



Perturbing binary black holes with effective field theory

Based on 2004.03570

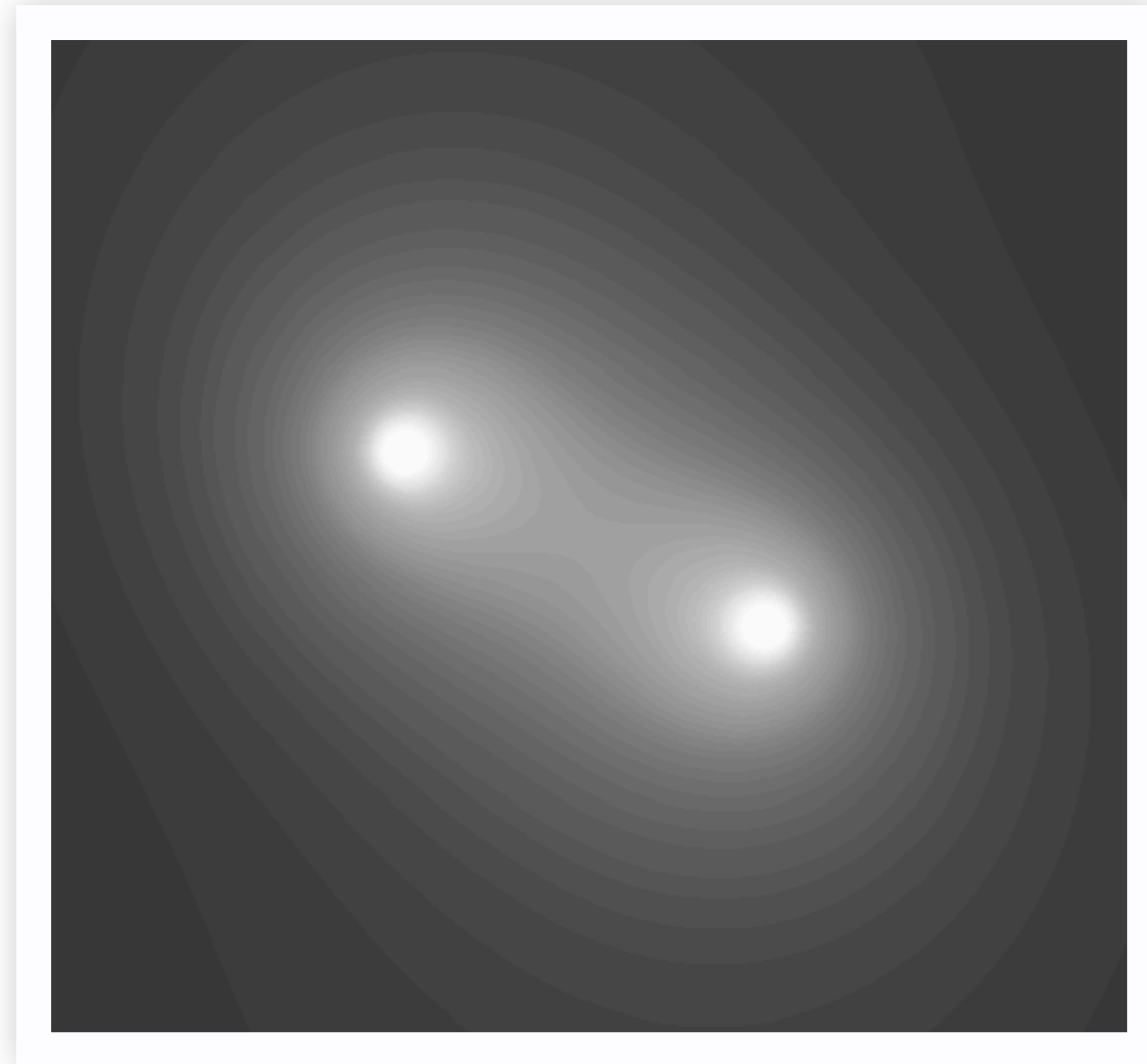
LEONG KHIM WONG

MEETING OF THE NATIONAL RESEARCH
GROUP ON GRAVITATIONAL WAVES

1 APRIL 2021

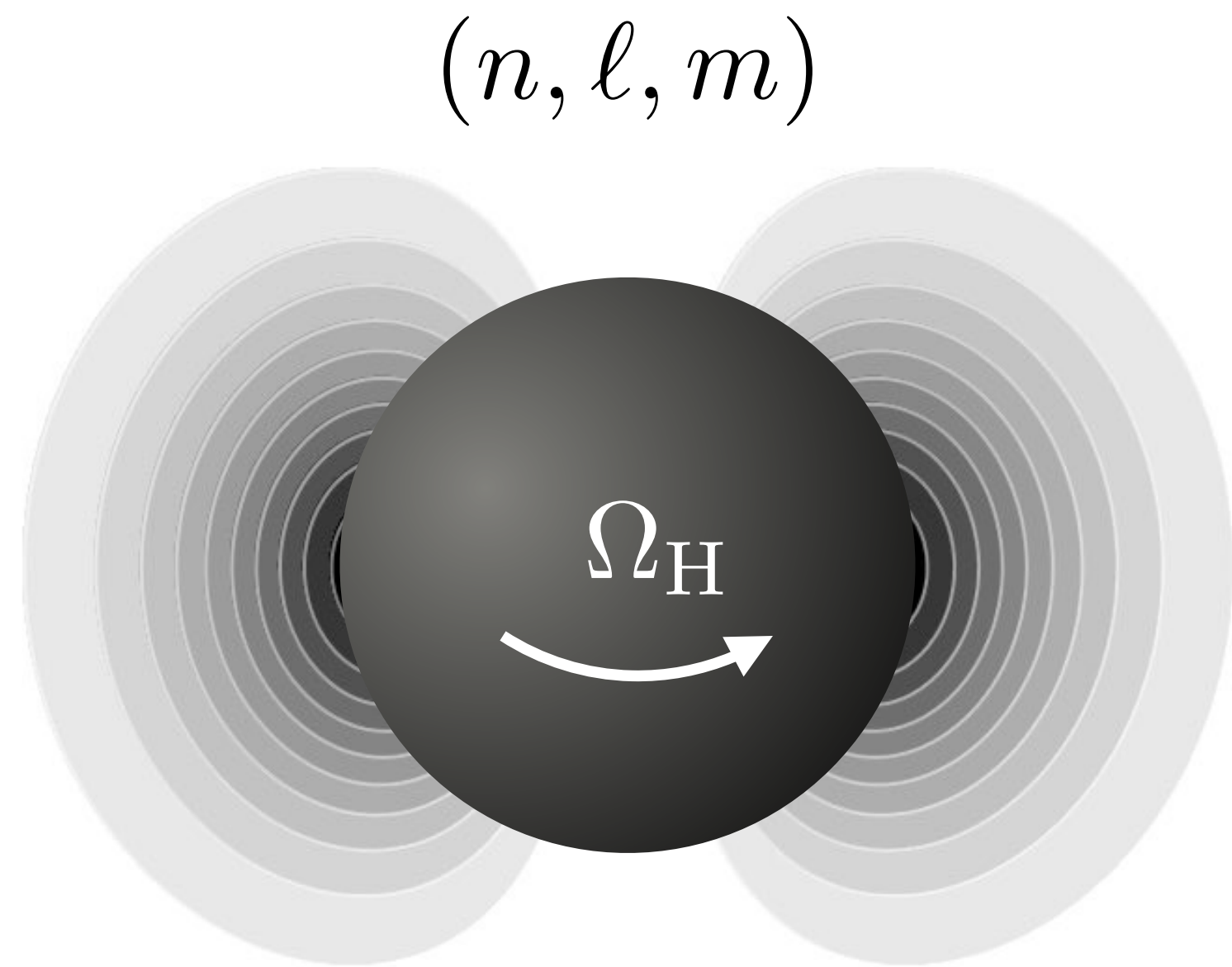
GRAVITATIONAL MOLECULES

