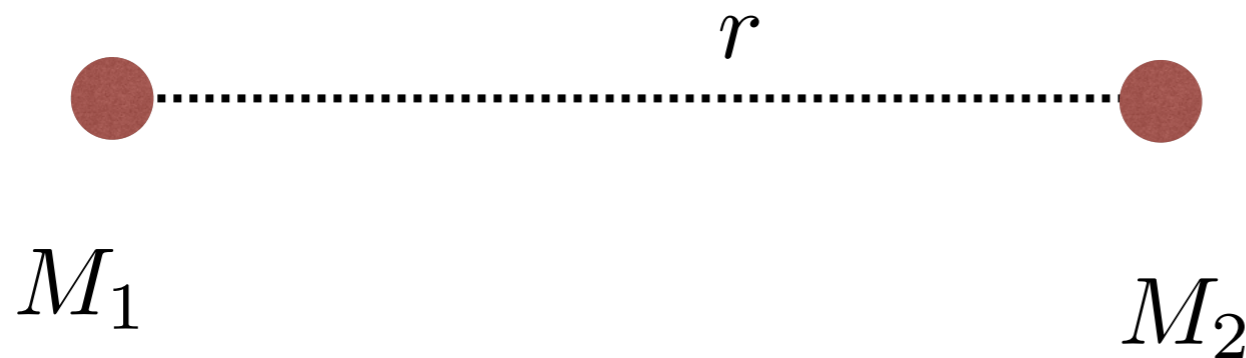
metric



torsion

$$L_{\text{interaction}} = G_0 T_{00} \left(\frac{-4\pi}{\Delta} \right) T_{00} + G_m T_{00} \left(\frac{-4\pi}{\Delta - \kappa^2} \right) T_{00}$$

Newtonian limit



$$V_{\text{int}}^{\text{Newtonian}} = -G_0 \frac{M_1 M_2}{r} - G_m \frac{M_1 M_2}{r} e^{-\kappa r}$$

massless
spin-2

massive
spin-2

$$\frac{G_m}{G_0} = \frac{4}{3}\eta$$

$$\eta \lesssim 3 \times 10^{-4} \quad \text{for } \kappa^{-1} \lesssim 10 \text{ km} \quad (\kappa \gtrsim 10^{-10} \text{ eV})$$



