

Radiative and Quantum Electrodynamics processes in **SMILEI**

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I. Presentation of the radiative and QED physical processes relevant for the next generation of extreme intensity lasers

II. Regimes of interaction

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IV. Presentation of the multiphoton Breit-Wheeler pair creation model

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Radiative and QED physical processes relevant for the next generation of extreme intensity lasers

Forthcoming short-pulse extreme intensity laser will soon reach the threshold of 10²²-10²³ W/cm²





Radiative and quantum electrodynamics effects will progressively come into play



Classical radiation cooling

Continuous emission of relatively low-energy photons in comparison with the electron energy



Nonlinear inverse Compton scattering

Stochastic emission of high-energy photon



Radiative and quantum electrodynamics effects will progressively come into play



Multiphoton Breit-wheeler pair creation

Creation a pair of electron-positron from the decay of a γ photon interacting with an electromagnetic field



Radiative and QED effects will affect the particle dynamic and the overall energy balance of simulations

Example of the laser thin-foil interaction

- 10^{22} W/cm²: ~ 1 % absorption by radiation
- 10^{23} W/cm²: ~ 10 % absorption by radiation
- 10^{24} W/cm² : > 50 % absorption by radiation, significant production of pairs



Novel experimental configurations to study QED processes in laboratories will be soon accessible with extreme-intensity lasers





Regimes of interaction



Radiative and QED effects depend on the electron/photon energy, the field strength and the interaction geometry

To determine the right regime of interaction, pay attention to these parameters



- $arepsilon_{oldsymbol{\gamma}}$ The photon energy
- m The electron mass
- c Speed of light in vacuum
- e Electron charge
- v_{\pm} Electron or positron velocity



To determine the right regime of interaction, pay attention to these parameters

The strength of the QED effects on an electron of Lorent γ_{\pm} actor and a photon of normalized energy propagating in an electromagnetic field and depends on the following Lorentz invariant quantum parameters:

Particle quantum parameter:

$$\chi_{\pm} = \frac{\gamma_{\pm}}{E_s} \sqrt{\mathbf{E}_{\parallel}^2 / \gamma_{\pm}^2 + (\mathbf{E}_{\perp} + \mathbf{v}_{\pm} \times \mathbf{B})^2}$$

Photon quantum parameter:

$$\chi_{\gamma} = \frac{\gamma_{\gamma}}{E_s} \sqrt{\left(\mathbf{E}_{\perp} + \mathbf{c} \times \mathbf{B}\right)^2}$$

- E_{\parallel} Electric field parallel to the propagation direction
- E_{\perp} Electric field orthogonal to the propagation direction
- B Magnetic field
- v_+ Particle velocity
- E_s Schwinger field

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Electromagnetic fields are constant during the emission process

$$\xi_1 = \frac{\mathbf{B}^2 - \mathbf{E}^2}{E_s^2} \ll 1 \qquad \qquad \xi_2 = \frac{\mathbf{B} \cdot \mathbf{E}}{E_s^2} \ll 1 \qquad \qquad \chi_{\pm} \gg \max\left(\xi_1, \xi_2\right)$$

- E Electric field
- *B* Magnetic field
- E_s Schwinger field



$$\chi_{\pm} \le 10^{-3}$$

- Radiation reaction negligible
- No impact on particle energy



$$\chi_{\pm} \in \left[10^{-3}, 10^{-2}\right]$$

- Radiation reaction can be seen as a slow continuous energy damping

- Emission of many photons of negligible energy in comparison with the emitting particle energy

Laser-thin foil interaction $I_0 \sim 10^{22}$ W/cm², $a_0 \sim 90$ at $\lambda = 1 \ \mu m$ - Weak impact on the target heating and ion acceleration

Laser-electron beam collision

 $I_{_0}{\sim}10^{_{21}}\,W/cm^2,\,a_{_0}\sim 1$ at $\lambda=1\,\mu m,$ electron beam energy ~ 100 MeV - $\,1$ GeV



$$\chi_{\pm} \in \left[10^{-2}, 10^{-1}\right]$$

- Radiation reaction can still be seen as a continuous energy damping

- Stronger slow down

Laser-thin foil interaction $I_0 \sim 10^{23}$ W/cm², $a_0 \sim 270$ at $\lambda = 1 \mu m$ - Stronger impact on the target heating and ion acceleration - Energy loss in radiation can represent several percents