General relativity + ultralight boson



[Ikeda, Bernard, Cardoso, Zilhão 2020]

GRAVITATIONAL ATOM

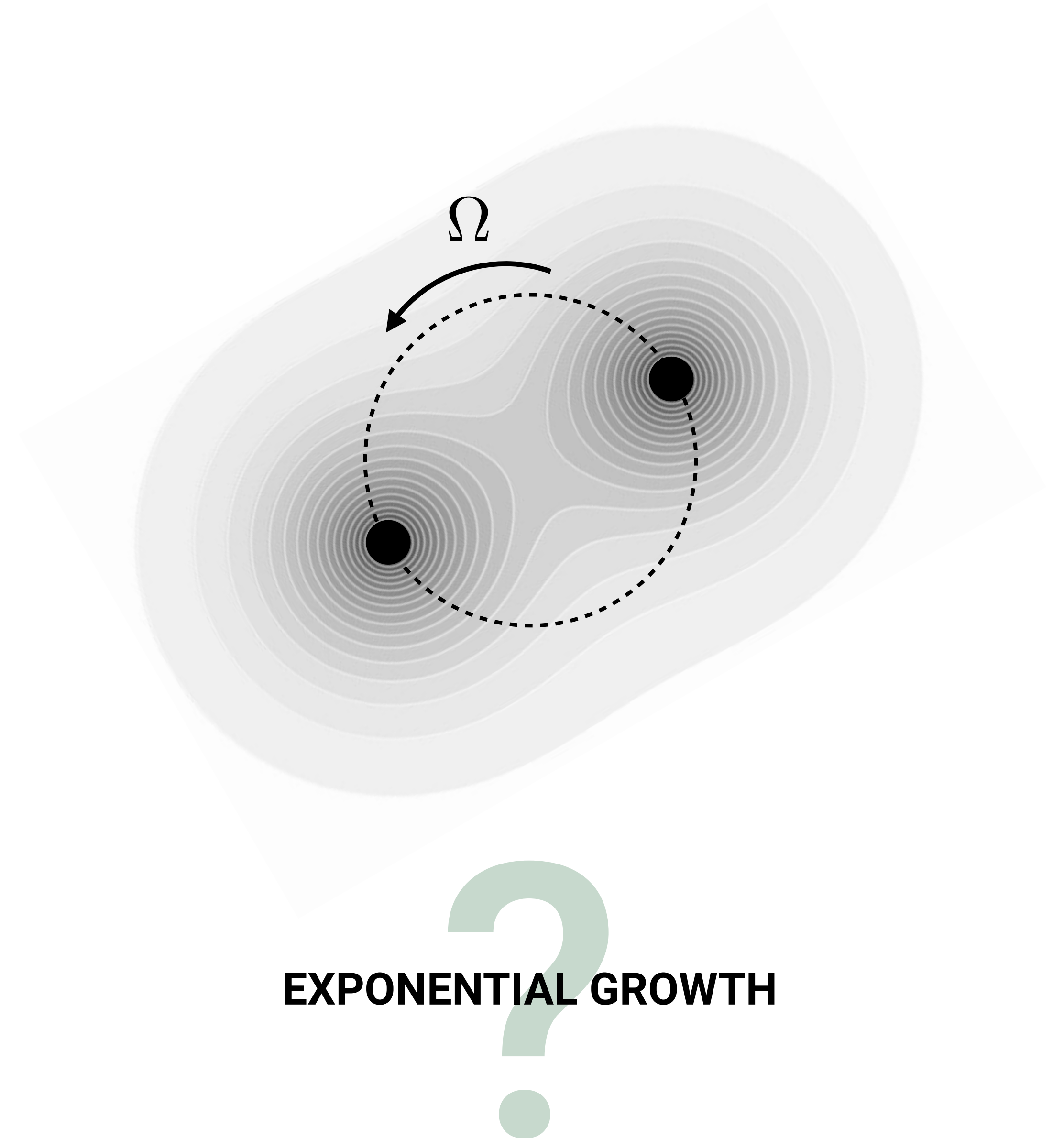


$$0 < E_{n\ell m} < m\Omega_H$$

EXPONENTIAL GROWTH
($E_{n\ell m} = \mu + \text{binding energy}$)

[Detweiler 1980; Dolan 2007; etc.]