star

Stars in torsion bigravity

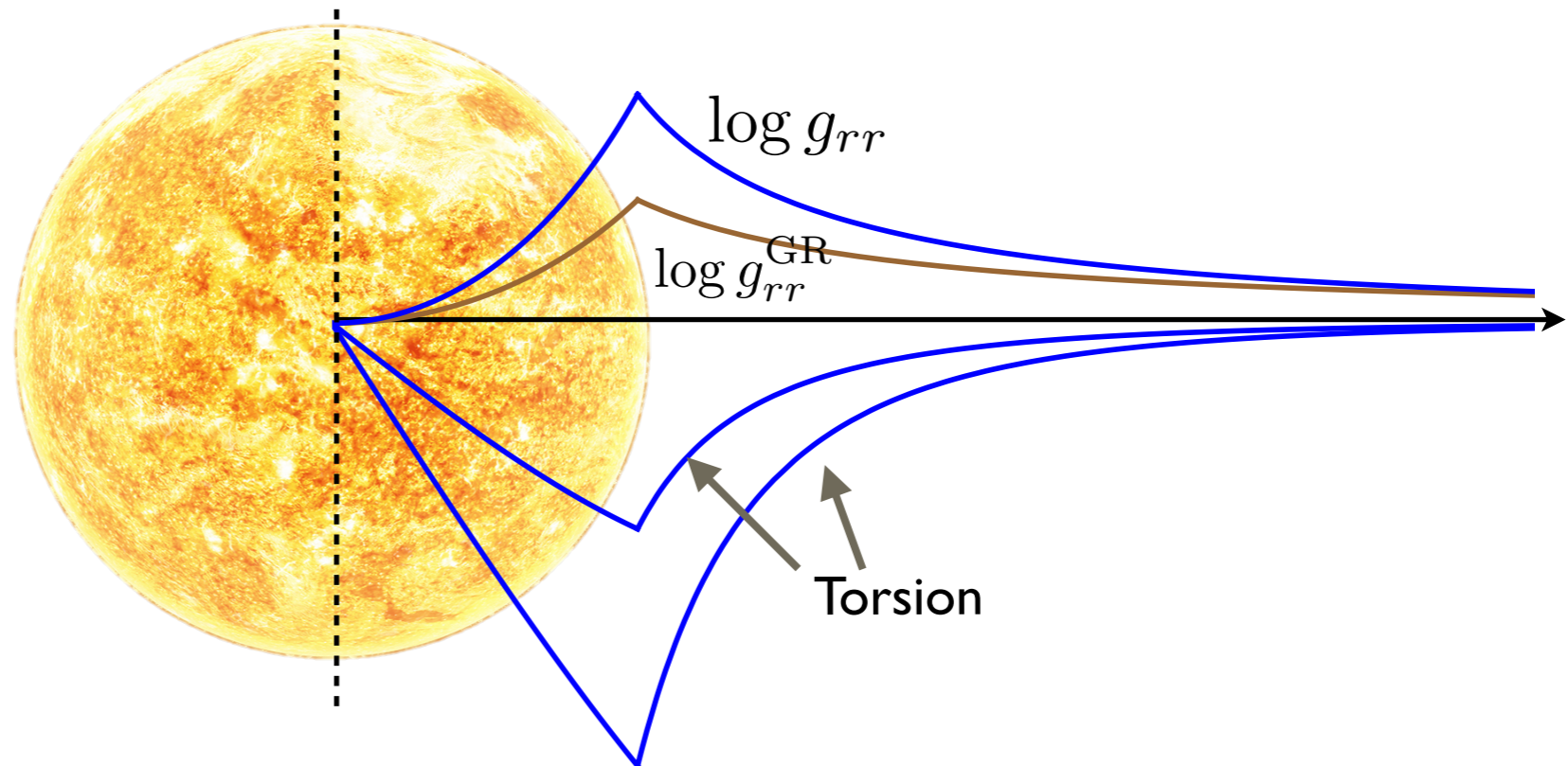
Spherically symmetric star solutions

Damour, Nikiforova'19

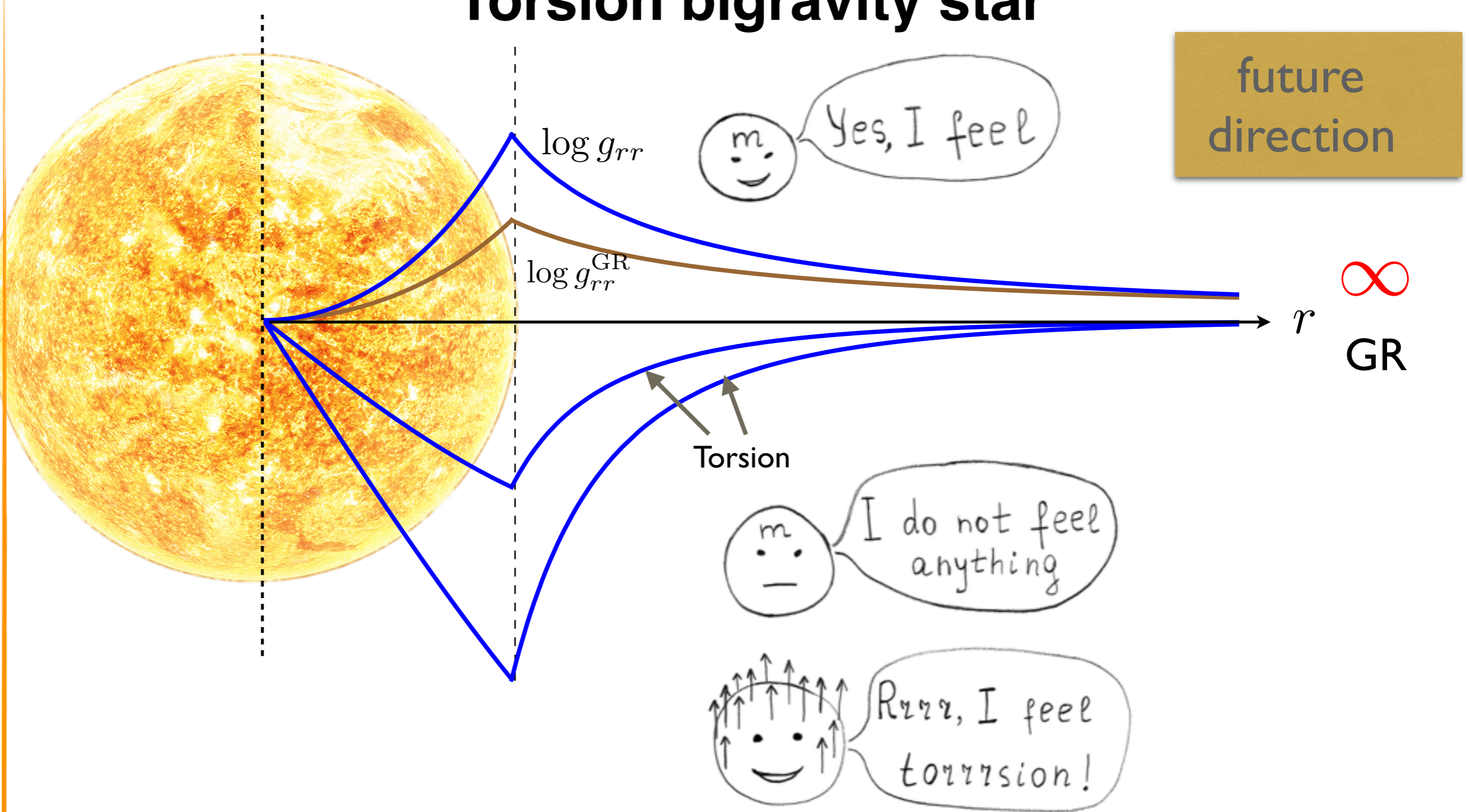
dynamical torsion is generated by $T_{\mu\nu}$ of matter !



metric is different from the one predicted in GR
(compactness is different)



Torsion bigravity star



Direct source of torsion, direct torsion probe

Dirac fermions

$$\mathcal{L}_m = \frac{i\hbar c}{2} (\bar{\psi} \gamma^\mu \nabla_\mu \psi - \nabla_\mu \bar{\psi} \gamma^\mu \psi) - mc^2 \bar{\psi} \psi$$

torsionfull connection



Direct source of torsion, direct torsion probe

Dirac fermions

$$\mathcal{L}_m = \frac{i\hbar c}{2} (\bar{\psi} \gamma^\mu \nabla_\mu \psi - \nabla_\mu \bar{\psi} \gamma^\mu \psi) - mc^2 \bar{\psi} \psi$$

