# Random loops and $T$-algebras 

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IHÉS (virtual), 16 November 2020

## Stochastic quantisation

Basic idea: Consider discrete approximation to "Euclidean QFT" $e^{-\beta S(\varphi)} D \varphi$ so $\varphi$ belongs to finite-dimensional vector space. This is invariant for stochastic evolution

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$1 D \sigma$-model: Field configurations given by loops on Riemannian manifold: $u: S^{1} \rightarrow \mathcal{M}, S(u)=\int_{S^{1}} g_{u}\left(\partial_{x} u, \partial_{x} u\right) d x$, usual Dirichlet energy.

## Formal Gibbs measure

Brownian loop measure on manifold $(\mathcal{M}, g)$ formally given (for some $c$ ) by

$$
\begin{gathered}
\mathbf{P}(D u) \propto \exp \left(-\int_{S^{1}}\left(\frac{1}{2} g_{u}\left(\partial_{x} u, \partial_{x} u\right)-c R(u)\right) d x\right) \text { " } D u " . \\
\text { Scalar curvature }
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In local coordinates

$$
\left(\partial_{t}-\partial_{x}^{2}\right) u^{\alpha}=\Gamma_{\beta \gamma}^{\alpha}(u) \partial_{x} u^{\beta} \partial_{x} u^{\gamma}+c g^{\alpha \beta}(u) \partial_{\beta} R(u)+\sqrt{2} \sigma_{i}^{\alpha}(u) \xi_{i}
$$

with $\sigma_{i}^{\alpha} \sigma_{i}^{\beta}=g^{\alpha \beta}, \Gamma$ Christoffel symbols for Levi-Civita.

## A general result

Given $H \in \mathcal{C}^{\infty}\left(\mathbf{R}^{d}, \mathbf{R}^{d}\right)$, write $U^{\varepsilon}(\Gamma, \sigma, H)$ for some (formal) $\varepsilon$-approximation to

$$
\left(\partial_{t}-\partial_{x}^{2}\right) u^{\alpha}=\Gamma_{\beta \gamma}^{\alpha}(u) \partial_{x} u^{\beta} \partial_{x} u^{\gamma}+H^{\alpha}(u)+\sigma_{i}^{\alpha}(u) \xi_{i}
$$

RST yields a collection $\mathcal{S}=\{\Omega, \infty, \&, \& \in\},, \& \%, \ldots\}$ of 54 symbols and a valuation map $\Upsilon_{\Gamma, \sigma}: \mathcal{S} \rightarrow \mathcal{C}^{\infty}\left(\mathbf{R}^{d}, \mathbf{R}^{d}\right)$ s.t.:

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Comodule structure for two Hopf algebras, allowing to recenter in probability space (renormalisation) and in real space (Taylor-like expansions).

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Theorem A (H., Bruned, Chandra, Chevyrev, Zambotti): For every choice of $\Gamma, \sigma, H$ and every truncation of heat kernel there exist constants $C_{\varepsilon}^{\text {InP1I }} \in\langle\mathcal{S}\rangle$ such that

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U(\Gamma, \sigma, H)=\lim _{\varepsilon \rightarrow 0} U^{\varepsilon}\left(\Gamma, \sigma, H+\Upsilon_{\Gamma, \sigma} C_{\varepsilon}^{\text {Beprit }}\right) .
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## Preservation of symmetries

Metatheorem: If, for some approximation procedure, $U^{\varepsilon}$ satisfies a symmetry, then one can find constants $C_{\varepsilon}$ such that

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U^{s \mathrm{sm}}(\Gamma, \sigma, H)=\lim _{\varepsilon \rightarrow 0} U^{\varepsilon}\left(\Gamma, \sigma, H+\Upsilon_{\Gamma, \sigma} C_{\varepsilon}\right)
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also satisfies the symmetry in question. (Also $C_{\varepsilon}-C_{\varepsilon}^{\text {iprly }} \rightarrow$ const.)