GRAVITATIONAL MOLECULE



EXPONENTIAL GROWTH

SIMPLIFYING ASSUMPTIONS

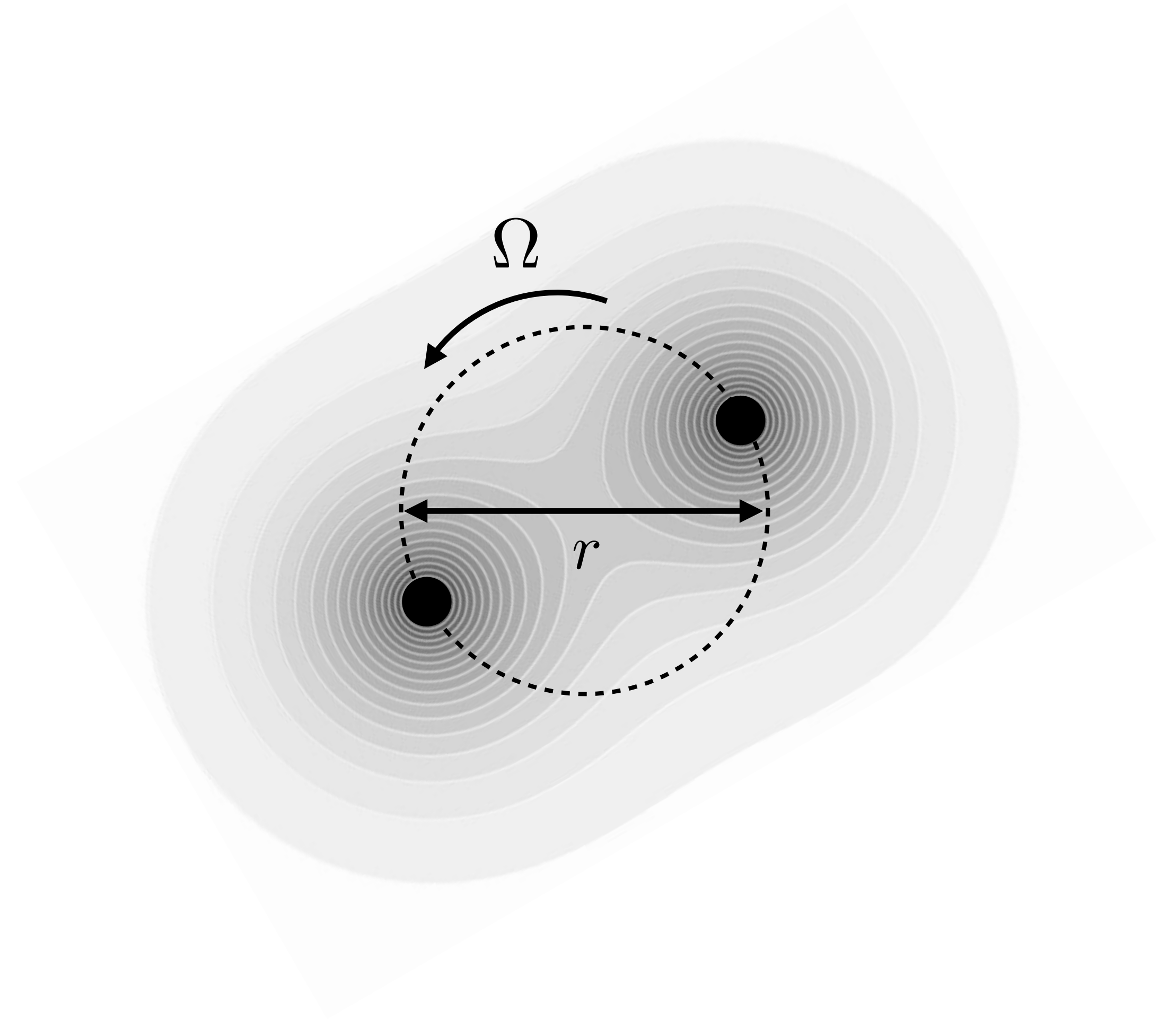
- 1 Low energy density:** Neglect backreaction of the cloud to a first approximation
- 2 Early inspiral stage:** Post-Newtonian approximation of the (circular) binary

$$v^2 \sim \frac{GM}{r} \ll 1$$

- 3 Long-wavelength limit:** Treat the entire binary as an effective point particle

$$r/\lambda \ll 1 \quad \Leftrightarrow \quad \mu \ll \Omega/v^2$$

$$\mu \ll 10^{-13} \text{ eV} \left(\frac{v}{0.1} \right) \left(\frac{100 M_\odot}{M} \right)$$