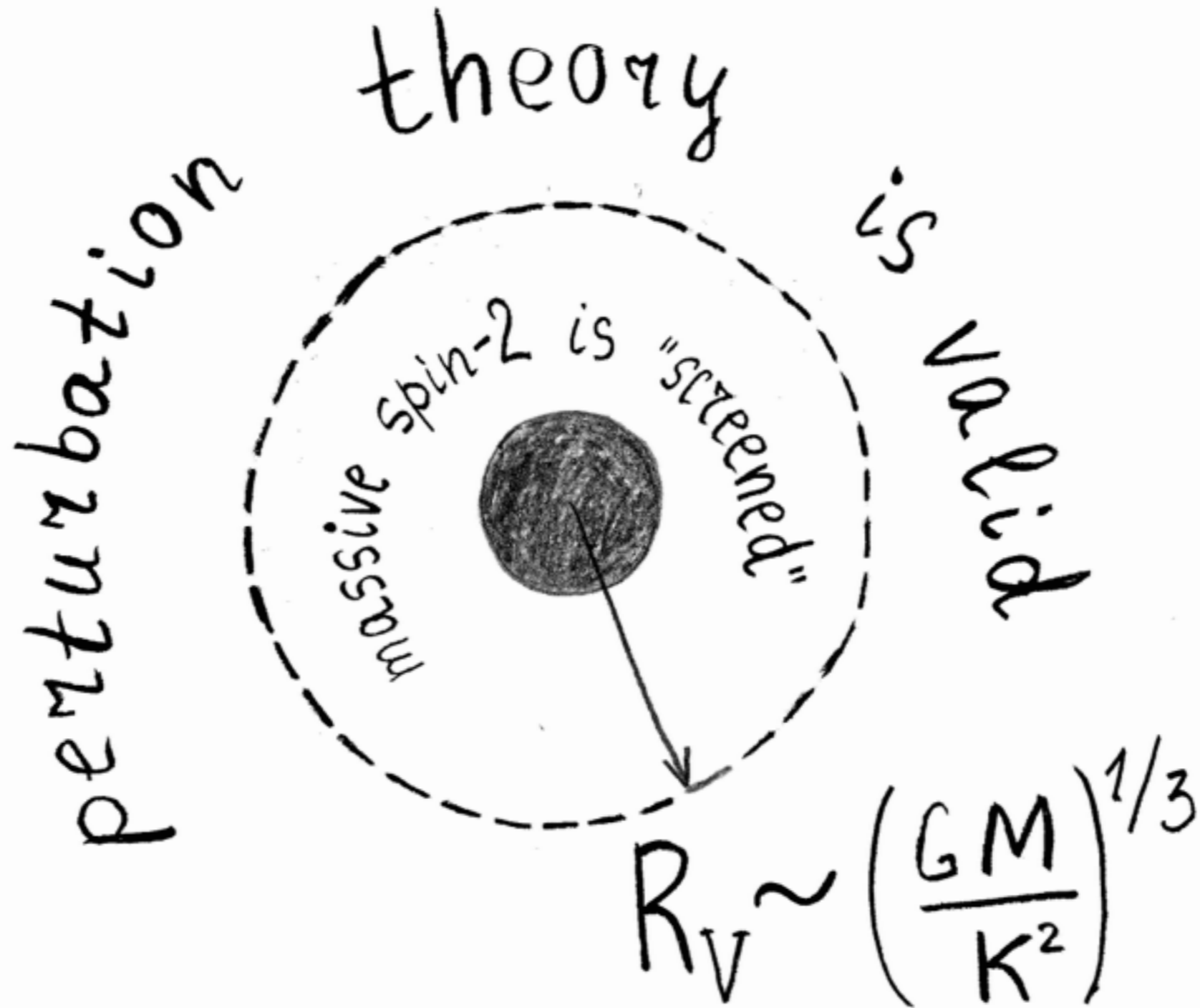
torsionfull connection



$$\bar{\psi} \gamma^{[\lambda\mu\nu]} \psi T_{\lambda\mu\nu}$$

torsion is not easy to detect directly

Vainshtein mechanism



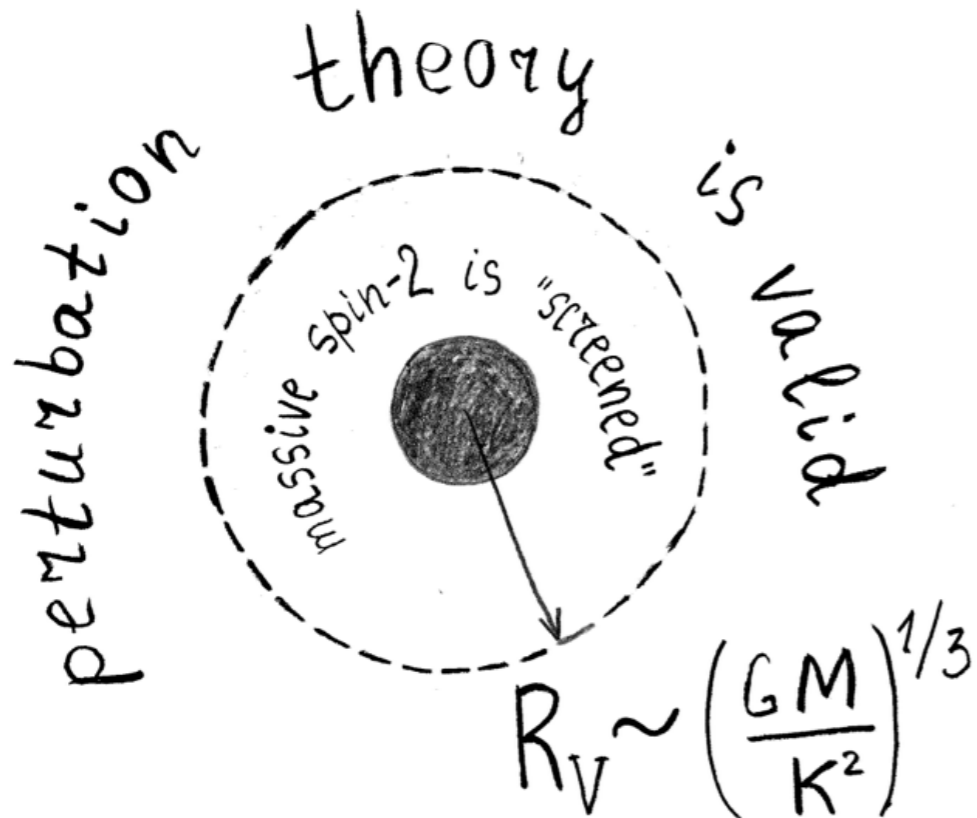
Theories with massive spin-2:

Vainshtein mechanism

bimetric gravity

torsion bigravity

Vainshtein screening

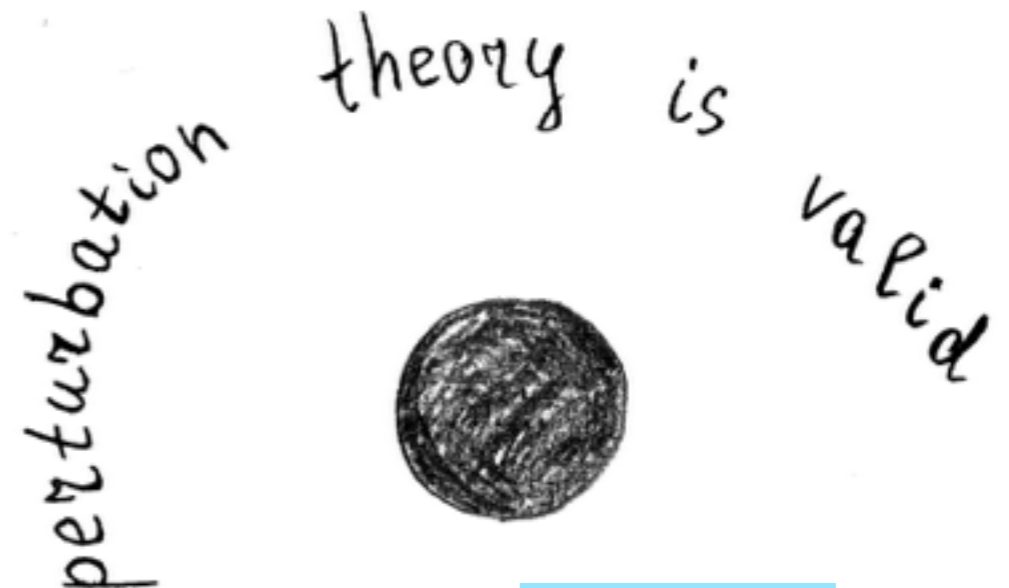


Babichev-Deffayet-Ziour'10,
Babichev-Deffayet'13

no $\frac{1}{\kappa^2}$ denominators

absence of Vainshtein radius

no Vainshtein screening



Nikiforova'20

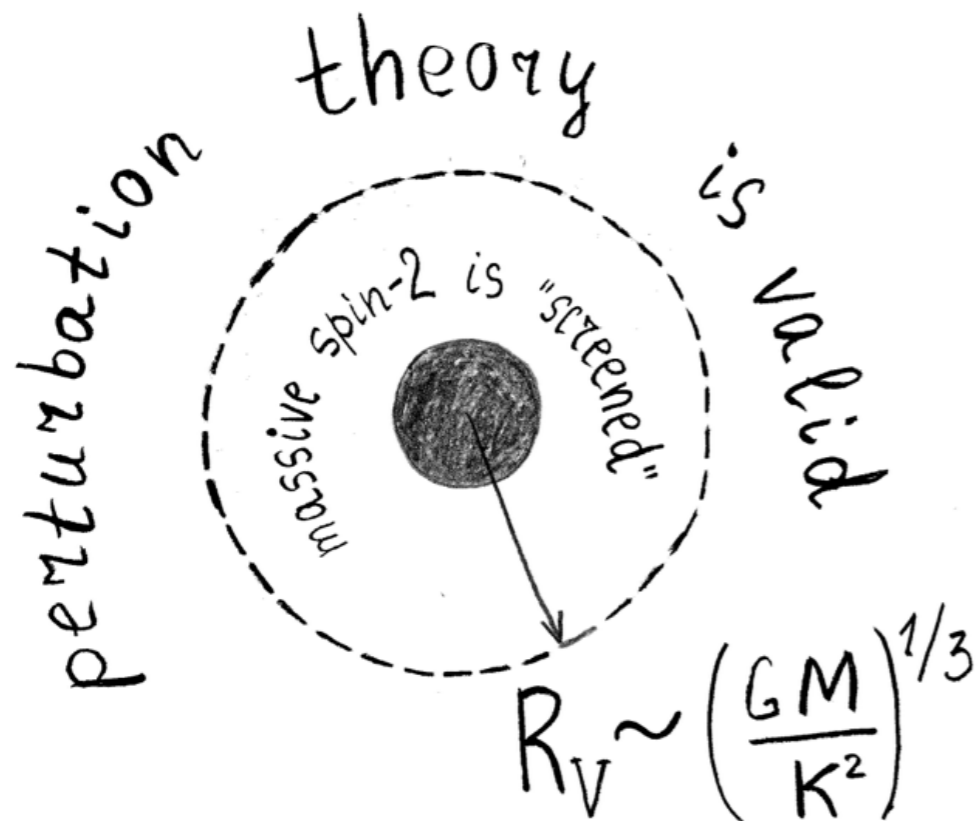
Theories with massive spin-2:

Vainshtein mechanism

bimetric gravity

torsion bigravity

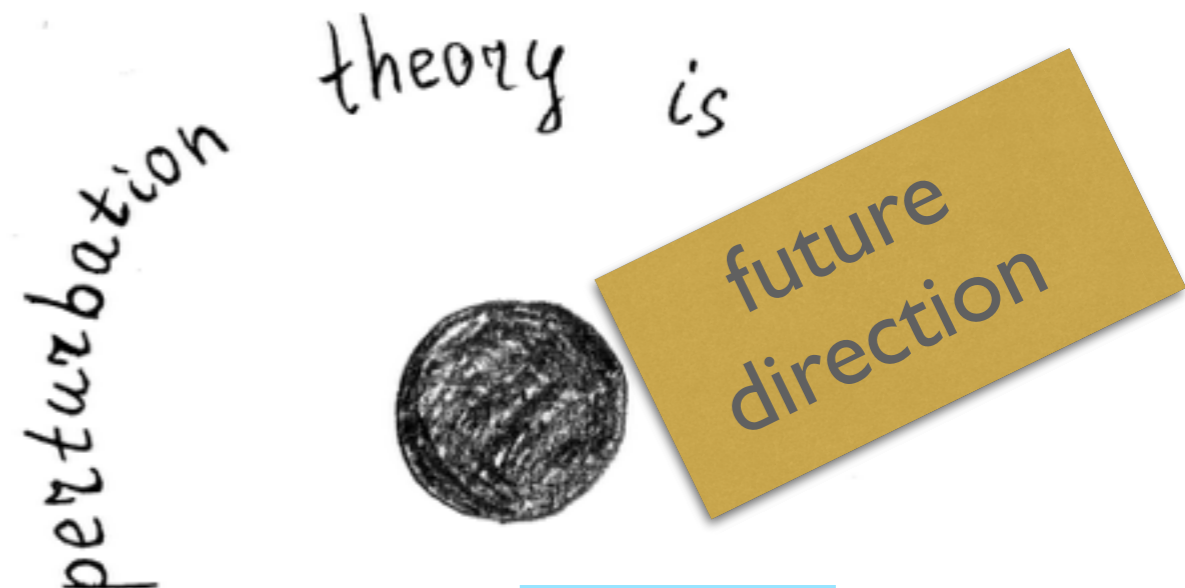
Vainshtein screening



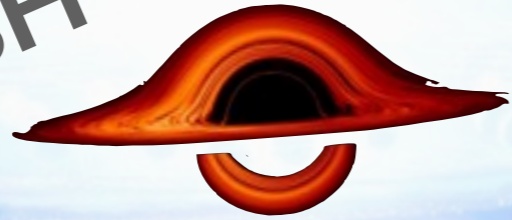
no $\frac{1}{\kappa^2}$ denominators

absence of Vainshtein radius

no Vainshtein screening



BH



Black Holes in torsion bigravity

Theories with massive spin-2:

(Asymptotically flat) black hole ZOO

bimetric gravity

Einstein BHs (incl. Kerr)
are exact solutions

Einstein BHs



(Volkov'12)

BHs
with massive graviton hair



(Brito-Cardoso-Pani'13,
Gervalle-Volkov'20)

Theories with massive spin-2:

(Asymptotically flat) black hole ZOO

torsion bigravity

Einstein BHs (incl. Kerr)
are exact solutions
(with zero torsion)

Einstein BHs



≠ Schwarzschild BHs
with linearized torsion hair

in the limit of
infinite range

$$\kappa \rightarrow 0$$

(different
from
the Einstein
BHs)

asymp. flat
hairy BHs



Theories with massive spin-2:

(Asymptotically flat) black hole ZOO

torsion bigravity

Einstein BHs (incl. Kerr)
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Einstein BHs



⚡ Schwarzschild BHs
with linearized torsion hair
(Nikiforova-Damour'20)

in the limit of
infinite range
 $\kappa \rightarrow 0$

(different
from
the Einstein
BHs)

asymp. flat
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Theories with massive spin-2:

(Asymptotically flat) black hole ZOO

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Einstein BHs



⚡ Schwarzschild BHs
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(Nikiforova-Damour'20)