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1. Yields equivariant ('Stratonovich') solution theories $U^{\text {seo }}$ parametrised by a 15-dimensional affine subspace $\mathcal{S}^{860}$ of vector fields.
2. Yields ('Itô') solution theories $U^{\text {thồ }}$ satisfying Itô isometry (law depends only on $\sigma_{i}^{\alpha} \sigma_{i}^{\beta}=g^{\alpha \beta}$ ) parametrised by a 19 -dimensional affine subspace $\mathcal{S}^{\text {rio }}$.

## Itô = Stratonovich!

Theorem (Bruned, Gabriel, H., Zambotti): There exists a two-parameter family of solution theories $U$ satisfying both symmetries simultaneously. All of them coincide with existing notions of solution in all previously studied cases.

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Recall solution given by

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Expect $\left(\Upsilon_{\Gamma, \sigma} C_{\varepsilon}\right)(u)=0$ whenever $\Gamma(u)=0$ and $(\partial \sigma)(u)=0$ (pointwise).

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Theorem: There exists a one-parameter family of solution theories $U$ satisfying 'equivariance / Stratonovich', 'Itô isometry', and 'minimality'.

## Back to geometry

In geometric case when $\Gamma$ are Christoffel symbols for Levi-Civita, all elements in that one-parameter family coincide! Yields completely natural notion of solution.

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Explains previously observed fact that different approximations to Brownian bridge measure are of form

$$
\exp \left(-\frac{1}{2} \int\left(g_{u}\left(\partial_{x} u, \partial_{x} u\right)+c R(u)\right) d t\right) D u
$$

for different c's: Onsager-Machlup $\left(-\frac{1}{6}\right)$, DeWitt $\left(\frac{1}{6},-\frac{1}{4}\right)$, Dekker $\left(\frac{1}{4}\right)$, Inoue, Maeda $\left(-\frac{1}{6}\right)$, Andersson, Driver $\left(0,-\frac{1}{3}\right)$, etc. Our choice suggests $c=-\frac{1}{4}$.

## Main step in the proof

One shows that 'geometric' and 'Itô' solutions differ by a counterterm in $\mathcal{S}^{\text {both }}$ : terms $\tau \in\langle\mathcal{S}\rangle$ such that $\left(\Upsilon_{\Gamma, \sigma}-\Upsilon_{\Gamma, \bar{\sigma}}\right) \tau$ is a vector field. One obviously has $\mathcal{S}^{\text {tho }}+\mathcal{S}^{\text {geo }} \subset \mathcal{S}^{\text {both. }}$. Non-trivial fact:

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For SDEs, $\mathcal{S}=\{\mathrm{B}\}$ and $\Upsilon_{\sigma}{ }^{\mathrm{g}}=\sigma_{i} D \bar{\sigma}_{i}$, so $\mathcal{S}^{1 \mathrm{toj}}=\mathcal{S}^{\mathrm{g} 80}=0$. But $\left(\Upsilon_{\sigma}-\Upsilon_{\bar{\sigma}}\right)$ ) $=\nabla_{\sigma_{i}} \sigma_{i}-\nabla_{\bar{\sigma}_{i}} \bar{\sigma}_{i}$, so $\mathcal{S}^{\text {both }}=\mathcal{S}!$

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For SPDEs, one has $\mathcal{S}_{(2)}=\{\varepsilon, \infty\}$ and

$$
\mathcal{S}_{(2)}^{1+\hat{\omega}}=\langle\alpha\rangle, \quad \mathcal{S}_{(2)}^{\mathrm{geo}}=\langle\alpha+q\rangle, \quad \mathcal{S}_{(2)}^{\text {both }}=\mathcal{S}_{(2)} .
$$

Much harder to check at level 4 , requires systematic approach.

## $T$-algebras

Motivation: abstraction of functions with multiple 'upper' and 'lower' indices.

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Definition: A $T$-algebra is a bigraded vector space $\mathcal{V}=\bigoplus\left\{\mathcal{V}_{\ell}^{u}: u, \ell \geq 0\right\}$ with

- An action of $\operatorname{Sym}(u, \ell)=\operatorname{Sym}(u) \times \operatorname{Sym}(\ell)$ on each $\mathcal{V}_{\ell}^{u}$.