SIMPLIFYING ASSUMPTIONS

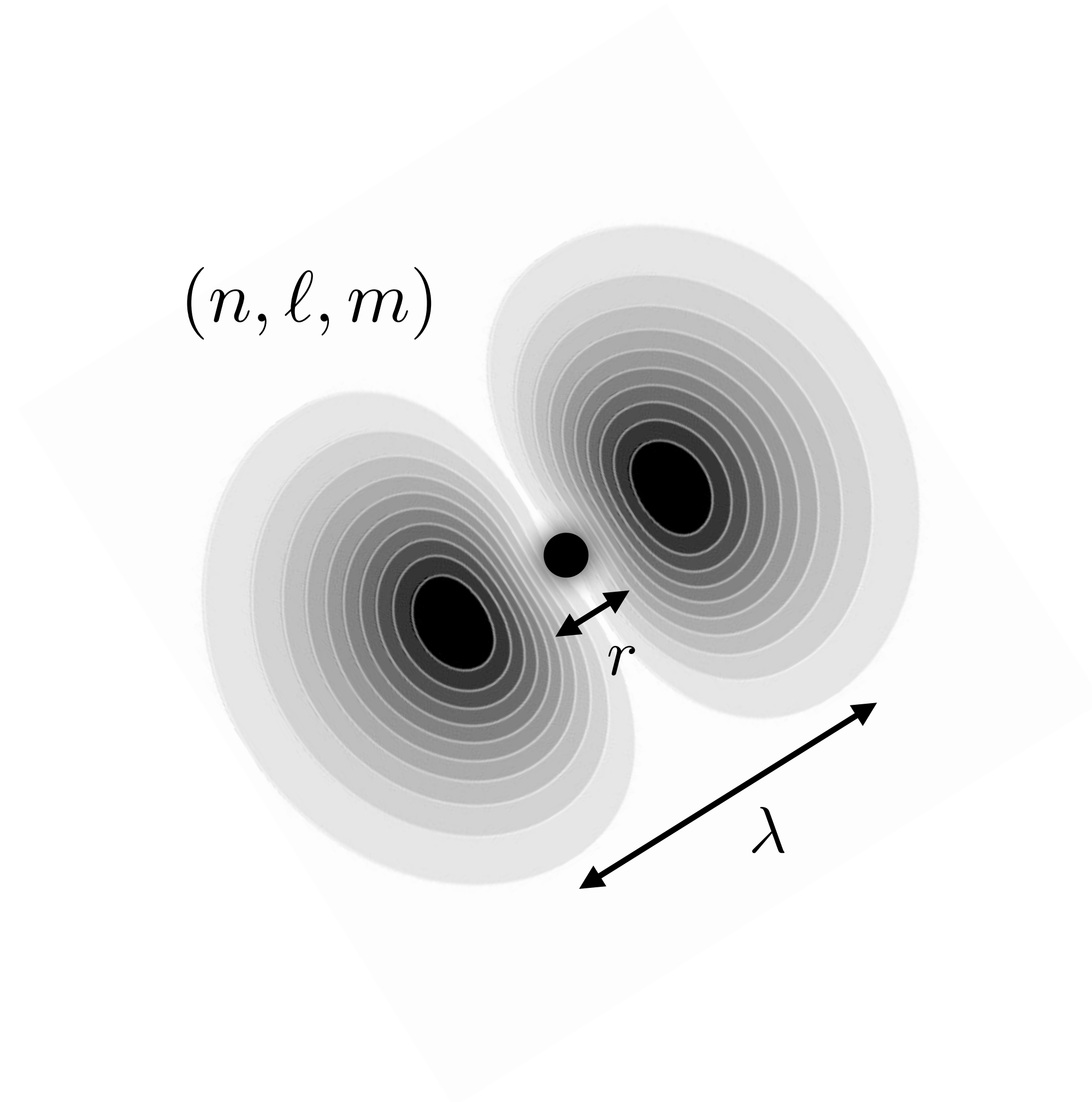
- 1 Low energy density:** Neglect backreaction of the cloud to a first approximation
- 2 Early inspiral stage:** Post-Newtonian approximation of the (circular) binary

$$v^2 \sim \frac{GM}{r} \ll 1$$

- 3 Long-wavelength limit:** Treat the entire binary as an effective point particle

$$r/\lambda \ll 1 \quad \Leftrightarrow \quad \mu \ll \Omega/v^2$$

$$\mu \ll 10^{-13} \text{ eV} \left(\frac{v}{0.1} \right) \left(\frac{100 M_\odot}{M} \right)$$



EFFECTIVE EQUATION OF MOTION

The evolution of the cloud is given by solving $\square_{\text{eff}}\phi = 0$, where

$$\square_{\text{eff}} = \underbrace{\left(\eta^{\mu\nu} \partial_\mu \partial_\nu - \mu^2 + \frac{2GM\mu^2}{r} \right)}_{\square_{\text{free}}} + \underbrace{\left(\sum_{\ell=0}^{\infty} \partial_{i_1} \cdots \partial_{i_\ell} \delta^{(3)}(\mathbf{x}) \mathcal{G}^{i_1 \cdots i_\ell}(t) \right)}_{\square_{\text{int}}} + \cdots$$

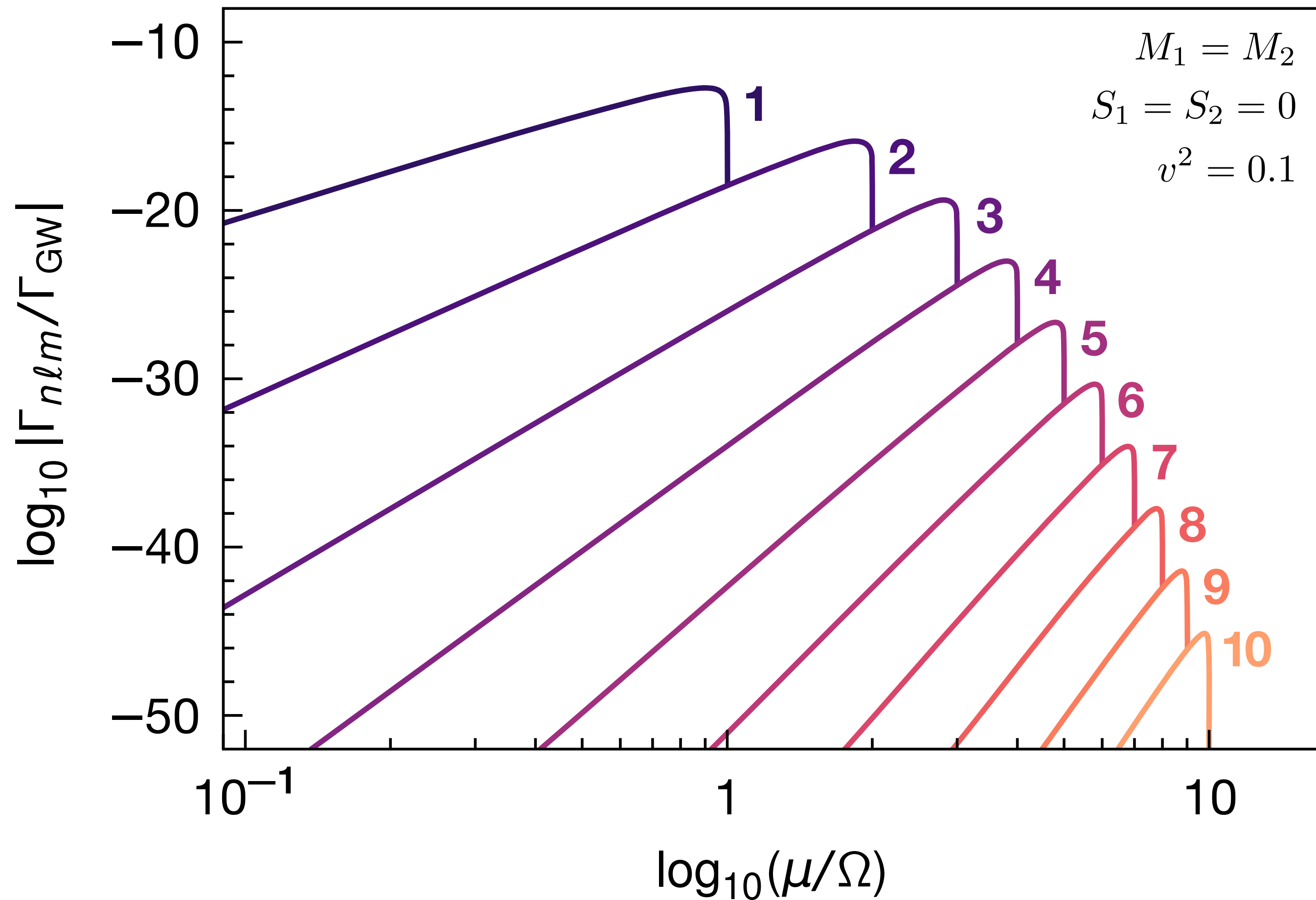
Leading dynamics of the cloud
in the bulk of the spacetime

