Imperial College London Claudia de Rham

The Speed of Gravity



Meeting of the National Research Group on Gravitational Waves

Institut Poincaré







Thanks to some of amazing collaborators



Lasma Alberte (@ Imperial)



Calvin Chen (@ Imperial)



Jeremie Francfort (@ Geneva)



Sumer Jaitly (@ Imperial)



Aoibheann Margalit (@ Imperial)



Scott Melville (@ Cambridge)



Johannes Noller (@ Portsmouth)



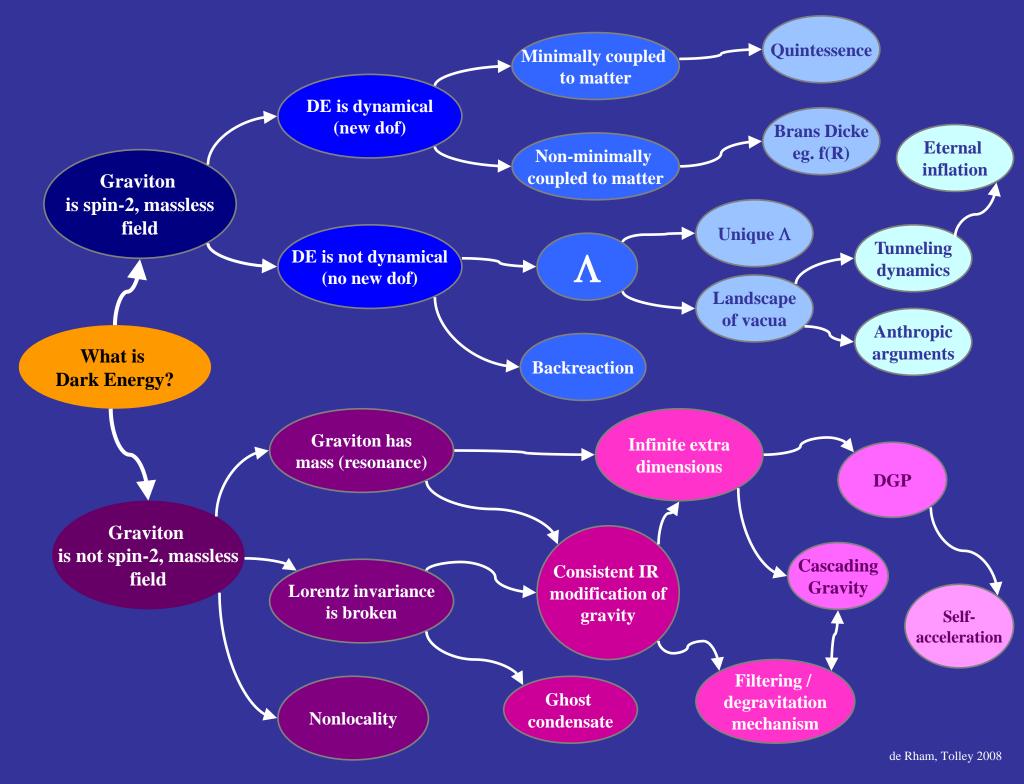
Andrew Tolley (@ Imperial)

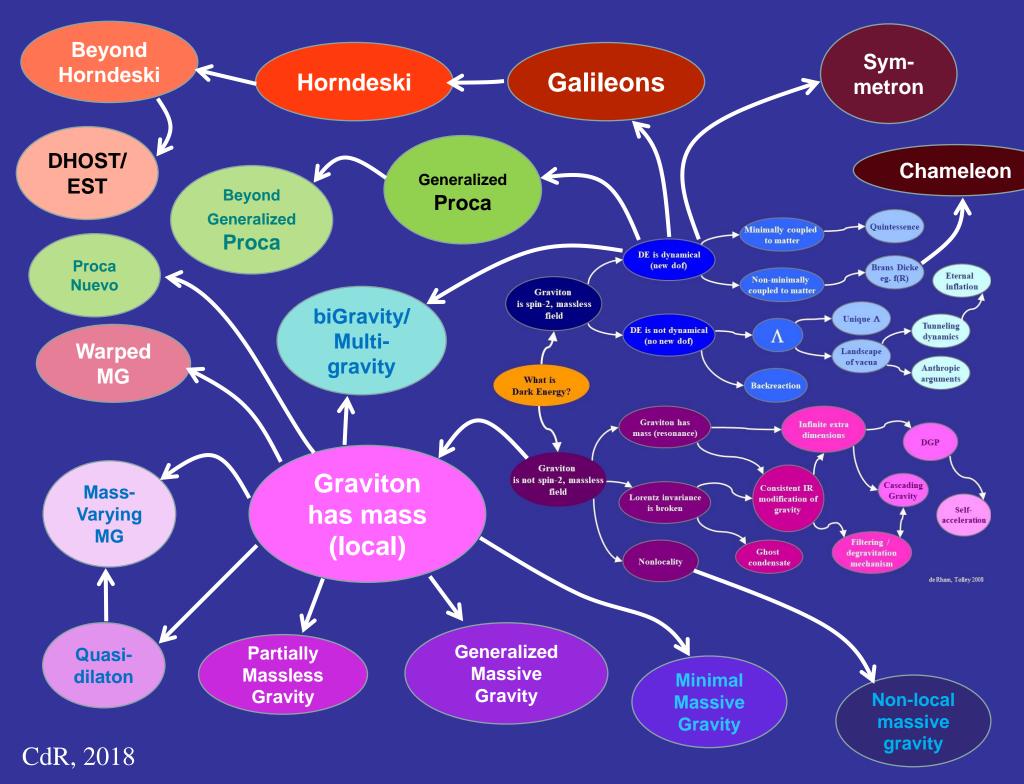


Jun Zhang (@ Imperial)



Shuang-Yong Zhou (@ USTC)





Diagnostic

• Does the model even make sense? (is it stable? Classically? QM??)



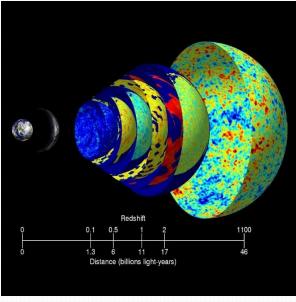
Diagnostic

O Low-energy criteria

- Does the model even make sense? (is it stable? Classically? QM??)
- Does it fit the multitude of existing observational constraints?
 (signature for future ones?)

