

QUANTUM PHYSICS : OUTLINE OF A PROGRAMME

Ph. Baudrand

- QM incredibly successful \rightarrow Standard QM clearly works pragmatically
 \rightarrow Justification FAPP (for all practical purposes) of the Quantum Rules starting from von Neumann's epistemic principle (1932)

- A "Unified Algebraic Framework" valid just as well for C- and Q- Physics (Segal Heisenberg, Jordan, von Neumann, Segal, Haag-Kastler, Araki)
 Algebra of observables

Physical system $\Leftrightarrow (\mathcal{A}, \varphi, \alpha_{\varphi})$

States (preparation procedure)
 $\varphi: \mathcal{A} \rightarrow \mathbb{C}$ $\varphi(\mathbb{I}) = 1$
 Dynamic linear, positive $\varphi(A^*A) \geq 0$, continuous
 Automorphism $\alpha_{\varphi}: \mathcal{A} \rightarrow \mathcal{A}$

Connections to experiments

Experiments more and more sophisticated
 History Non-demolition experiments

$$A = A^* \rightarrow \varphi(A) \text{ real} \quad \text{Expectation } (A \text{ in state } \varphi)$$

To be computed

$$\text{Prob} [A, \varphi \mid \varphi(A) \in I \subset \mathbb{R}]$$

History in the sense of the history books

$$\text{Prob} [\text{History}] = \text{Prob} [P_n P_{n-1} \dots P_1]$$

C^* -Algebra \mathcal{A}

V. Neumann algebra $\mathcal{M} = \mathcal{M}'' \subset \mathcal{B}(\mathcal{H})$

Algebraic Probability Theory (\mathcal{M} , φ)

φ faithful normal state \approx non-commutative measure

GNS $\pi_{\varphi}(A)$ in $\mathcal{B}(\mathcal{H}_{\varphi}) \Rightarrow$ recovering the Hilbert space
 $\langle \cdot, \cdot \rangle_{\mathcal{H}_{\varphi}} = \varphi(A^*B)$

\mathcal{M} commutative \Leftrightarrow Kolmogorov probability theory (1933)

The 2 logics of science

\mathcal{M} non-commutative

"Quantum" Probability Theory

From Q to C : Decoherence

(emergent phenomenon: multiscal environment, time)

S + E closed system

$$\mathcal{H}_{S+E} = \mathcal{H}_S \otimes \mathcal{H}_E$$

$$\mathcal{N} = \mathcal{M} \otimes \mathcal{M}_E$$

$\alpha_E: \mathcal{N} \rightarrow \mathcal{N}$ automorphism

ω_E Reference state of E

$$i: \mathcal{M} \rightarrow \mathcal{N} \text{ embedding } i(A) = A \otimes \mathbb{I}_E$$

$E_{\omega_E}: \mathcal{N} \rightarrow \mathcal{M}$ conditional expectation

$$T_t = E_{\omega_E} \circ \alpha_E \circ i$$

g \otimes ω_E state on \mathcal{N}

$$T_t(s) = (E_{\omega_E} \circ \alpha_E \circ i)(s \otimes \omega_E)$$

T_t is completely positive as combination of 3 CP-maps

① Mathematical definition of decoherence B.O. 15 2003

Splitting of $M = M_1 \oplus M_2$
 $\subset v.N.$ subalgebra of M

$$A \in M \quad A = A_1 \oplus A_2$$

M_1 \approx effective observable, decoherence free part, can be detected in practice

$\beta_t = \alpha_t |_{M_1}$ is an automorphism

$M_2 \approx$ not detectable by measurements for $t \rightarrow \infty$

$$\lim_{t \rightarrow \infty} \varphi(\alpha_t(A_2)) = 0$$

• Automorphism $\alpha_t : M \otimes M_E$

\downarrow E_{ω_E} : conditional expectation \sim averaging over M_E

• CP-map $T_t : M \rightarrow M$

\downarrow asymptotic decoherence (Happ, ...)

$$\lim_{t \rightarrow \infty}$$

• Automorphism
FAPP

$$\beta_t = T_t |_{M_1}$$

Decoherence is nothing else as the application of the standard formulation of Q-Physics

② Scenarios of Decoherence Attractors of the Dynamics

$$\mathcal{Z}(M) = M \cap M' = \mathbb{C}\mathbb{I} \Leftrightarrow S \text{ genuine (purely) } Q\text{-system}$$

• Pointer states M_1 commutative $\approx L^\infty(S^2)$

$$\{\beta_t\}_{t \in \mathbb{R}} \text{ trivial } \beta_t = \mathbb{I}$$

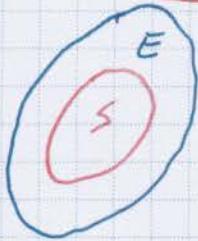
• SSR M_1 non-commutative
 $\mathcal{Z}(M_1)$ non-trivial \Leftrightarrow superselection observables

• From Q to q M_1 non-commutative and
 $\mathcal{Z}(M_1) = \mathbb{C}\mathbb{I}$

• Classical system M_1 commutative

• Ergodicity $M_1 = \mathbb{C}\mathbb{I}$ M_1 is trivial

The notion of Decoherence in the algebraic framework



Closed system $S \times E$ described by

- M von Neumann algebra
- continuous group of automorphisms $\{\beta_t\}_{t \in \mathbb{R}}$
- S described by a v.N. subalgebra $M \subseteq M$
- \exists a normal conditional expectation $E : M \rightarrow M$

\Rightarrow Define the reduced dynamics as follows

$$T_t(x) = E \circ \beta_t(x) \quad x \in M \quad t \geq 0$$

This is the Heisenberg picture time evolution for an observer with experimental capabilities limited to M .