Laser-electron beam collision

 $I_{_0}{\sim}10^{22}$ W/cm², $a_{_0}{\sim}$ 1 at λ = 1 $\mu m,$ electron beam energy ${\sim}$ 100 MeV - 1 GeV



$$\chi_{\pm} \in \left[10^{-1}, 1\right]$$

- Radiation reaction can not longer be seen as a continuous energy damping

- Emission of photons with individual energies no longer negligible
- Stochastic effects come into play (recoil effect)

Laser-thin foil
interaction
$$I_0 \sim 10^{24}$$
 W/cm², $a_0 \sim 850$ at $\lambda = 1 \ \mu m$
- Strong impact on the hole
boring + TNSA, reduced ion
energy

Laser-electron beam collision

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 $I_{_0}{\sim}10^{_{23}}\,W/cm^2,\,a_{_0}\sim270$ at $\lambda=1$ $\mu m,$ electron beam energy \sim few GeV



$\chi_{\pm} \ge 1$

- Emission of photons with high individual energies, potentially of the order of the emitting particle energy

- Important Stochastic effects (strong recoil)

Laser-thin foil interaction $I_0 > 10^{24} \text{ W/cm}^2$, $a_0 > 850 \text{ at } \lambda = 1$ μm - Strong impact on the hole boring + TNSA, reduced ion energy

Laser-electron beam collision

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 $I_{_0}{\sim}10^{_{23}}\,W/cm^2,\,a_{_0}\sim270$ at $\lambda=1$ $\mu m,$ electron beam energy \sim few GeV



$$\chi_{\gamma} \ge 0.1$$

- Nonlinear Breit-Wheeler pair creation start to be important $\chi_{\gamma} \sim 1$ - Maximum production rate around $\chi_\gamma \sim 10$ around





Regime / Quantum parameter	Impacts	Thin foil - laser interaction	Electron beam - laser collision
Quantum regime/pair creation $\chi_{\gamma} \geq 0.1$	Important creation of pairs	$I_0 \sim 10^{24}$ W/cm ² , $a_0 \sim 850$ at $\lambda = 1 \ \mu m$ Creation of a pair plasma on the rear target	$I_0 \sim 10^{23} \text{ W/cm}^2$, $a_0 \sim 270 \text{ at } \lambda = 1 \ \mu\text{m}$, electron beam energy \sim few GeV
Quantum regime $\chi_{\pm} \geq 1$	Strong recoil Emission of high-energy gamma photons	$I_0 \sim 10^{24}$ W/cm ² , $a_0 \sim 850$ at $\lambda = 1 \ \mu m$ Strong impact on the hole boring + TNSA, reduced ion energy	
$\label{eq:constraint} \begin{array}{c} {\rm Semi-quantum}\\ {\rm regime} \end{array}$ $\chi_{\pm} \sim 10^{-2} - 10^{-1}$	Strong continuous damping	$I_0 \sim 10^{23} \text{ W/cm}^2$, $a_0 \sim 270$ at $\lambda = 1 \ \mu \text{m}$	$\begin{split} I_{_0} \sim & 10^{22} \text{ W/cm}^2 \text{, } a_{_0} \sim 1 \\ \text{at } \lambda = 1 \mu\text{m} \text{, electron} \\ \text{beam energy} \sim & 100 \\ \text{MeV} - 1 \text{ GeV} \end{split}$
Classical regime $\chi_{\pm} \sim 10^{-3} - 10^{-2}$	Continuous damping Emission of many neglegible- energy photon	$I_0 \sim 10^{22}$ W/cm ² , $a_0 \sim 90$ at $\lambda = 1 \ \mu m$ Weak impact on the target heating and ion acceleration	$I_0 \sim 10^{21}$ W/cm ² , $a_0 \sim 1$ at $\lambda = 1 \mu$ m, electron beam energy ~ 100 MeV - 1 GeV



Presentation of the radiative models



<u>Validity</u> : $\chi_{\pm} \leq 10^{-2} \quad \gamma_{\pm} \gg 1$

<u>Principles</u>: Modification of the momentum equation by the addition of a friction term to model the radiation damping

acceleration in the laser field charged particle

$$\frac{d\mathbf{p}}{dt} = \mathbf{F}_{Lorentz} + \mathbf{F}_{rad} \qquad \mathbf{F}_{rad} = -P_{rad} \frac{\mathbf{u}}{u^2 c}$$