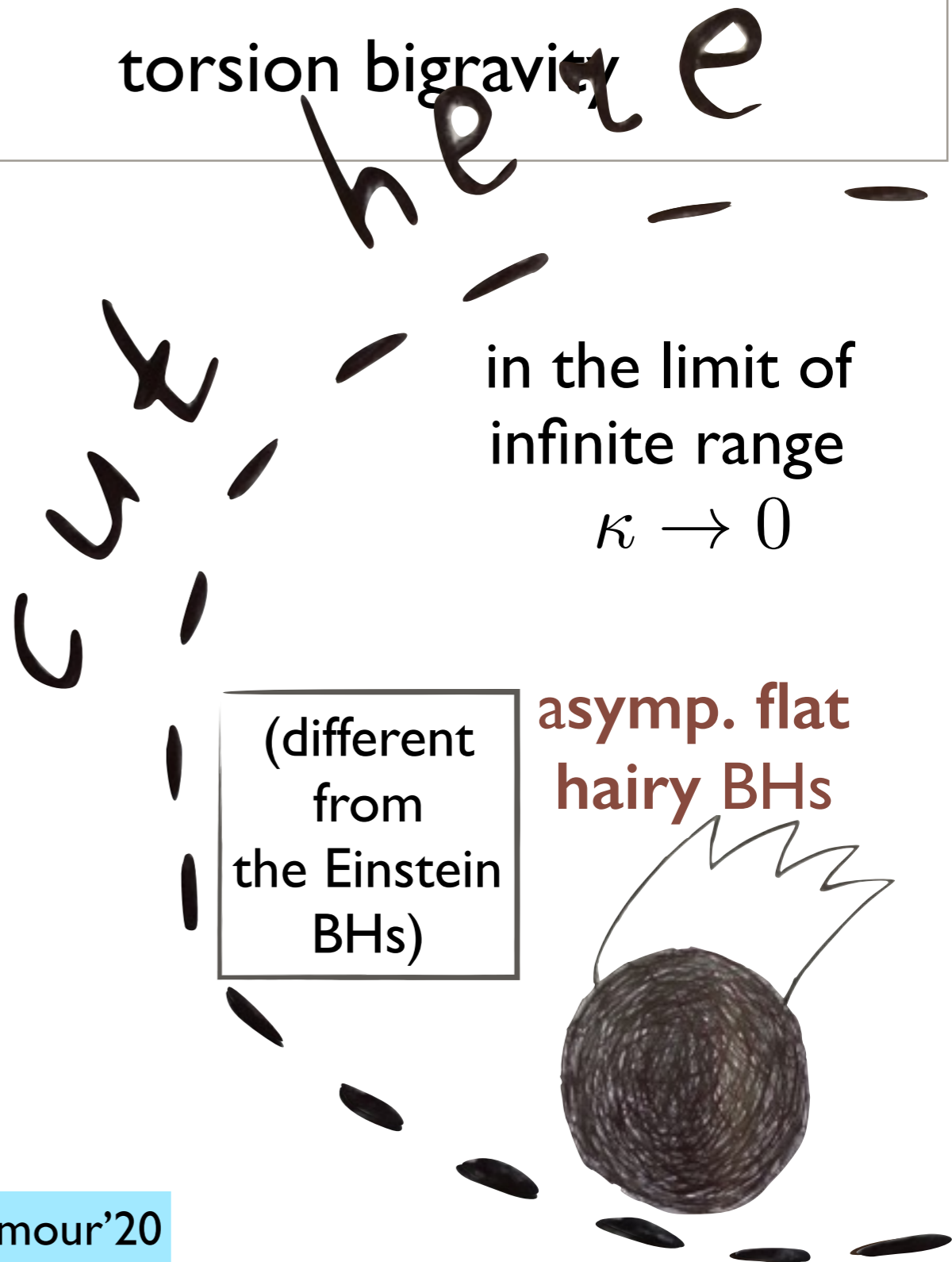
Nikiforova-Damour'20

torsion bigravity

(different
from
the Einstein
BHs)

asymp. flat
hairy BHs

in the limit of
infinite range
 $\kappa \rightarrow 0$



Theories with massive spin-2:

(Asymptotically flat) black hole ZOO

torsion bigravity

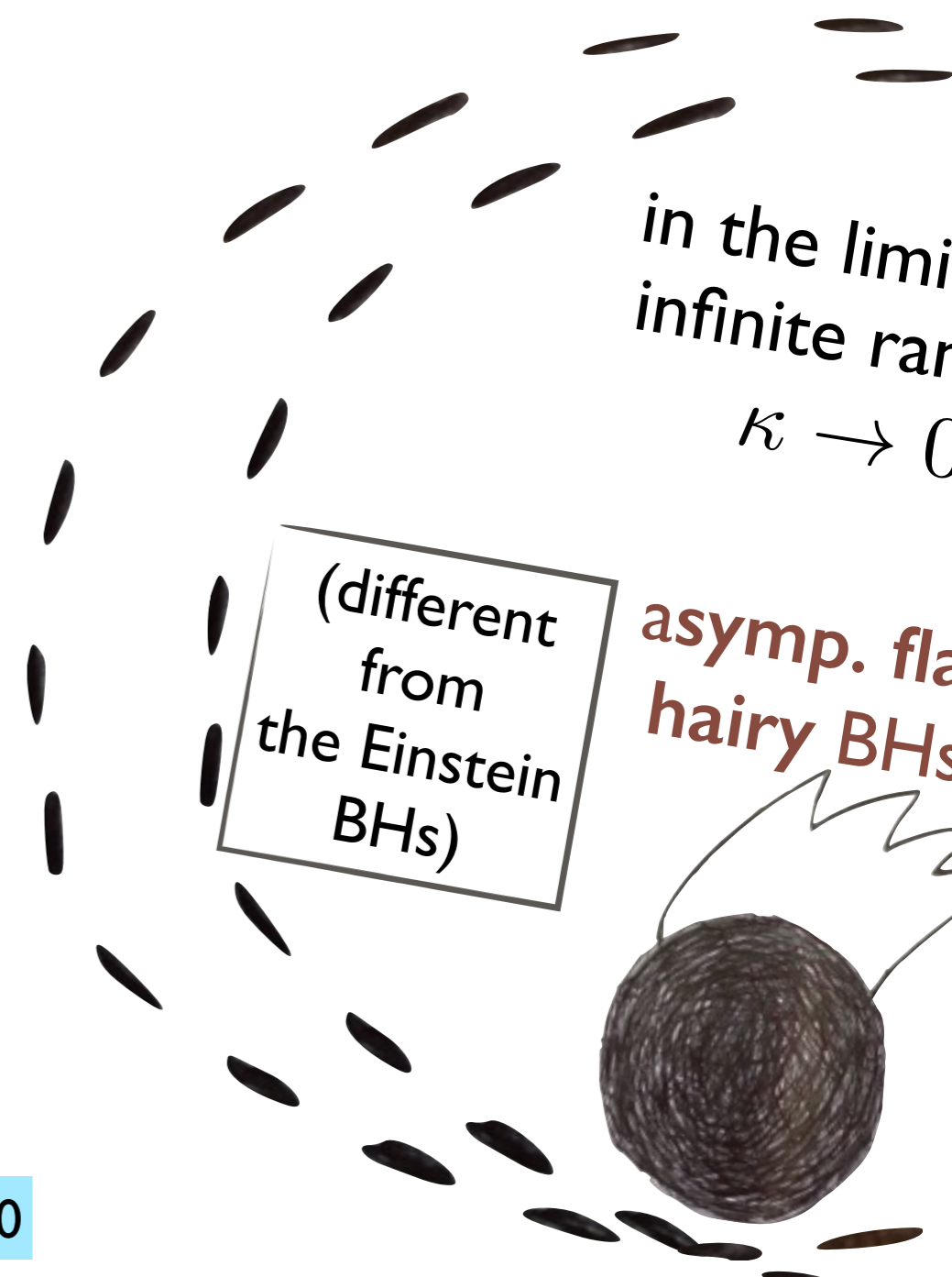
Einstein BHs (incl. Kerr)
are exact solutions
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Einstein BHs