Contact interactions describing absorption
by and energy-momentum transfer with the
moving black hole horizons

[Wong et al. 1903.07080, 1905.08543, 2004.03570]

$$\mathcal{G}^{i_1 \cdots i_\ell}(t) = - \lim_{\mathbf{x} \rightarrow 0} \sum_{N=1}^2 \sum_{\ell'=0}^{\infty} \frac{A_N}{\ell'!} \mathbf{z}_N^{\langle i_1} \cdots \mathbf{z}_N^{i_\ell \rangle} \frac{d}{dt} \left[\mathbf{z}_N^{\langle j_1} \cdots \mathbf{z}_N^{j_{\ell'} \rangle} \partial_{j_1} \cdots \partial_{j_{\ell'}} \cdot \right]$$

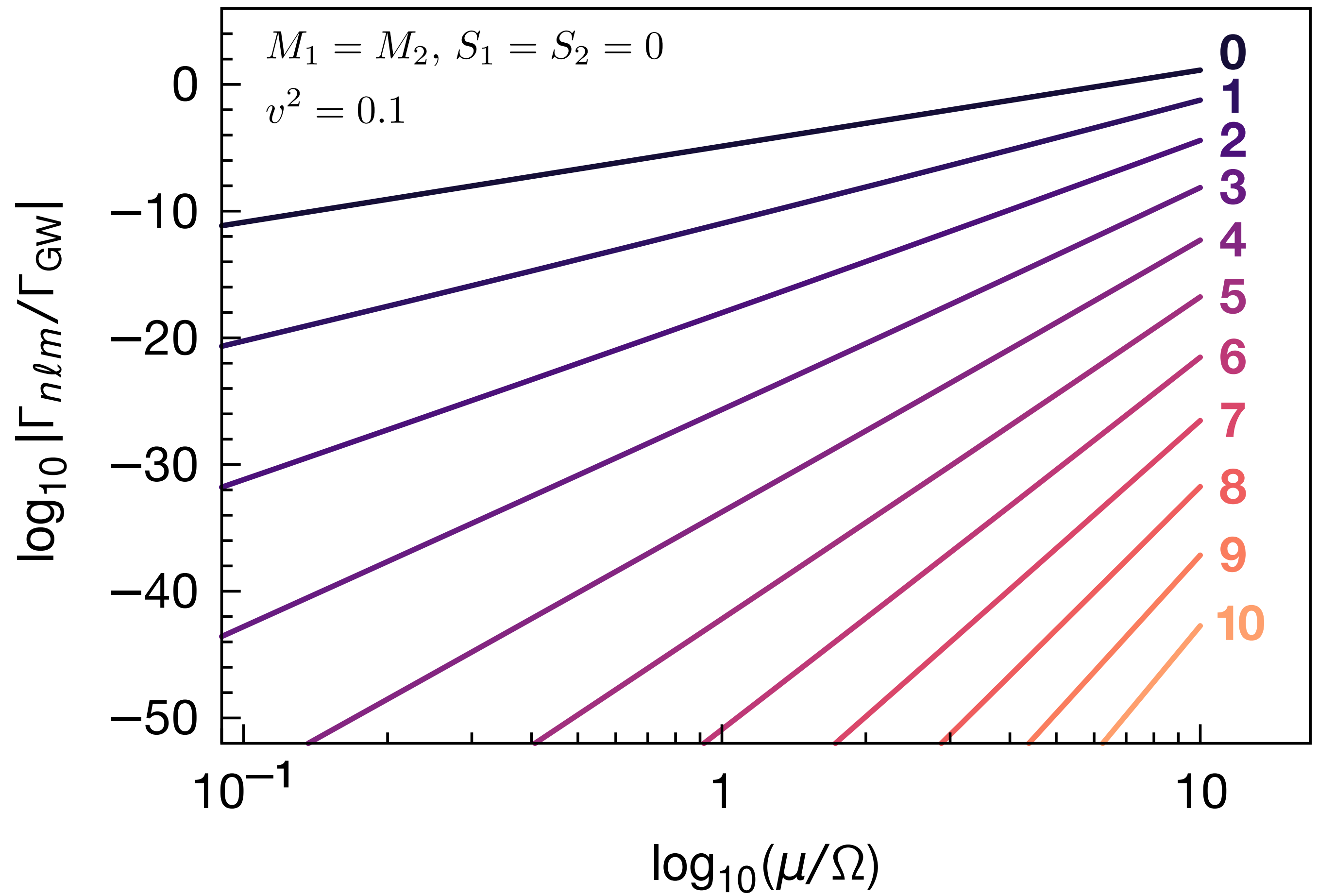
$$(n, \ell, m) = (\ell + 1, \ell, \ell)$$