T_t is in general irreversible (S is an open system)

Mathematical properties of T_t

1. $\{T_t\}_{t \geq 0}$ family of CP normal linear maps on M
2. $T_t(\mathbb{1}) = \mathbb{1} \quad \forall t \geq 0 \quad \|T_t x\| \leq \|x\|$
3. $t \rightarrow T_t(x)$ is (weak) continuous
4. In general $\{T_t\}_{t \geq 0}$ is not Markovian.

Remark In some physically relevant situations \exists good approximation to describe T_t by a semigroup \tilde{T}_t

$\tilde{T}_t \Rightarrow Q$ dynamical semi group

Models

$$M = M_0 \otimes M_0 \quad \text{acting on } \mathcal{H} \otimes \mathcal{H}_0$$

$$\begin{matrix} & \uparrow & \uparrow \\ & S & E \\ \downarrow & & \downarrow \\ & \tilde{S} & \tilde{E} \end{matrix}$$

$$\beta_t(x) = e^{-itH} x e^{itH}$$

$$H = H_0 \otimes \mathbb{1} + \mathbb{1} \otimes H_0 + H_{int}$$

ω reference state of E

$$\varphi \otimes \omega(x) = \varphi(E_\omega(x)) \quad x \in M$$

φ state on M

Time evolution

$$T_t(\varphi) = \text{tr}_2 [e^{-itH} (\varphi \otimes \omega) e^{itH}]$$

φ normal state on M

tr_2 partial trace with respect to the degrees of freedom of E

T_t not reversible \Rightarrow new phenomena can appear

- approach to equilibrium

Definition of Decoherence

The reduced dynamics $\{T_t\}_{t \geq 0}$ displays decoherence if

- \exists a T_t invariant v. Neumann subalgebra M_1 of M and a group of automorphisms $\{\alpha_t\}_{t \geq 0}$ on M_1 such that

$$T_t|_{M_1} = \alpha_t \quad \text{for } t \geq 0 \quad \text{The effective dynamics}$$

- \exists a T_t invariant Banach subspace M_2 of M such that

$$M = M_1 \oplus M_2$$

with $\lim_{t \rightarrow \infty} T_t(x) = 0 \quad \forall x \in M_2$

We require that M_1 is a maximal v. N. subalgebra of M (not contained in any larger von Neumann subalgebra on which $\{T_t\}_{t \geq 0}$ extends to an automorphism)

M_1 is called the algebra of effective observables

Physical interpretation

$$x \in M \quad x = x_1 \otimes x_2 \quad x_1 \in M_1, x_2 \in M_2$$

$$\lim_{t \rightarrow \infty} \varphi(T_t(x_2)) = 0 \quad \Rightarrow \quad M_2 = 0 \quad \text{FAPP}$$

After sufficiently long time S behaves effectively (FAPP) like a closed system described by M_1 and $\{T_t\}_{t \geq 0}$

Question : Is M_1 unique ?

Yes if T_t is a QDS

In certain cases the splitting $M = M_1 \oplus M_2$ is always given by a conditional expectation

Proposition $\{T_t\}_{t \geq 0}$ displays decoherence and $T_t|_{M_1} = \pi_{M_1}$ then there exists a normal conditional expectation E from M into M_1

Problem Given T_t under which conditions decoherence appears ?

- Markovian case CMP 239 241-255 (2003) LO

Theorem let $\{T_t\}_{t \geq 0}$ be a continuous one parameter SG with a faithful normal state ω and that

$$\omega \circ T_t = \omega \quad \forall t \geq 0$$

and assume that

1. T_t $t \geq 0$ is strongly positive and unital

2. $\{\sigma_t^\omega\}_{t \in \mathbb{R}}$ the modular group corresponding to

then $[T_t, \sigma_s^\omega] = 0 \quad \forall s \in \mathbb{R}$ and $t \geq 0$

Then $\{T_t\}_{t \geq 0}$ displays decoherence and \exists a normal conditional expectation $E : M \rightarrow M$ and that $[T_t, E] = 0$ and $\omega \circ E = \omega$

References sur la décohérence

Wroclaw { Olkiewicz
Lugiewicz

Wroclaw - Bielefeld

Bielefeld { R. B.
Hellmich

5. II. 2015

- B.O Decoherence induced transition from Q to C dynamics Rev. Math. Phys 15 217-243
- B.L.O From Q to C via decoherence Phys Lett A 314 29-36 (2003)
- B.O Decoherence induced continuous pointer states PRL 90 010403 (2003)
- B.H.L.O Quantum dynamical semigroup for finite and in finite Bose systems J. Math. Phys 48 012106 (2007)
- B.H.L.O Continuity and generators of dynamical semigroups for infinite Bose systems J. Funct. Analysis 256 1453-1475 (2005) JFA 259 2455-2456 (2010)
- O Environment - induced Superselection rules in Markovian regime CMP 208 245-265 (1999)
- L.O Decoherence in infinite quantum systems J. Phys. A Math. Gen. 35 6695-6712 (2002)
- L.O Classical properties of infinite quantum open systems
- B.O Decoherence as Irreversible Dynamical Process in Open Quantum Systems LMP 1882 Springer 117-160 (2006)
- B.H Decoherence in Infinite Quantum Systems Quantum Africa 2010 AIP Conference Proceedings 1469 2-15 (2012)
- B.H.L.O Theory of the decoherence effect in finite and infinite Open Quantum Systems using the algebraic approach