⚡ Schwarzschild BHs
with linearized torsion hair
(Nikiforova-Damour'20)

Nikiforova-Damour'20



Black hole perturbations

Nikiforova'21 + Nikiforova'22 + future works



Black hole perturbations

Nikiforova'21 + Nikiforova'22 + future works



Black hole perturbations

Nikiforova'21 + Nikiforova'22 + future works

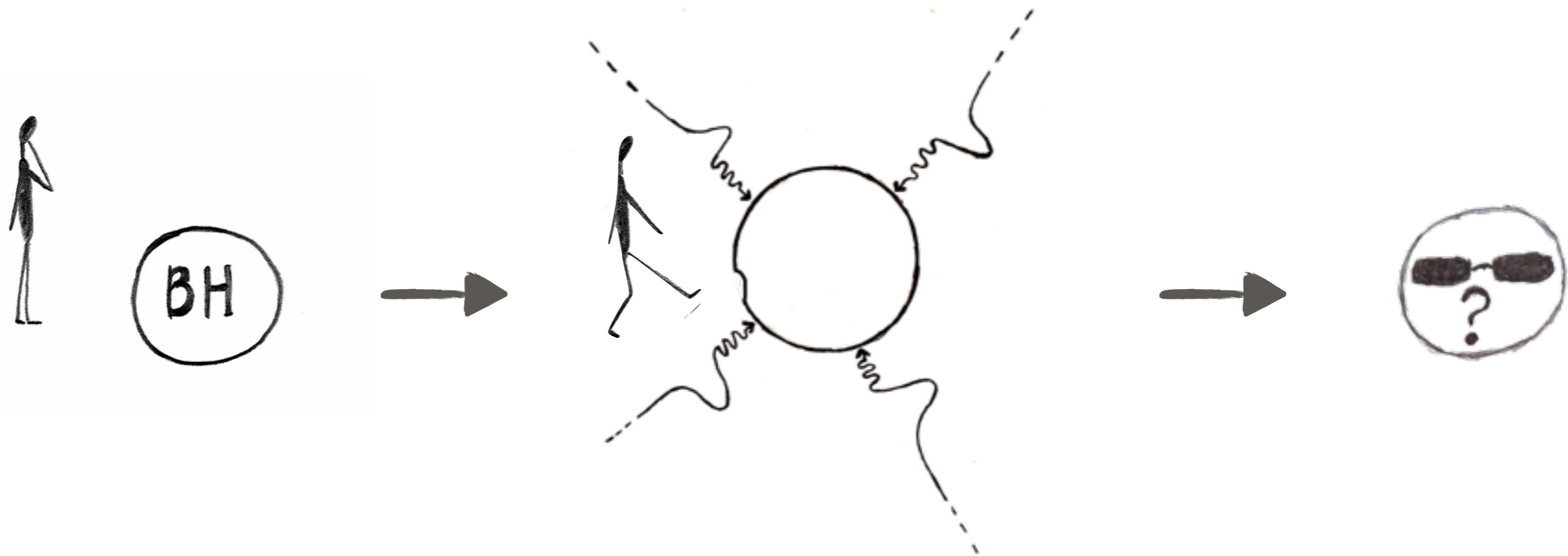
- quasi-bound states and stability
- quasi-normal modes



Black hole perturbations

Nikiforova'21 + Nikiforova'22 + future works

- quasi-bound states and stability
- quasi-normal modes



Theories with massive spin-2: Stability of Schwarzschild BHs

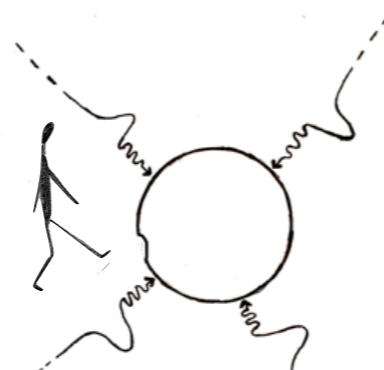
bimetric gravity

Einstein BHs



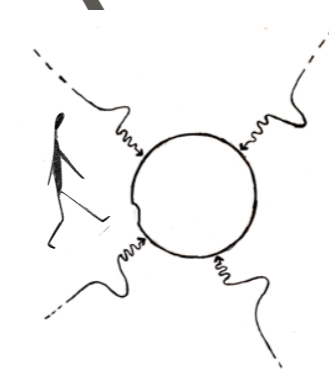
unstable for $\kappa r_h < 0.86$

Babichev-Fabbri'13,
Brito-Cardoso-Pani'13



stable

Gervalle-Volkov'20



unstable
???

Theories with massive spin-2: Stability of Schwarzschild BHs

torsion bigravity

Schw. BH is linearly **stable**

Theories with massive spin-2: Stability of Schwarzschild BHs

torsion bigravity

lower bound on the mass of the massive spin-2
to avoid singularities

$$\kappa > \frac{\sqrt{1 + \eta}}{r_h}$$

Schw. BH is linearly **stable**

Theories with massive spin-2: Stability of Schwarzschild BHs

torsion bigravity

lower bound on the mass of the massive spin-2
to avoid singularities

$$\kappa > \frac{\sqrt{1 + \eta}}{r_h}$$

for $2M_\odot$ black hole:

$$\kappa \gtrsim 10^{-10} \text{ eV} \quad \kappa^{-1} < 6 \text{ km}$$

Superradiant instabilities: NO

Damour-Deruelle-Ruffini'75



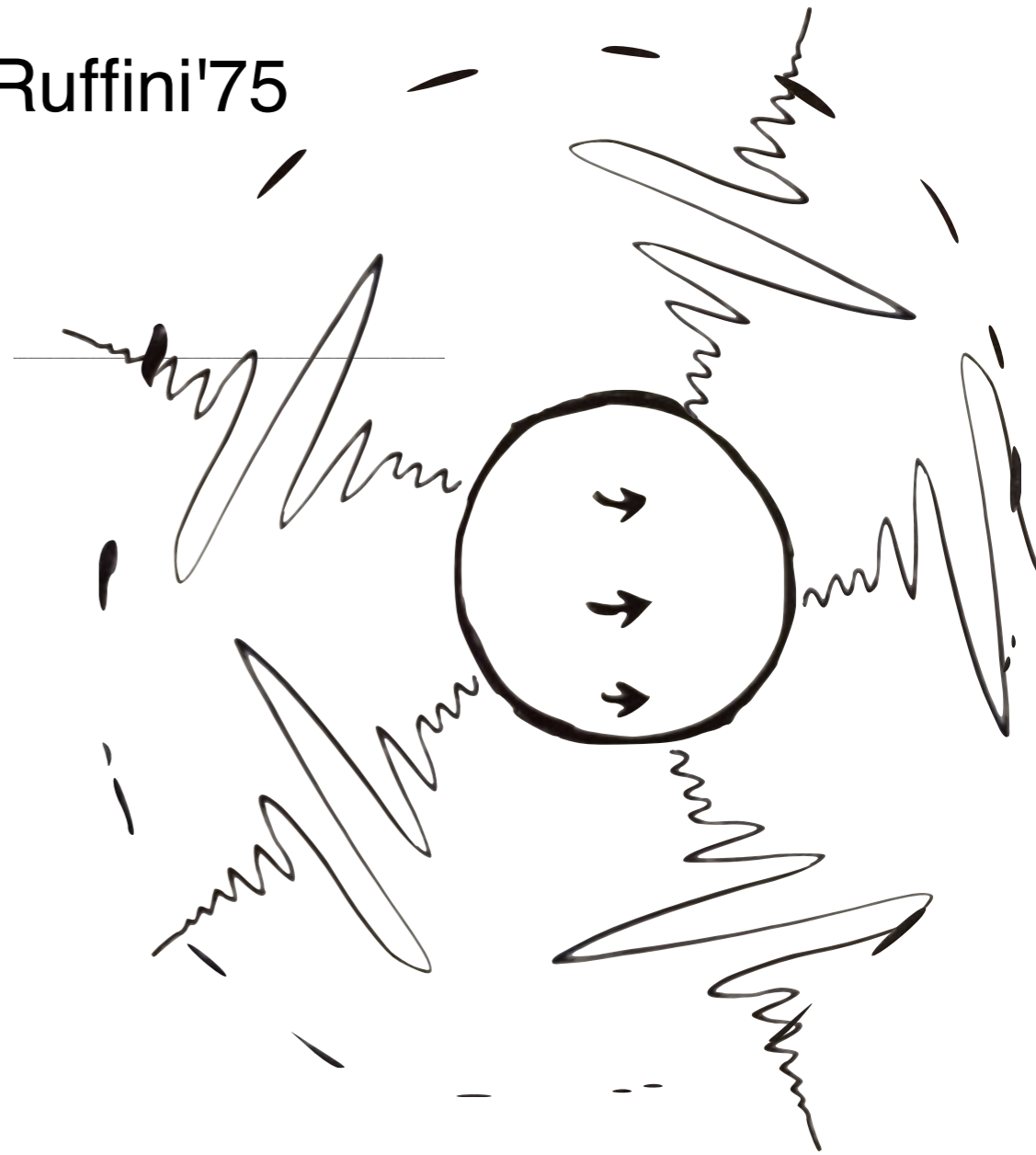
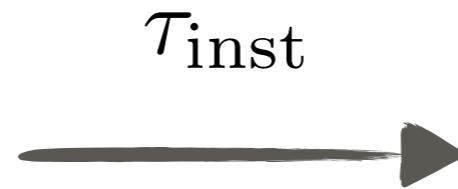
Kerr

Superradiant instabilities: NO

Damour-Deruelle-Ruffini'75

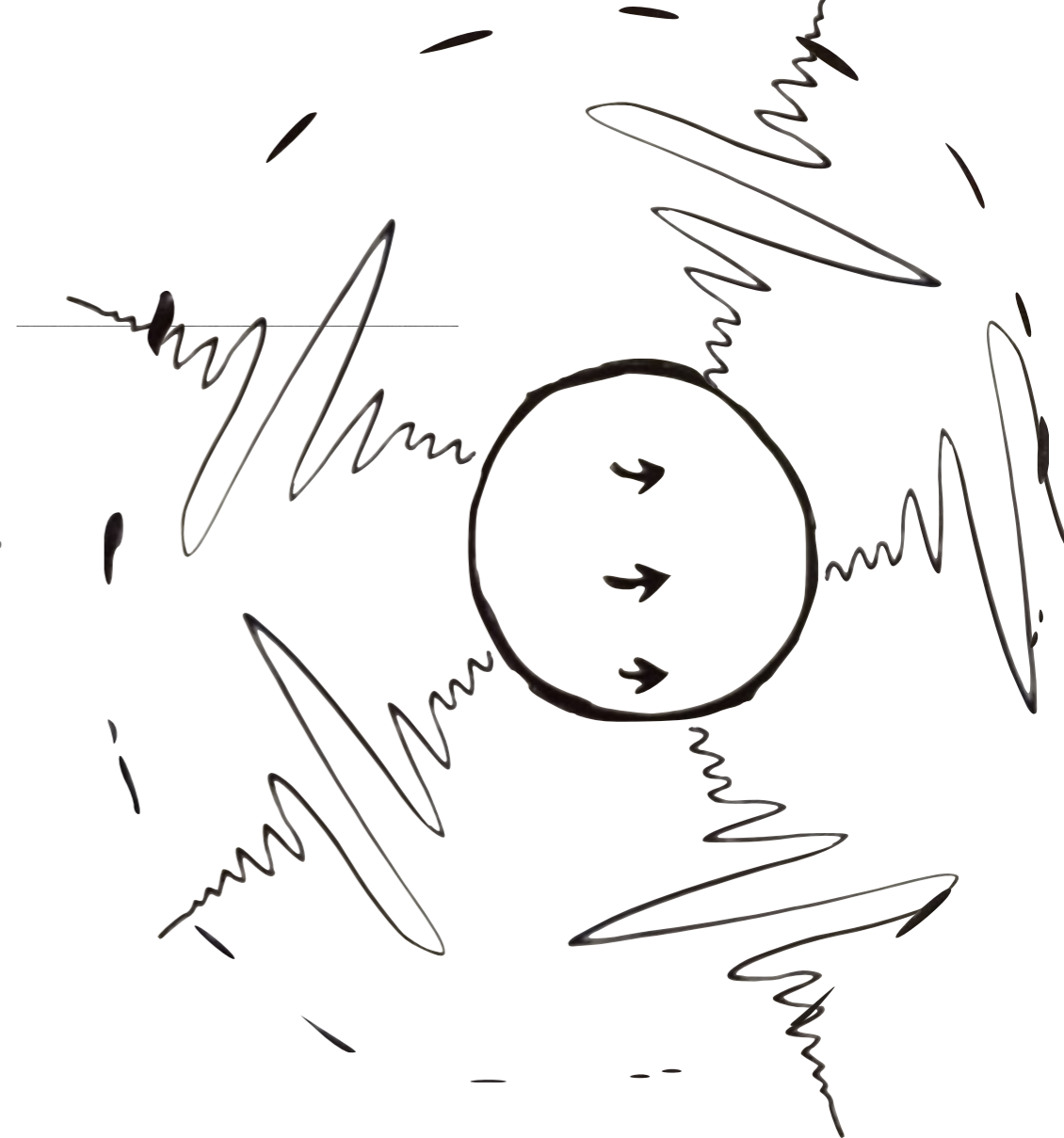
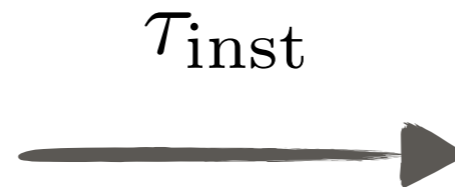
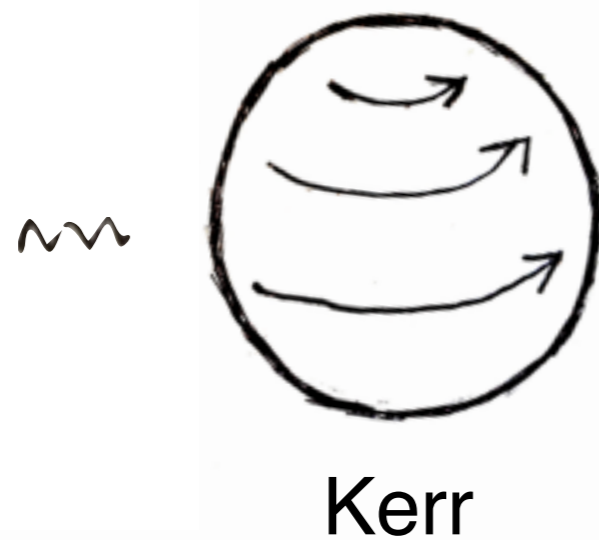