GROWTH RATES

- ◆ Bound state grows exponentially like $\psi_{nlm} \sim e^{\Gamma_{nlm} t}$ when $0 < E_{nlm} < m\Omega$
- ◆ Energy is predominantly extracted from the orbital motion
- ◆ But timescale is not observationally relevant

$$(n, \ell, m) = (\ell + 1, \ell, -\ell)$$

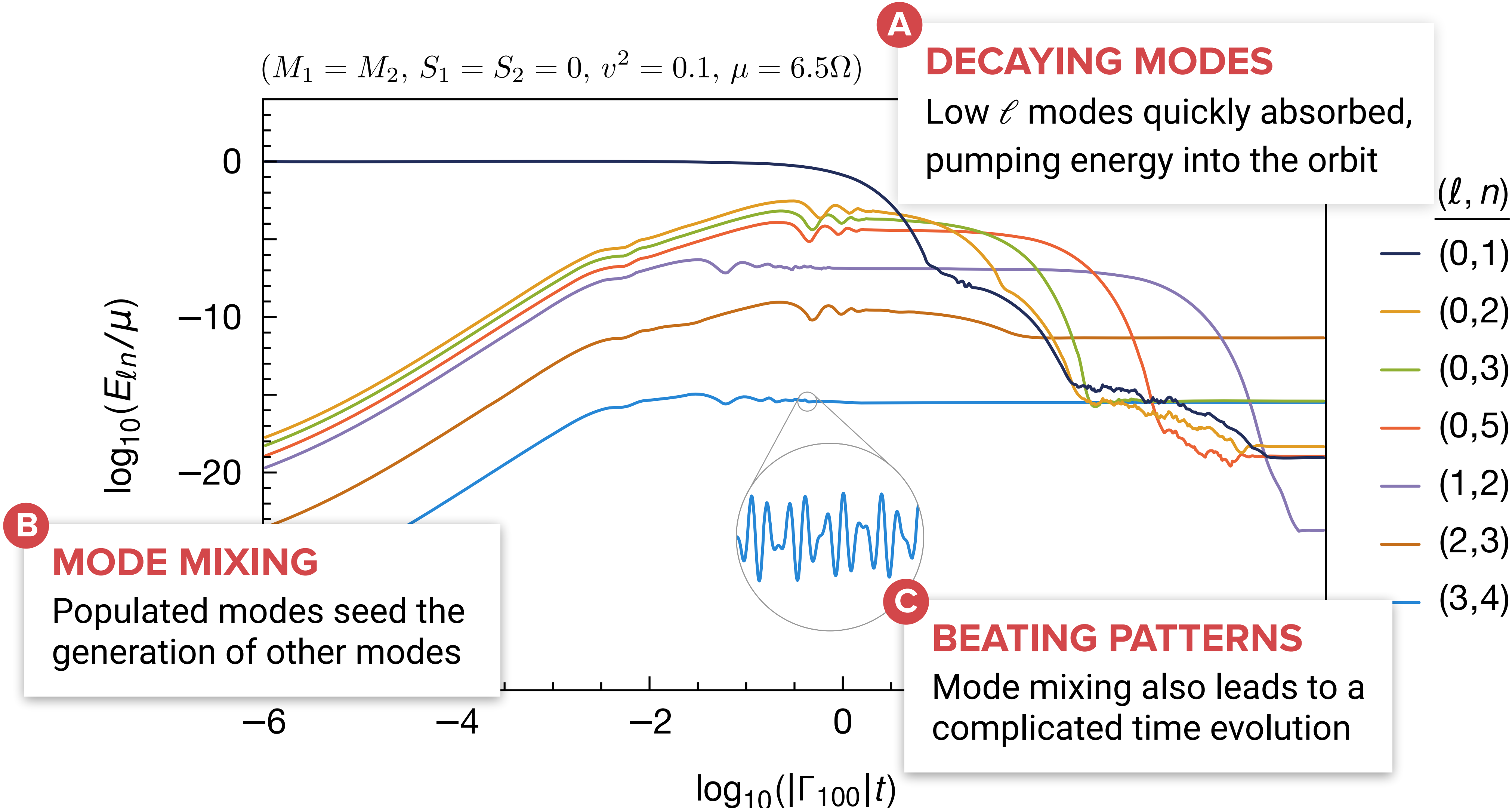


DECAY RATES

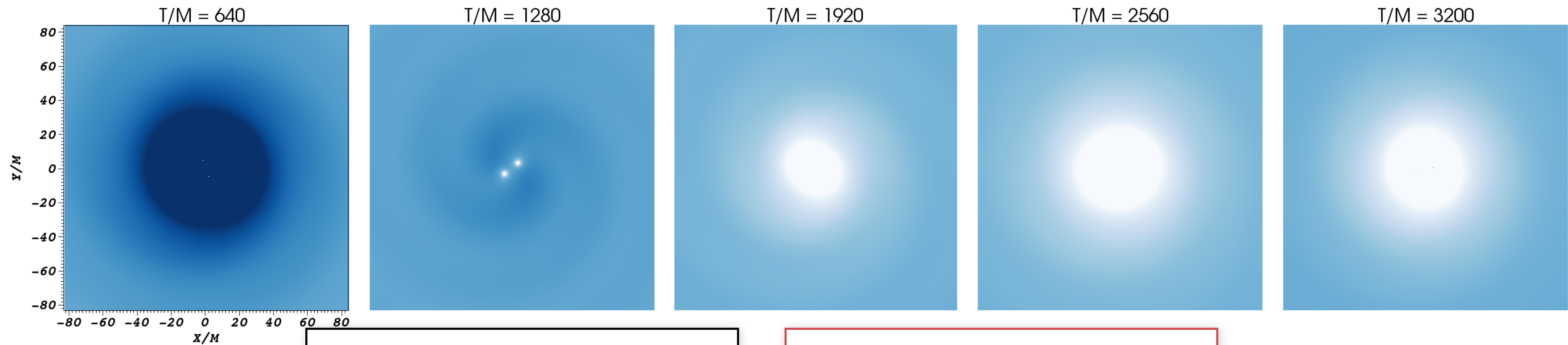
- ◆ Decaying modes pump energy into the orbit
- ◆ Upward trend suggests potential for observational signatures at larger values of μ

EVOLUTION OF THE CLOUD

Energy in the (ℓ, n) levels after a time t , with only the $(1,0,0)$ mode populated at $t=0$