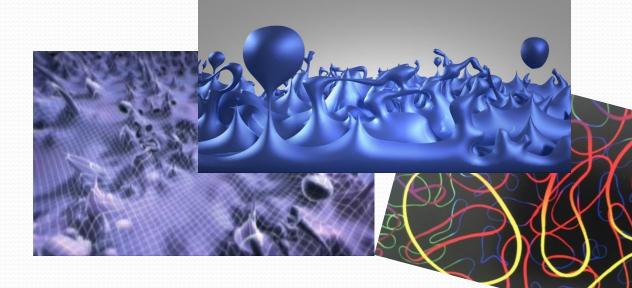
High energy completion?





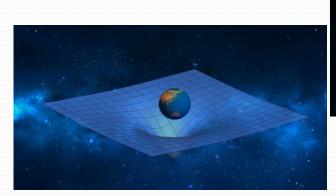


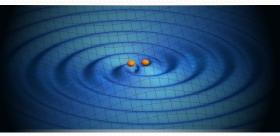
Energy



Will remain agnostic on the precise UV completion (field content, realization,...) so long as it is local, unitary, Lorentz invariant and CAUSAL









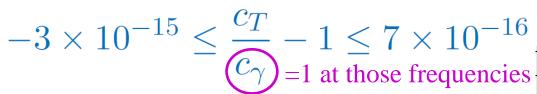
Within low-energy gravitational EFTs

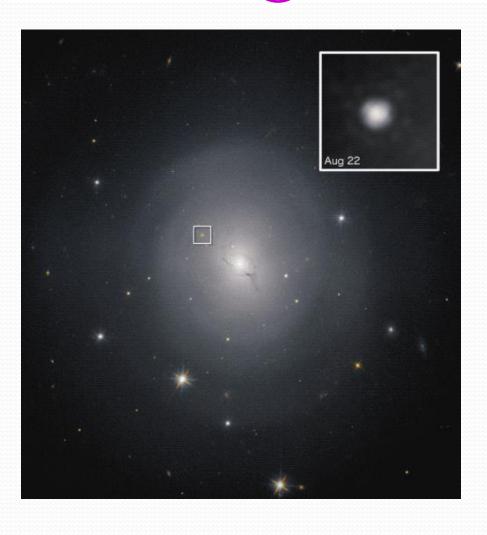
(relevant for EFT of Inflation, dark energy, dark matter, BSM,...)

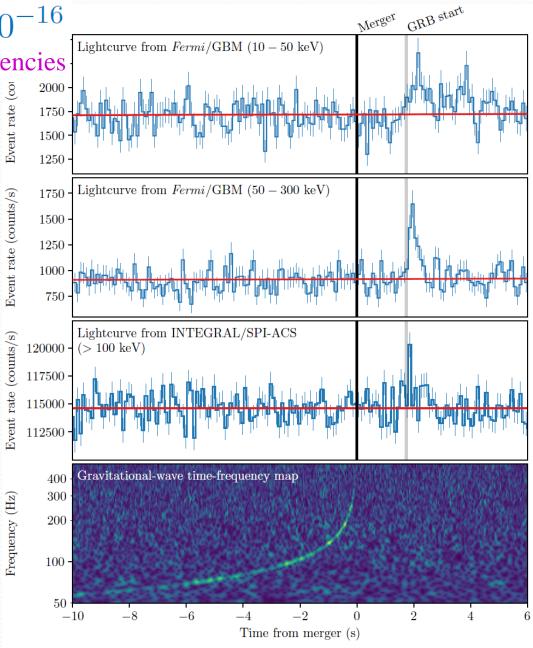
- Can Gravitational Waves be Superluminal?
- Can Other Species (eg light) be Superluminal?
- Is this consistent with Causality?
- Is this consistent with a Standard UV completion?