Kerr



Superradiant instabilities: NO

Nikiforova'22 based on Zouros-Eardley'79 , Stott'09 and Brito-Grillo-Pani'21



- Salpeter time
- time of observation of a spinning black hole

EXCLUDED: $10^{-13} \text{ eV} \lesssim \kappa \lesssim 10^{-11} \text{ eV}$

Superradiant instabilities: NO

Nikiforova'22 based on Zouros-Eardley'79 , Stott'09 and Brito-Grillo-Pani'21

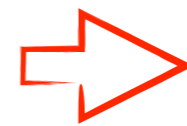


Kerr

τ_{inst}



$$\kappa > \frac{\sqrt{1 + \eta}}{r_h}$$



$$\kappa \gtrsim 10^{-10} \text{ eV}$$

no observable superradiant instabilities

- Salpeter time
- time of observation of a spinning black hole

EXCLUDED: $10^{-13} \text{ eV} \lesssim \kappa \lesssim 10^{-11} \text{ eV}$

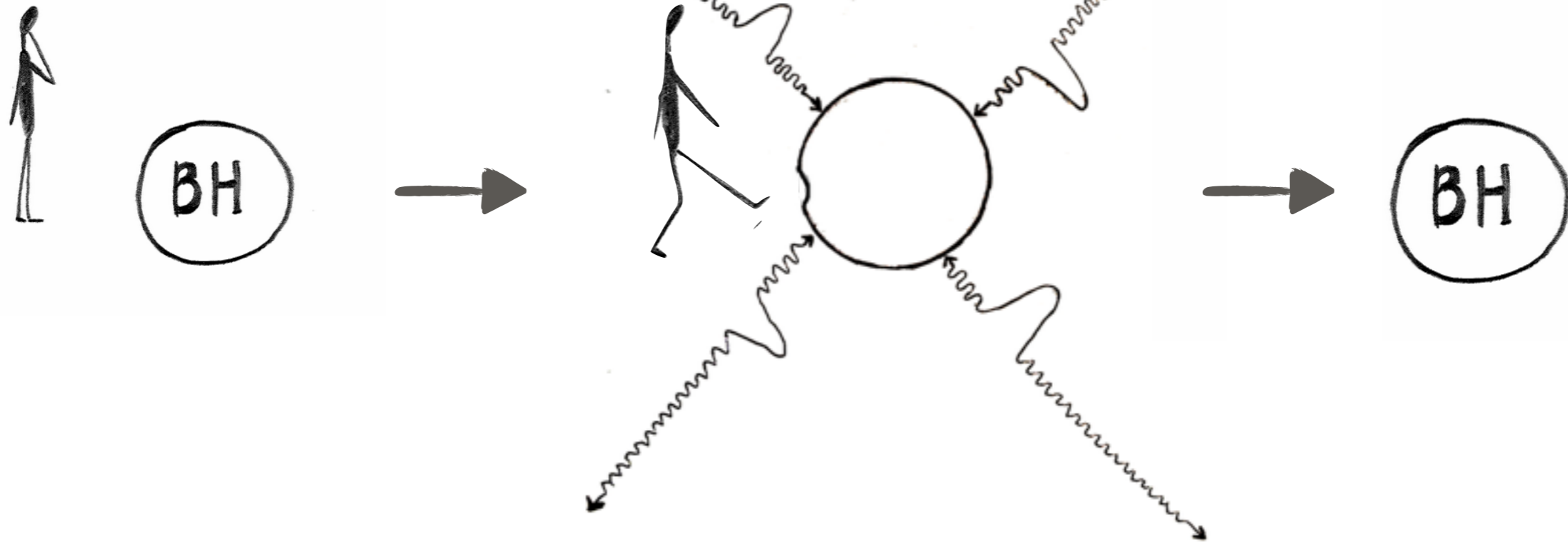
Black hole perturbations

Nikiforova'21 + work in progress

- quasi-bound states and stability
- **quasi-normal modes**

future
direction

LIGO-Virgo-Kagra



Take away message

- it's a good time to do gravity. good time to formulate questions, good time to look beyond GR
- gravity with dynamical torsion is a natural extension of GR
- gravity with dynamical torsion gives an interesting alternative to GR, which can be useful both for phenomenology and for theoretical development of gravity.
TO THIS END, I will continue study.

A 5-parameter class of dynamical torsion theories revived with cosmological motivation

Nair—Randjbar-Daemi—Rubakov (2009)
V. Nikiforova, S. Randjbar-Daemi, V. Rubakov (2009)
Deffayet—Randjbar-Daemi (2011)
V. Nikiforova, S. Randjbar-Daemi, V. Rubakov (2016)
V. Nikiforova (2017)
V. Nikiforova, T. Damour (2018)

self-accelerating solution

where torsion accelerates the Universe

but

instabilities found.