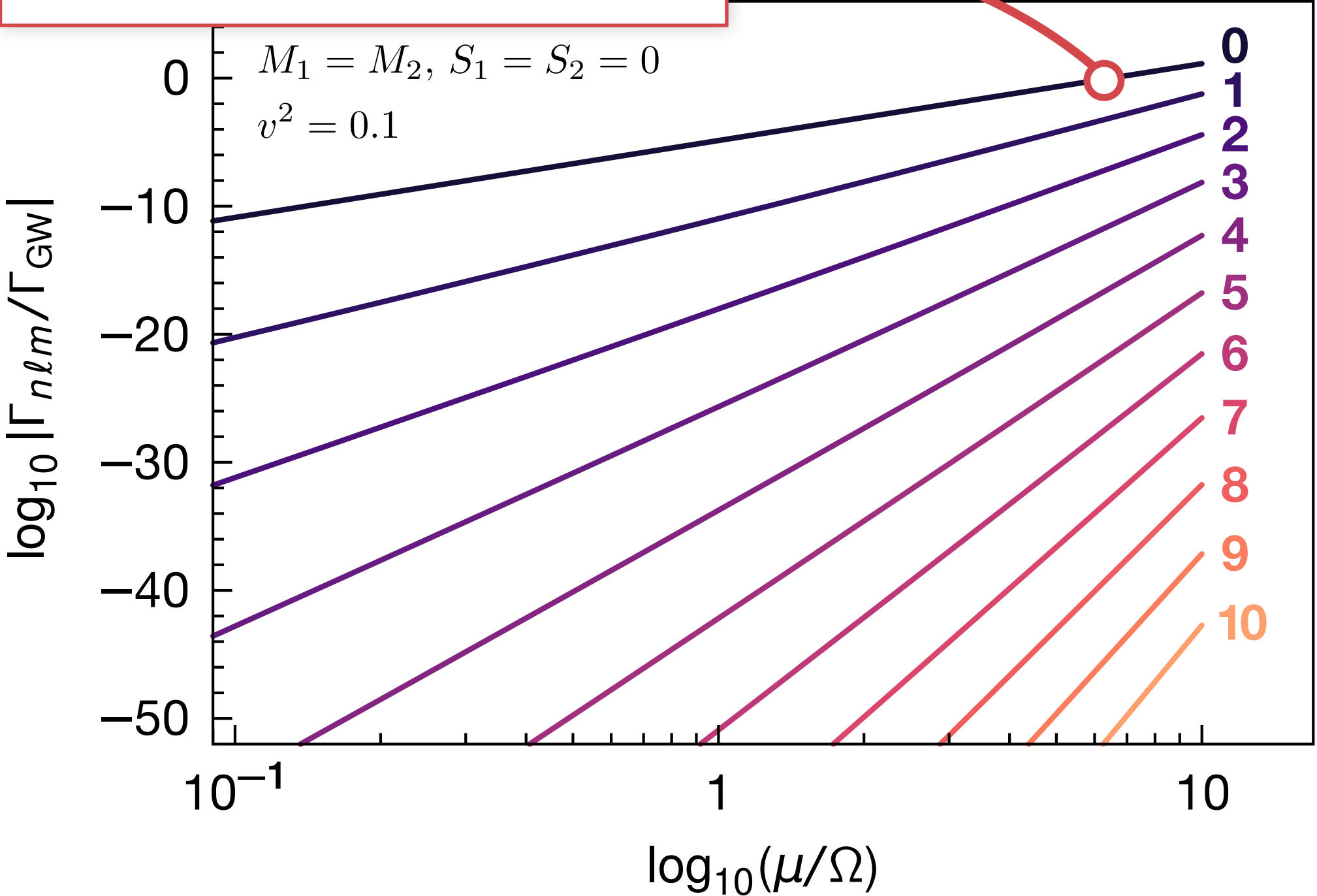
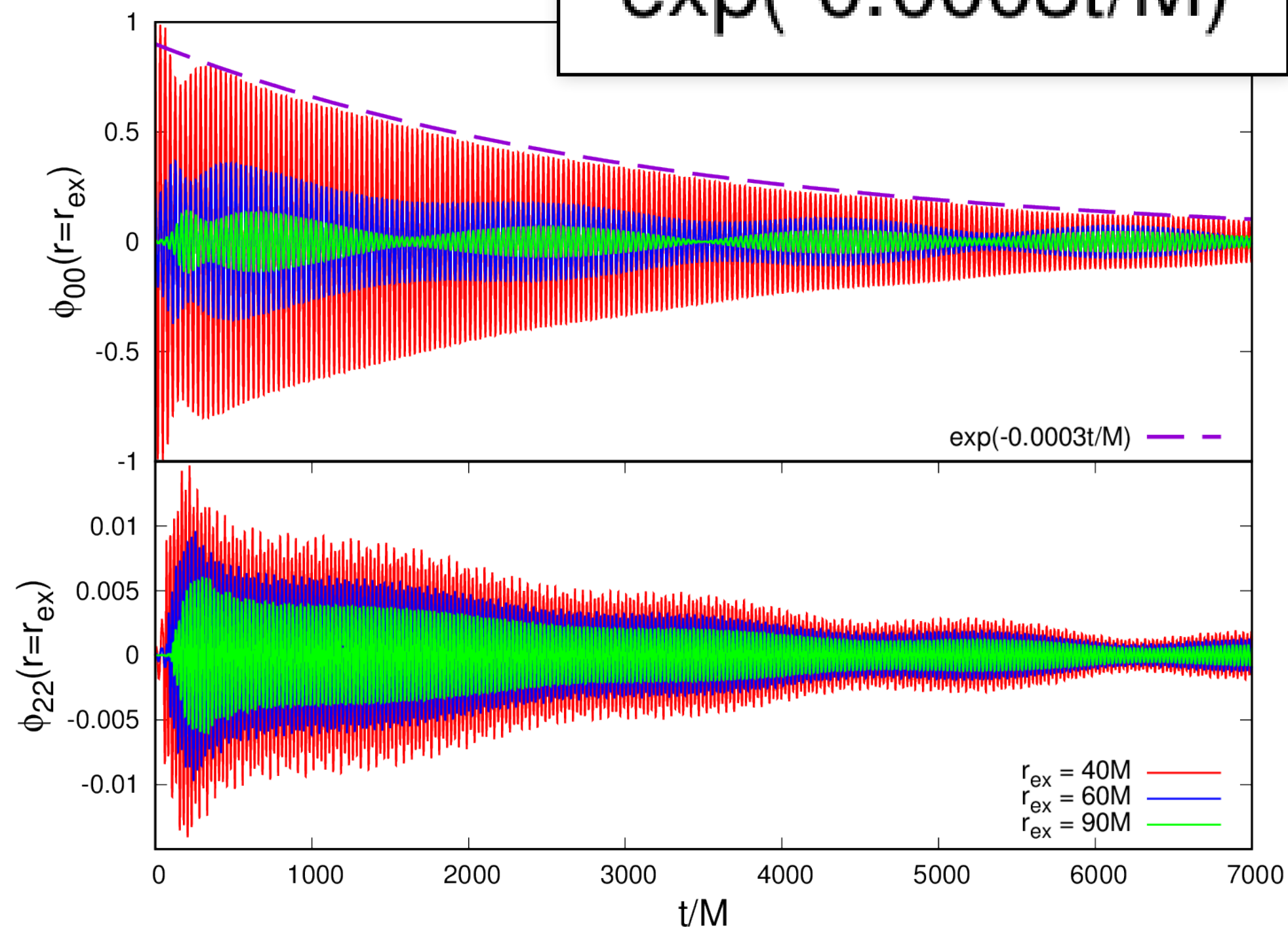
COMPARISON WITH NUMERICAL SIMULATIONS [Ikeda, Bernard, Cardoso, Zilhão 2020]