GW&GBR 170817

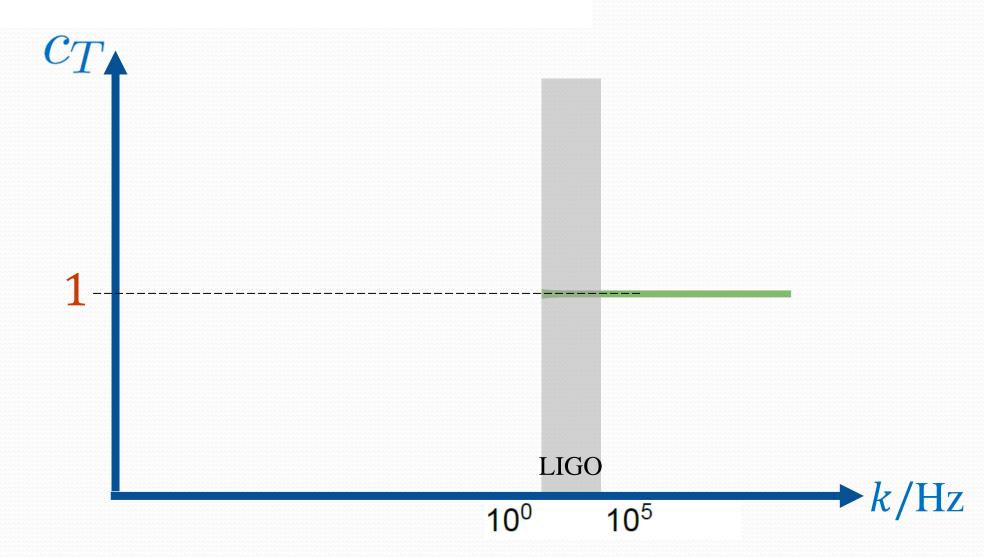




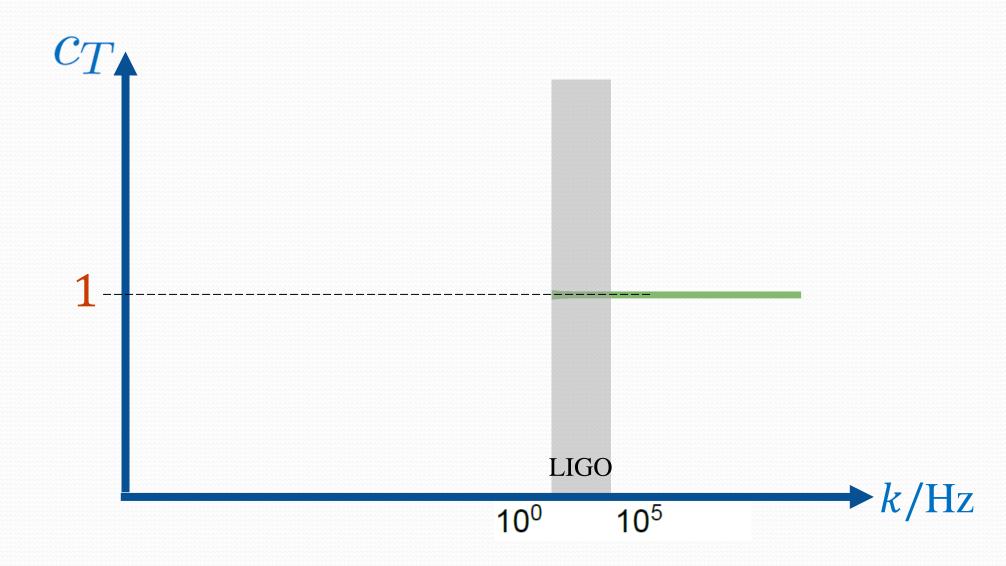


GW&GBR 170817

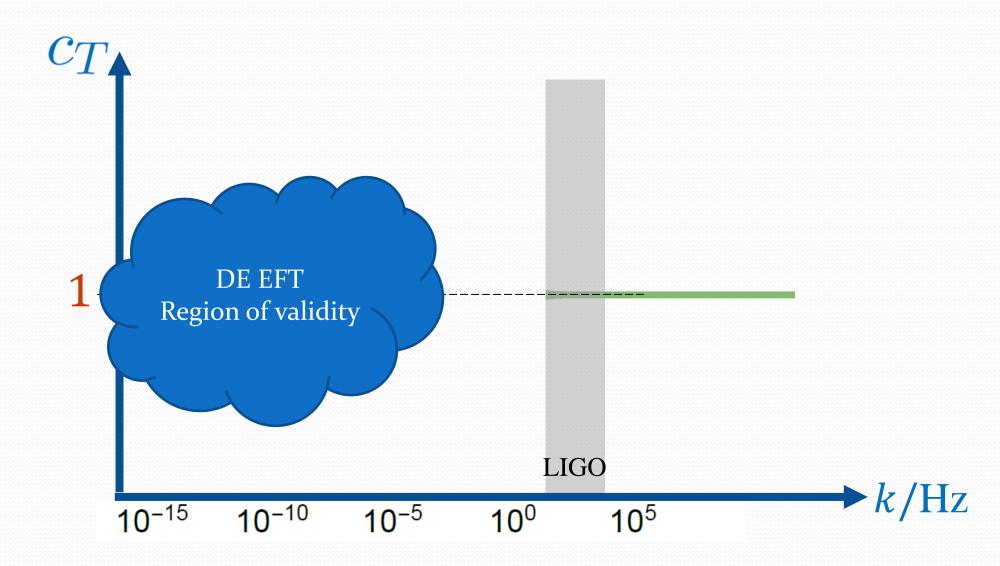
$$-\mathcal{O}\left(10^{-15}\right) < c_T - 1 < \mathcal{O}\left(10^{-16}\right)$$



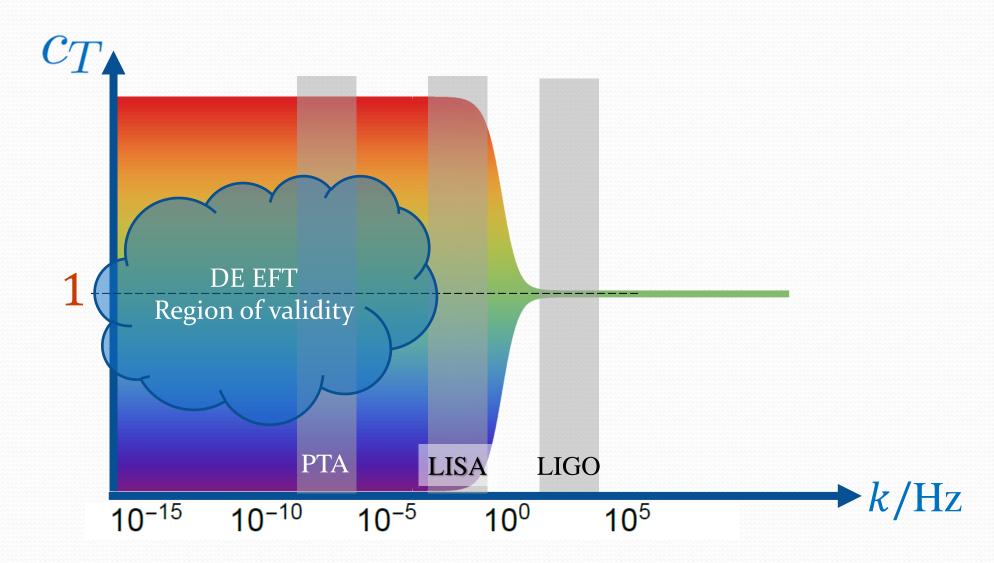
Many EFT for dark energy "predict" a non-luminal sound speed for GW on FLRW by an amount larger than 10⁻¹⁵



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From GW&GBR 170817

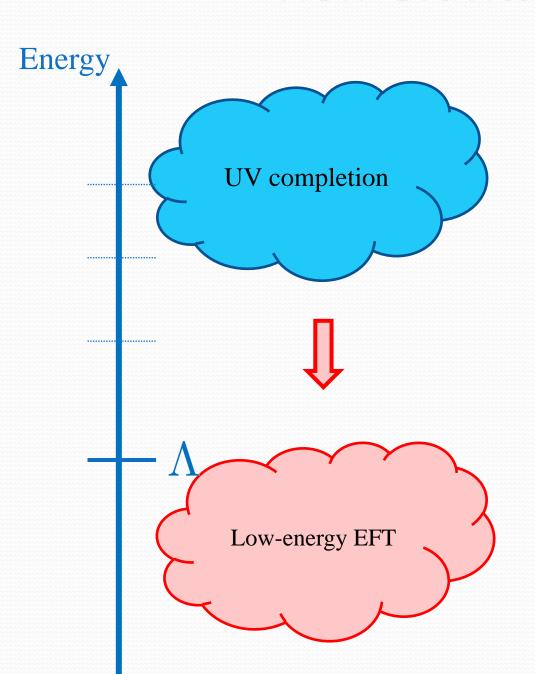
$$-\mathcal{O}(10^{-15}) < c_T - 1 < \mathcal{O}(10^{-16})$$