Derived from the Landau-Liftshitz model approximated at high gamma factors

Overhead : low (vectorized)



acceleration i the laser field

charged particle

<u>Validity</u> $\chi_{\pm} \leq 10^{-1} \quad \gamma_{\pm} \gg 1$

<u>Principles</u>: Modification of the momentum equation by the addition of a friction term to model the radiation daming

$$\frac{d\mathbf{p}}{dt} = \mathbf{F}_{Lorentz} + \mathbf{F}_{rad} \qquad \mathbf{F}_{rad} = -P_{rad} g\left(\chi_{\pm}\right) \frac{\mathbf{u}}{u^2 c}$$

Derived from the Landau-Liftshitz model at high gamma factors Include a quantum correction derived from the QED cross-section

Overhead : low (vectorized)

Fig. : Importance of the quantum correction when the quantum parameter reaches 0.1.





The Fokker-Planck radiation of Niel et al.

<u>Validity</u> $\chi_{\pm} \leq 1 \quad \gamma_{\pm} \gg 1$

Principles : Extension of the corrected Landau-Lifshitz model with an operator that takes into account diffusive stochastic effects

$$\mathbf{F}_{rad}dt = \left(-P_{rad} g\left(\chi_{\pm}\right) + mc^2 \sqrt{R\left(\chi_{\pm}, \gamma_{\pm}\right)} dW\right) \frac{\mathbf{u}}{u^2 c}$$

$$R(\chi,\gamma)=rac{2}{3}rac{lpha^2}{ au_e}h(\chi) ~~h(\chi)$$
 computed from the Compton cross-section

dW is a Wiener process of variance dt

Overhead : high (partly vectorized) / under improvement

F. Niel et al., ArXiv:1707.02618 (2017)



<u>Validity</u> $\chi_{\pm} \ge 10^{-2}$ $\gamma_{\pm} \gg 1$ $a_0 \gg 1$ $\chi_{\pm} \gg \max(\xi_1, \xi_2)$

Principles : Monte-Carlo process



Random drawing of an optical depth computed from the Compton photon emission rate

Cumulative1 distribution function computed from the Compton photon energy function for the random drawing of the emitted photon energy1 Independent of the PIC iteration (particle pusher





<u>Overhead</u> : high (not vectorized)

R. Duclous et al., Plasma Physics and Controlled Fusion, 53 (1), 015009 (2011) ; M. Lobet et al., J. Phys.: Conf. Ser. 688, 012058 (2016)

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Presentation of the Multiphoton Breit-Wheeler pair creation model

The Monte-Carlo pair creation emission process

Validity:
$$\chi_{\gamma} \ge 10^{-2}$$
 $\gamma_{\gamma} \ge 2$ $a_0 \gg 1$ $\chi_{\pm} \gg \max(\xi_1, \xi_2)$

<u>Principles</u>: Applied on macro-photon species (from the Compton or initialized)



Monte-Carlo process

Optical depth computed from the multiphoton Breit-Wheeler annihilation rate

Cumulative distribution function for the random drawing of the created electron and positron energy from the multiphoton Breit-Wheeler Monte Chate Independent of the PIC iteration / particle pusher

Overhead : high (not vectorized)



Examples of simulations



Radiative model comparison using particle tests





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Generation of pairs







Conclusion



Different models implemented for different regimes of interaction

Quantum parameters	Interaction regime	Models	
$\chi_{\gamma} \ge 0.1$	Pair creation regime	Emission/propagation of macro-photon as for macro-particles Monte-Carlo pair creation process on the photon species	
$\chi_{\pm} \ge 1$	Quantum radiation emission regime	Monte-Carlo emission process	
$\chi_{\pm} \sim 10^{-2} - 10^{-1}$	Near quantum regime	Monte-Carlo emission process Fokker-Planck radiation model of Niel <i>et al.</i>	
$\chi_{\pm} \sim 10^{-2} - 10^{-1}$	Semi-quantum regime	Fokker-Planck radiation model of Niel <i>et al.</i> Corrected Landau-Lifshitz model	
$\chi_{\pm} \sim 10^{-3} - 10^{-2}$	Classical regime	Corrected Landau-Lifshitz model Classical Landau-Lifshitz model	



Radiation reaction processes :

- Classical and corrected Landau-Lifshitz
- Niel *et al.*
- Monte-Carlo with or without macrophoton creation

Multi-photon Breit-Wheeler :

Monte-Carlo pair creation process

SMILEI version 3.3

Future related features :

• Particle and photon merging process

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