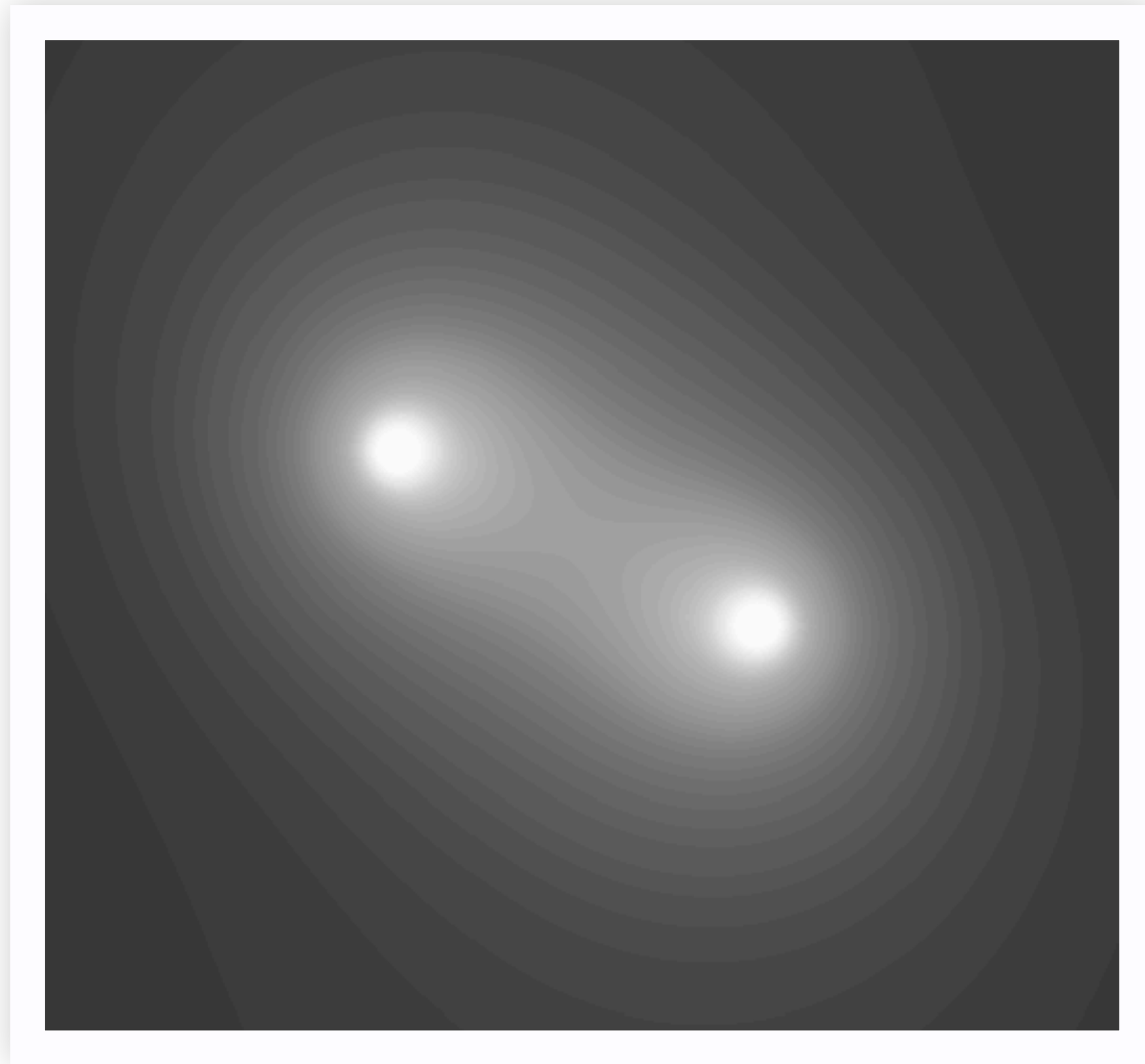


$$\exp(-0.0003t/M)$$

$$0.00026 (GM)^{-1}$$

$$(n, \ell, m) = (\ell + 1, \ell, -\ell)$$





[Ikeda, Bernard, Cardoso, Zilhão 2020]

OPEN QUESTIONS

- ◆ Does the cloud leave a measurable signal for larger μ ?
- ◆ Viable formation scenarios?

OUTLOOK

- ◆ EFT methods facilitate the systematic study of complicated systems
- ◆ Analytic results can help us interpret what we see in numerical simulations
- ◆ Similar techniques applicable to many other systems of astrophysical interest