Is this even a possibility Or should we disregard this option???

$$-\mathcal{O}\left(10^{-15}\right) < c_T - 1 \le 0$$

type of prior often imposed

Non-Gravitational EFT



- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ CAUSAL (analyticity)
- ✓ Local (Froissart Bound)



positivity bounds

(applied to low-energy scattering amplitude or refractive index)

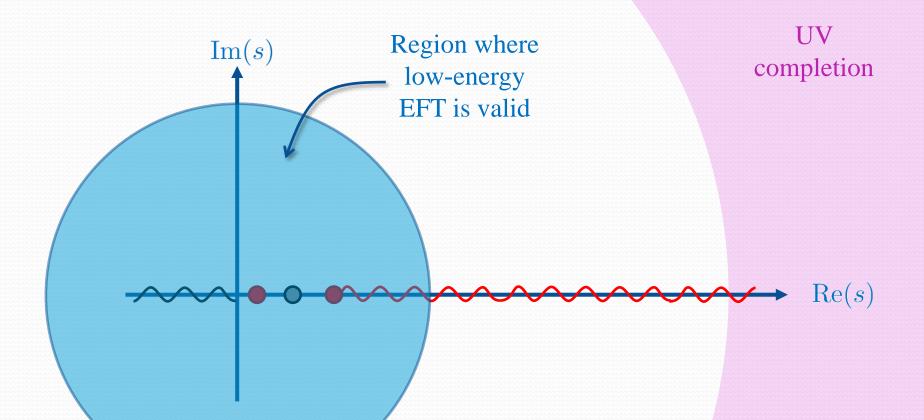


Low-energy EFT

 $\mathcal{A}(s)$: 2 – 2 elastic scattering amplitude

s: center of mass energy²

t: momentum transfer

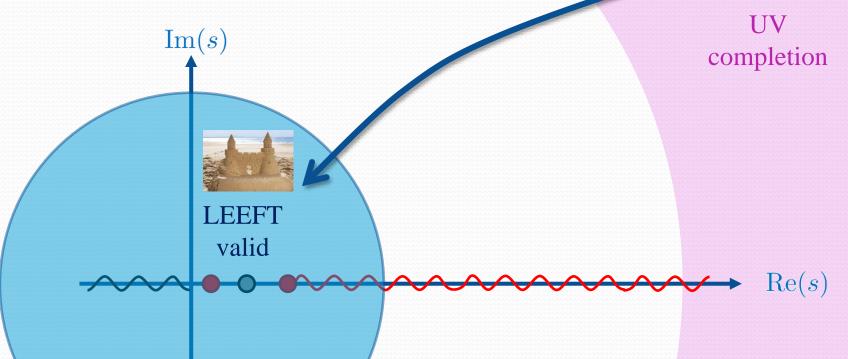


Analyticity - Causality





encoded by requirement of analyticity –
 is what connects UV to IR

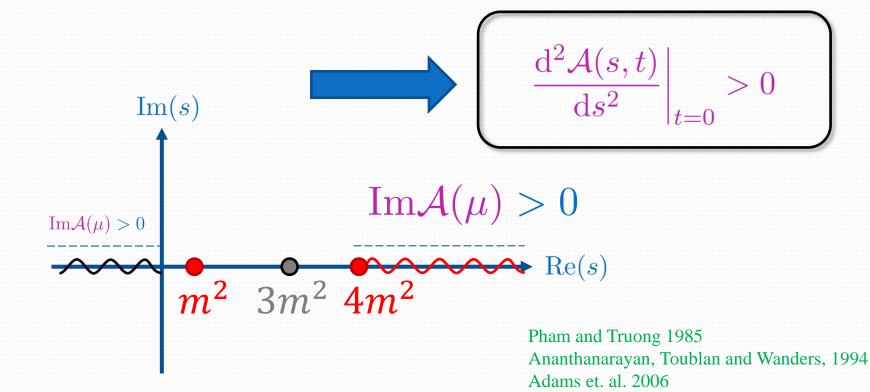


Non-Gravitational EFT

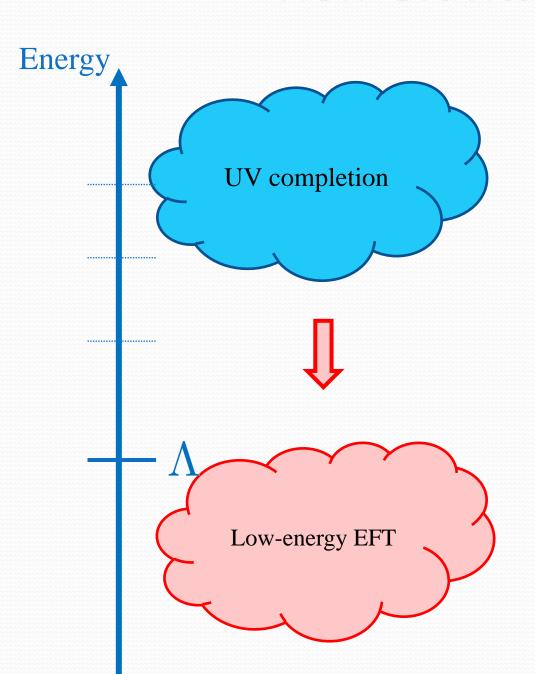
- UV completion
- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)

A: 2-2 elastic scattering amplitude

$$2\operatorname{Im} = \sum_{X} \left| \sum_{X} \right|^{2} \ge \left| \sum_{X} \right|^{2}$$