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- An associative product $\mathcal{V}_{\ell_{1}}^{u_{1}} \times \mathcal{V}_{\ell_{2}}^{u_{2}} \rightarrow \mathcal{V}_{\ell_{1}+\ell_{2}}^{u_{1}+u_{2}}$ satisfying

$$
B \cdot A=S_{\ell_{1}, \ell_{2}}^{u_{1}, u_{2}}(A \cdot B), \quad \alpha_{1} A \cdot \alpha_{2} B=\left(\alpha_{1} \cdot \alpha_{2}\right)(A \cdot B) .
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$$

$\triangleright$ A trace tr: $\mathcal{V}_{\ell+1}^{u+1} \rightarrow \mathcal{V}_{\ell}^{u}$ with $\operatorname{tr}(A \cdot B)=A \cdot \operatorname{tr} B($ if $\operatorname{deg} B \geq(1,1))$ and

$$
\alpha \operatorname{tr} A=\operatorname{tr}\left(\left(\alpha \cdot \mathrm{id}_{1}^{1}\right) A\right), \quad \operatorname{tr}^{2} A=\operatorname{tr}^{2}\left(\left(\mathrm{id}_{\ell}^{u} \cdot S_{1,1}^{1,1}\right) A\right) .
$$

## Examples

Canonical example: Given a vector space $V$, set

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\mathscr{V}[V]_{\ell}^{u}=\left(V^{*}\right)^{\otimes \ell} \otimes V^{u}
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Natural product and action of permutations. Trace pairs up last factors.

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Additional structure: derivation $\partial: \mathcal{V}_{\ell}^{u} \rightarrow \mathcal{V}_{\ell+1}^{u}$ via

$$
L\left(V, \mathscr{V}[V]_{\ell}^{u}\right) \simeq V^{*} \otimes \mathscr{V}[V]_{\ell}^{u} \simeq \mathscr{V}[V]_{\ell+1}^{u}
$$

if $V$ finite-dimensional. Satisfies $\partial^{2} A=\left(S_{1,1} \cdot \mathrm{id}_{\ell}^{u}\right) \partial^{2} A$, plus Leibniz rule and natural interaction with trace and symmetric group.

## Free $T$-algebras

Given by ' $T$-graphs' with nodes decorated by generators.

Given $W=\bigoplus\left\{W_{\ell}^{u}: u, \ell \geq 0\right\}$ with action of symmetric group, generates a $T$-algebra $\operatorname{Tr}(W)$. Every $T$-graph $g$ yields a subspace $\operatorname{Tr}_{g}(W) \subset \operatorname{Tr}(W)$.

## Non-degeneracy result

Fix $W=\bigoplus\left\{W_{\ell}^{u}: u, \ell \geq 0\right\}$ locally finite-dimensional with action of symmetric group, $\hat{W}_{\ell}^{u} \subset W_{\ell}^{u}$ invariant, and finite collection $G$ of connected anchored $T$-graphs.

Theorem: There exists $V$ finite-dimensional, $\Phi, \bar{\Phi} \in \operatorname{Hom}(\operatorname{Tr}(W), \mathscr{V}[V])$ injective on $\operatorname{Tr}_{G}(W)$ such that, for $\tau \in \operatorname{Tr}_{G}(W), \Phi \tau=\bar{\Phi} \tau$ if and only if $\tau \in \operatorname{Tr}(\hat{W})$.
If $W$ (and therefore $\operatorname{Tr}(W)$ ) admits a derivation, same holds with $\mathscr{V}[V]$ replaced by $\mathscr{W}[V]$ and $\Phi, \Phi \in \operatorname{Hom}_{\partial}(\operatorname{Tr}(W), \mathscr{W}[V])$.

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Remark: $\Phi$ certainly cannot be injective on all of $\operatorname{Tr}(W)$ since $\operatorname{dim} \operatorname{Tr}(W)_{\ell}^{u}=\infty$ but $\operatorname{dim} \mathscr{V}[V]_{\ell}^{u}<\infty!$

## Some open questions

- Minimal dimension required for $V$ in non-degeneracy result?
- More intrinsic "geometric" formulation of solution theory?
- Behaviour in sub-Riemannian case, notion of hypoellipticity?
- Large deviations between closed geodesics?
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## Thank you for your attention!