Non-Gravitational EFT



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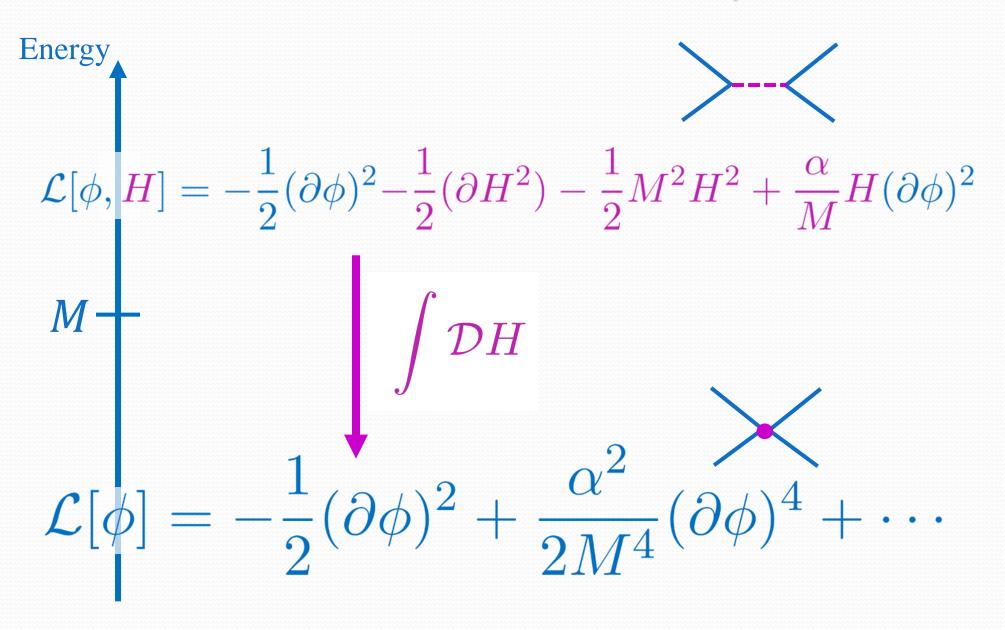


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(applied to low-energy scattering amplitude or refractive index)



EFT for a Scalar Field ϕ



EFT for a Scalar Field ϕ

$$\mathcal{L}[\phi] = -\frac{1}{2}(\partial\phi)^2 + \frac{\alpha^2}{2M^4}(\partial\phi)^4 + \cdots$$

On any Lorentz symmetric background, the sound speed is exactly luminal

$$\phi = f(\eta_{\mu\nu}x^{\mu}x^{\nu}) + \chi \qquad \Rightarrow \qquad c_{\chi} = 1$$

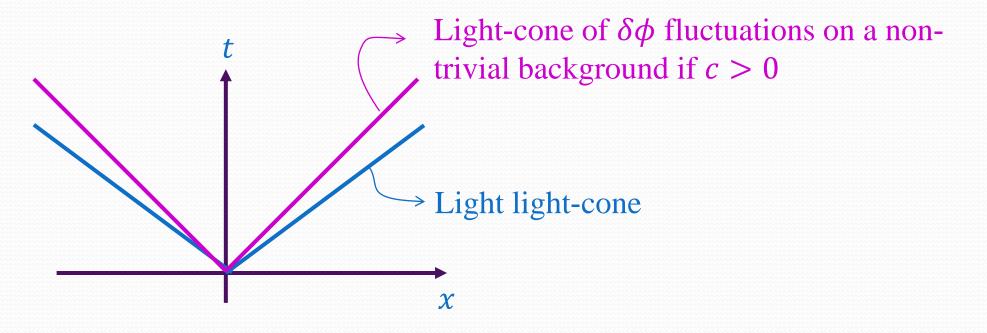
On backgrounds that spontaneously break Lorentz Invariance, the sound speed can be subluminal

$$\phi = \phi(t) + \chi \qquad \Rightarrow \qquad c_{\chi} = 1 - \frac{\alpha^2}{M^4} \dot{\phi}^2$$

Scalar Field minimally coupled to gravity

No gravity:
$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 + \frac{c}{\Lambda^4}(\partial\phi)^4 + \cdots$$

In the absence of gravity, positivity bounds assuming Unitarity, Analyticity, Causality require c > 0



Non-Gravitational EFT

UV completion

- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ CAUSAL (analyticity)
- ✓ Local (Froissart Bound)

Eg of LEEFT

$$\mathcal{L}[\phi] = -\frac{1}{2}(\partial\phi)^2 + \frac{\alpha^2}{2M^4}(\partial\phi)^4 + \cdots$$

$$\mathcal{A}''(s)\big|_{t=0} \sim \alpha^2 > 0$$



(improved) positivity bounds



$$\phi = \phi(t) + \chi$$
 \Rightarrow $c_{\chi} = 1 - \frac{\alpha^2}{M^4} \dot{\phi}^2$

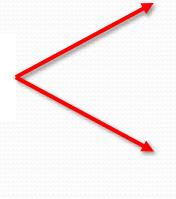
UV completion

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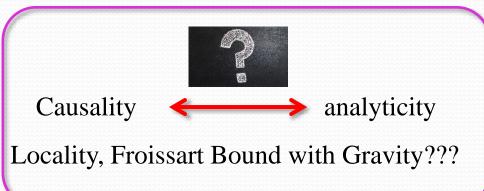


(improved) positivity bounds

More subtle for gravitational EFTs







- ✓ Unitary (optical theorem)
- ✓ Lorentz invariant (crossing symmetry)
- ✓ CAUSAL (analyticity)







(improved) positivity bounds

More subtle for gravitational EFTs







(sub)luminal sound speed

Connection more subtle with gravity for 2 reasons:

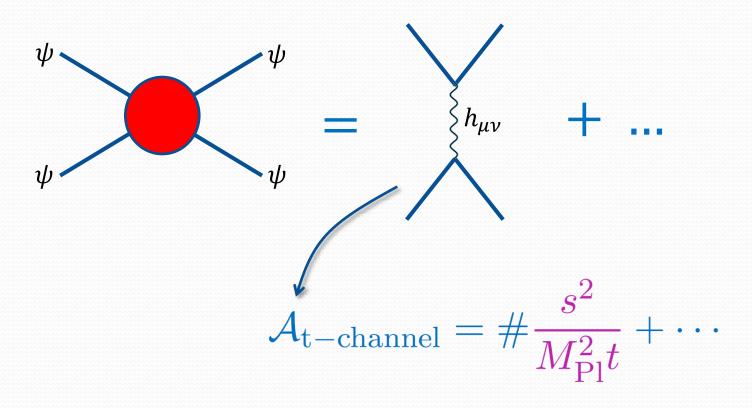
1. Gravitational exchange

cannot prevent gravity from coupling to everyone and having an effect

2. Frame artefacts

(speed not invariant under frame transformations, notion of causality is) trivial artefact yet important implications

Positivity Bounds in Gravitational LEEFT



t-channel pole from gravity exchange compromises positivity bound
$$\frac{\mathrm{d}^2 \mathcal{A}(s,t)}{\mathrm{d}s^2} \bigg|_{t=0} > 0$$

Causality in a Negative World

In non-gravitational EFTs

causality

positivity bounds

$$\mathcal{A} \sim \mathcal{C} \, s^2 + \dots \Rightarrow \mathcal{C} > 0$$

 \Rightarrow

(sub)luminal sound speed $0 < c_s \le 1$

In gravitational EFTs at scale *M*

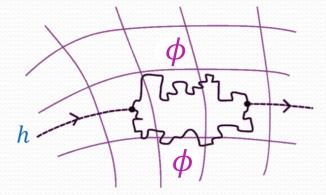
causality

Approximate positivity bounds $C > -\frac{\mathcal{O}(1)}{M^2 M^2}$

sound speed

$$C > -\frac{\mathcal{O}(1)}{M^2 M_{\rm Pl}^2} \implies 0 < c_s \le 1 + \mathcal{O}\left(\frac{M^2}{M_{\rm Pl}^2}\right)$$

EFT for Gravity



Energy

High-energy theory with gravity and light & heavy modes

 $\int \mathcal{D}H$

Integrate out heavy modes

Low-energy EFT of gravity

$$\mathcal{L}_{\mathrm{IR}} = \sqrt{-g} \left[-\Lambda^{\mathrm{IR}} + \frac{M_{\mathrm{Pl}}^2}{2} R + \mathcal{L}_{\psi}^{(\mathrm{light})}(g, \psi) + \mathcal{L}_{R^2} + \frac{1}{M^2} \mathcal{L}_{R^3} + \cdots \right]$$

Speed of Gravity

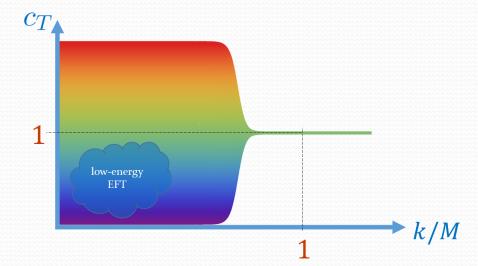
$$\mathcal{L}_{IR} = \sqrt{-g} \left[-\Lambda^{IR} + \frac{M_{Pl}^2}{2} R + \mathcal{L}_{\psi}^{(light)}(g, \psi) + \mathcal{L}_{R^2} + \frac{1}{M^2} \mathcal{L}_{R^3} + \cdots \right]$$

For GWs on curved background (e.g. FLRW, Schwarzschild,...) even in a frame where high-frequency travel luminally, at low-frequency speed may be qualified as *super*luminal

$$c_s^2 = 1 + \mathcal{O}(\pm 1) \frac{(-\dot{H})}{M_{\rm Pl}^2} + \cdots$$

 $c_s^2 = 1 + \mathcal{O}(\mp 1) \frac{1}{M^2 M_{\rm Pl}^2 r_g^4} + \cdots$

+ corrections at high energy



Speed of Gravity

$$\mathcal{L}_{IR} = \sqrt{-g} \left[-\Lambda^{IR} + \frac{M_{\rm Pl}^2}{2} R + \mathcal{L}_{\psi}^{({\rm light})}(g, \psi) + \mathcal{L}_{R^2} + \frac{1}{M^2} \mathcal{L}_{R^3} + \cdots \right]$$

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$$c_s^2 = 1 + \mathcal{O}(\pm 1) \frac{(-\dot{H})}{M_{\rm Pl}^2} + \cdots + \text{corrections at high energy}$$
 $c_s^2 = 1 + \mathcal{O}(\mp 1) \frac{1}{M^2 M_{\rm Pl}^2 r_g^4} + \cdots$

Typical approach (eg. for EFT of DE) is to constrain low-energy EFT so as to ensure SUB luminality!

Speed of Gravity

$$\mathcal{L}_{IR} = \sqrt{-g} \left[-\Lambda^{IR} + \frac{M_{\rm Pl}^2}{2} R + \mathcal{L}_{\psi}^{({\rm light})}(g, \psi) + \mathcal{L}_{R^2} + \frac{1}{M^2} \mathcal{L}_{R^3} + \cdots \right]$$

For GWs on curved background (e.g. FLRW, Schwarzschild,...) even in a frame where high-frequency travel luminally, at low-frequency speed may be qualified as *super*luminal

$$c_s^2 = 1 + \mathcal{O}(\pm 1) \frac{(-\dot{H})}{M_{\rm Pl}^2} + \cdots + \text{corrections at high energy}$$
 $c_s^2 = 1 + \mathcal{O}(\mp 1) \frac{1}{M^2 M_{\rm Pl}^2 r_q^4} + \cdots$

Such levels of Superluminality are not in tension with causality and should not be used as a way to discard some operators in the low-energy EFT of gravity

CAUSALITY



(sub)luminal sound speed

Connection more subtle with gravity for 2 reasons:

1. Gravitational exchange cannot prevent gravity from coupling to everyone and having an effect

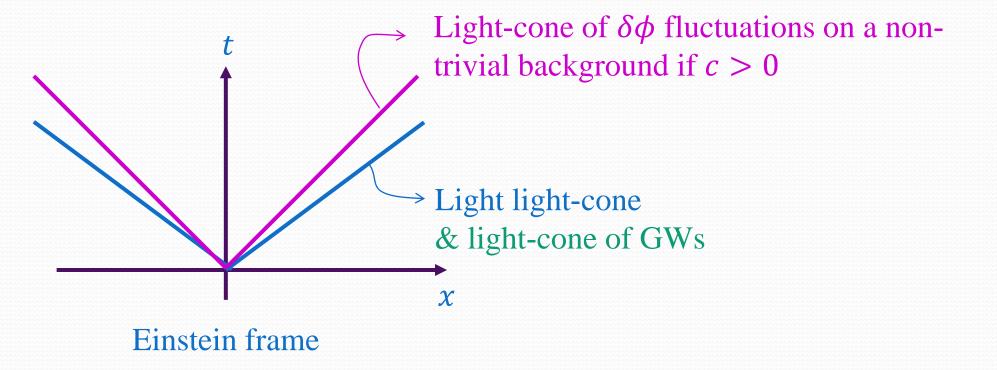
2. Frame artefacts

(speed not invariant under frame transformations, notion of causality is) trivial artefact yet important implications

Scalar Field minimally coupled to gravity

Minimally-coupled to gravity:

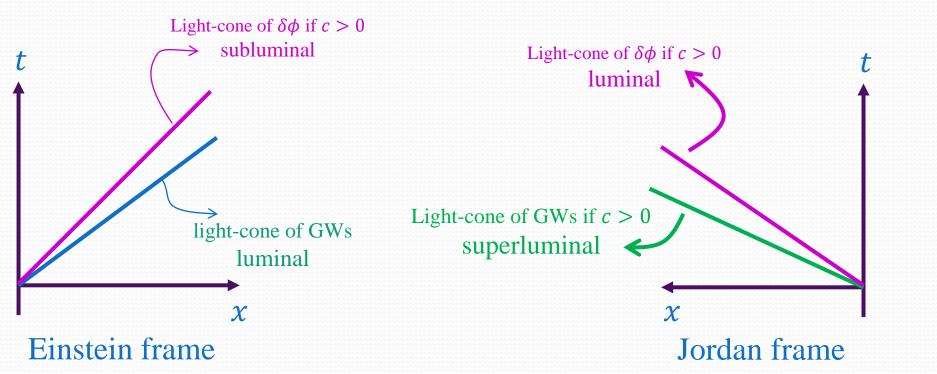
$$\mathcal{L} = \sqrt{-g} \left| \frac{M_{\text{Pl}}^2}{2} R - \frac{1}{2} (\partial \phi)^2 + \frac{c}{\Lambda^4} (\partial \phi)^4 + \cdots \right|$$



Scalar Field minimally coupled to gravity

Einstein Frame:
$$\mathcal{L} = \sqrt{-g} \left[\frac{M_{\rm Pl}^2}{2} R - \frac{1}{2} (\partial \phi)^2 + \frac{c}{\Lambda^4} (\partial \phi)^4 + \cdots \right]$$

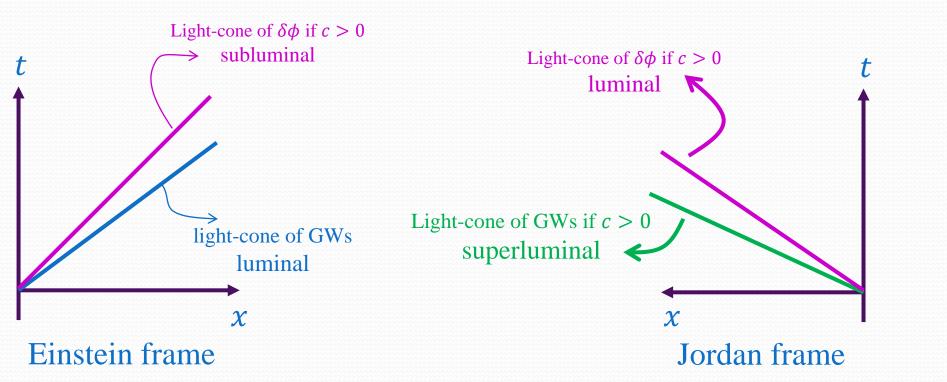
$$g_{\mu\nu} = \tilde{g}_{\mu\nu} - \frac{c}{\Lambda^4} \partial_{\mu} \phi \partial_{\nu} \phi + \frac{c M_{\rm Pl}^2}{\Lambda^4} G_{\mu\nu}$$
 Jordan Frame:
$$\mathbf{L} = \sqrt{-\tilde{g}} \left[\frac{M_{\rm Pl}^2}{2} \tilde{R} - \frac{1}{2} (\tilde{\partial} \phi)^2 + \tilde{c} \ \tilde{G}^{\mu\nu} \tilde{R}_{\mu\nu} + \cdots \right]$$



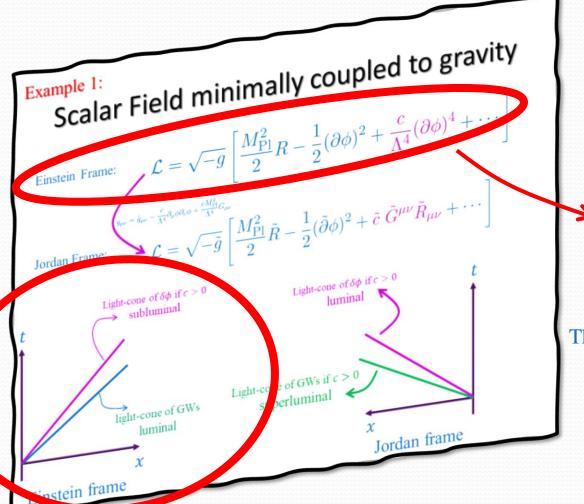
Frame artefacts

In some low-energy EFTs, causality imposes superluminal GWs... just a trivial (yet Important!) frame artefact

particularly important when dealing with potential modifications of Gravity (e.g. à la Horndeski,...) but not only



In some low-energy EFTs, causality imposes superluminal GWs... just a trivial (yet Important!) frame artefact

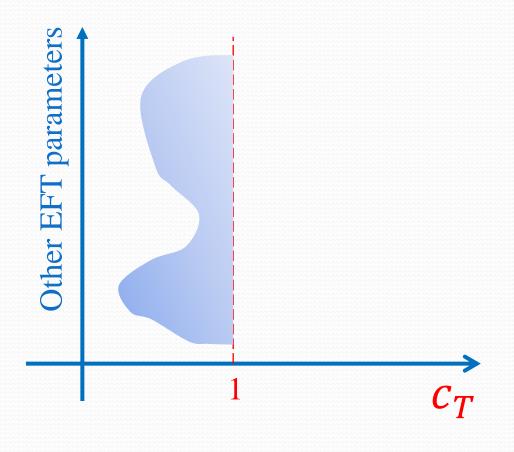


To 0th order, it is safe to impose subluminality for all the fields once we are in the frame where gravity can be decoupled

Frame where we can take a smooth limit $M_{Pl} \rightarrow \infty$

The change of frame is singular in that limit

In some low-energy EFTs, causality imposes superluminal GWs... just a trivial (yet Important!) frame artefact

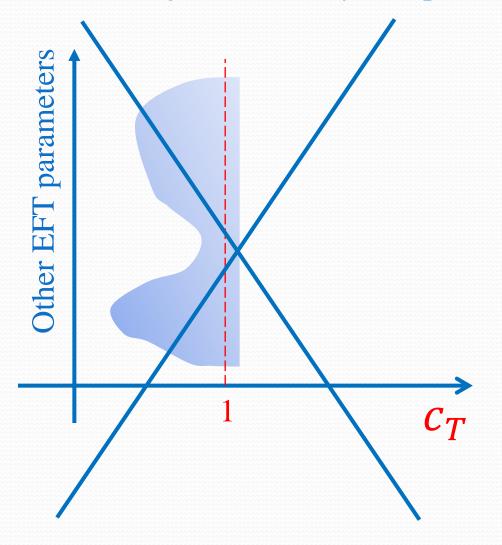


To 0th order, it is safe to impose subluminality for all the fields once we are in the frame where gravity can and is decoupled

Frame where we can take a smooth limit $M_{Pl} \rightarrow \infty$

This may imply a large amount of superluminality in original frame

In some low-energy EFTs, causality imposes superluminal GWs... just a trivial (yet Important!) frame artefact

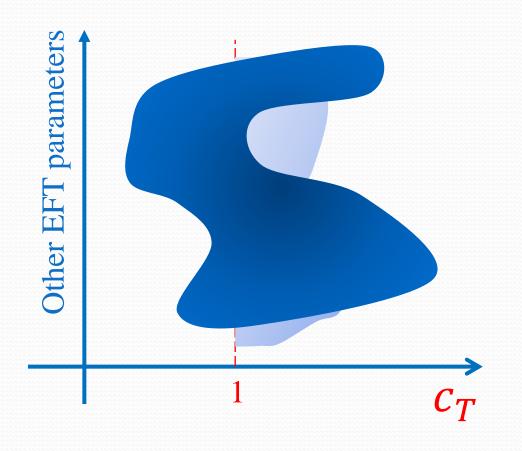


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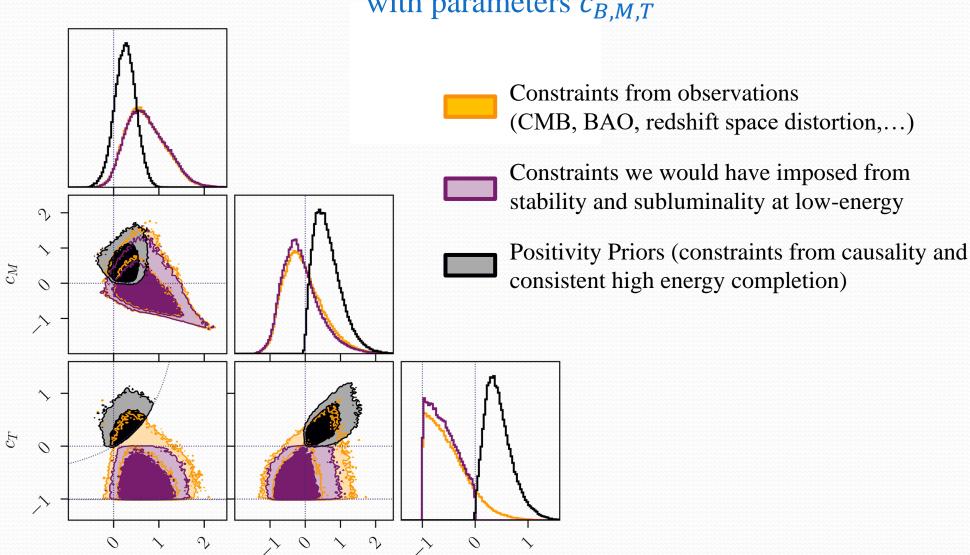
Frame where we can take a smooth limit $M_{Pl} \rightarrow \infty$

This may imply a large amount of superluminality in original frame

Positivity bounds are frame independent – agnostic to such considerations In doubts, apply (approximate) positivity bounds (with allowed negativity from t-channel pole)

Constraints on Low-energy Models

Example of Dark Energy model (quartic Horndeski) with parameters $c_{B,M,T}$



 c_T

 c_B

 c_M

Living with Superluminality

- Gravitational Waves are luminal to a (VERY) good accuracy at LIGO frequencies $-\mathcal{O}\left(10^{-15}\right) < c_T 1 < \mathcal{O}\left(10^{-16}\right)$
- Within the standard EFT of gravity, GWs are no longer perfectly luminal on backgrounds that spontaneously break Lorentz invariance (eg Schwarzschild, FLRW, the real world,...)

Lesson 1:

- In an arbitrary frame, GWs may be superluminal
- Imposing subluminality priors only makes sense in a frame where gravity can be decoupled
- In doubts, to be derived from positivity bounds
- In the original frame this may correspond to GWs being superluminal by a 'considerable' amount

Living with Superluminality

Lesson 2:

- Even in the frame where matter and gravity can decouple,
 a tiny amount of SL or a negative phase shift be it for GWs or
 other fields is not in conflict with causality.
 It may even follow from consistent causal and Lorentz invariant
 UV completions.
- In the frame where matter and gravity can decouple, superluminality is consistent with causality so long as

$$\lim_{M_{
m Pl} o \infty} |c_s^2 - 1| \sim M_{
m Pl}^{-lpha} \quad ext{with} \quad lpha \geq 2$$

Living with Superluminality

Lesson 2:

- Even in the frame where matter and gravity can decouple,
 a tiny amount of SL or a negative phase shift be it for GWs or
 other fields is not in conflict with causality.
 It may even follow from consistent causal and Lorentz invariant
 UV completions.
- superluminality not in conflict with causality so long as amplitude respects some (approximate) positivity bounds

$$\mathcal{A} \sim -\frac{s^2}{M_{\rm Pl}^2 t} + \frac{c \, s^2}{M^4}$$
 with $c > -\frac{M^2}{M_{\rm Pl}^2} \times \mathcal{O}(1)$