

What do Black Hole Microstates look like ?

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with

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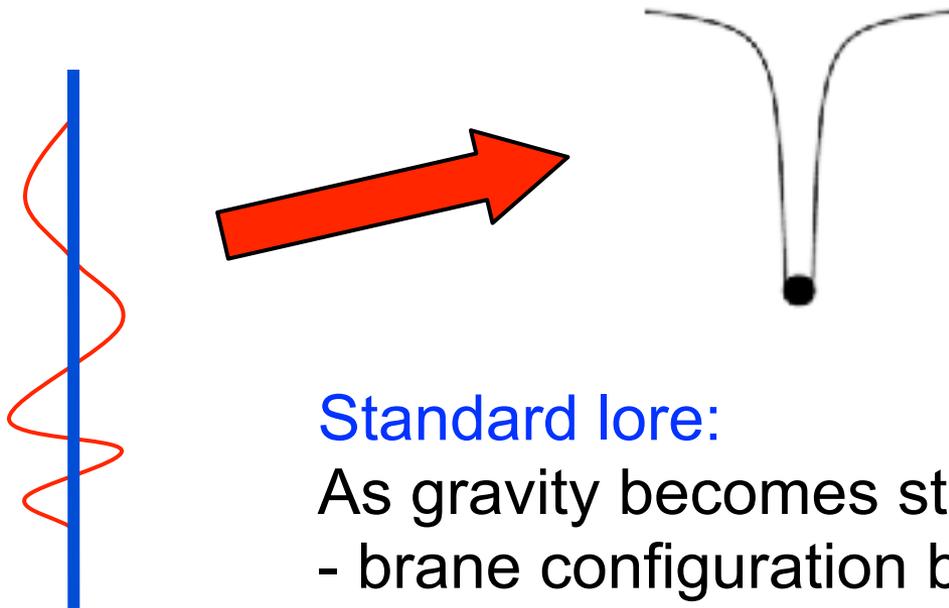
DE LA RECHERCHE À L'INDUSTRIE
cea
SACLAY

Strominger and Vafa (1996):

*Black Hole Microstates at **Zero Gravity*** (branes + strings)

Correctly match B.H. entropy !!!

One Particular Microstate at **Finite Gravity**:



Standard lore:

As gravity becomes stronger,

- brane configuration becomes smaller
- horizon develops and engulfs it
- recover standard black hole

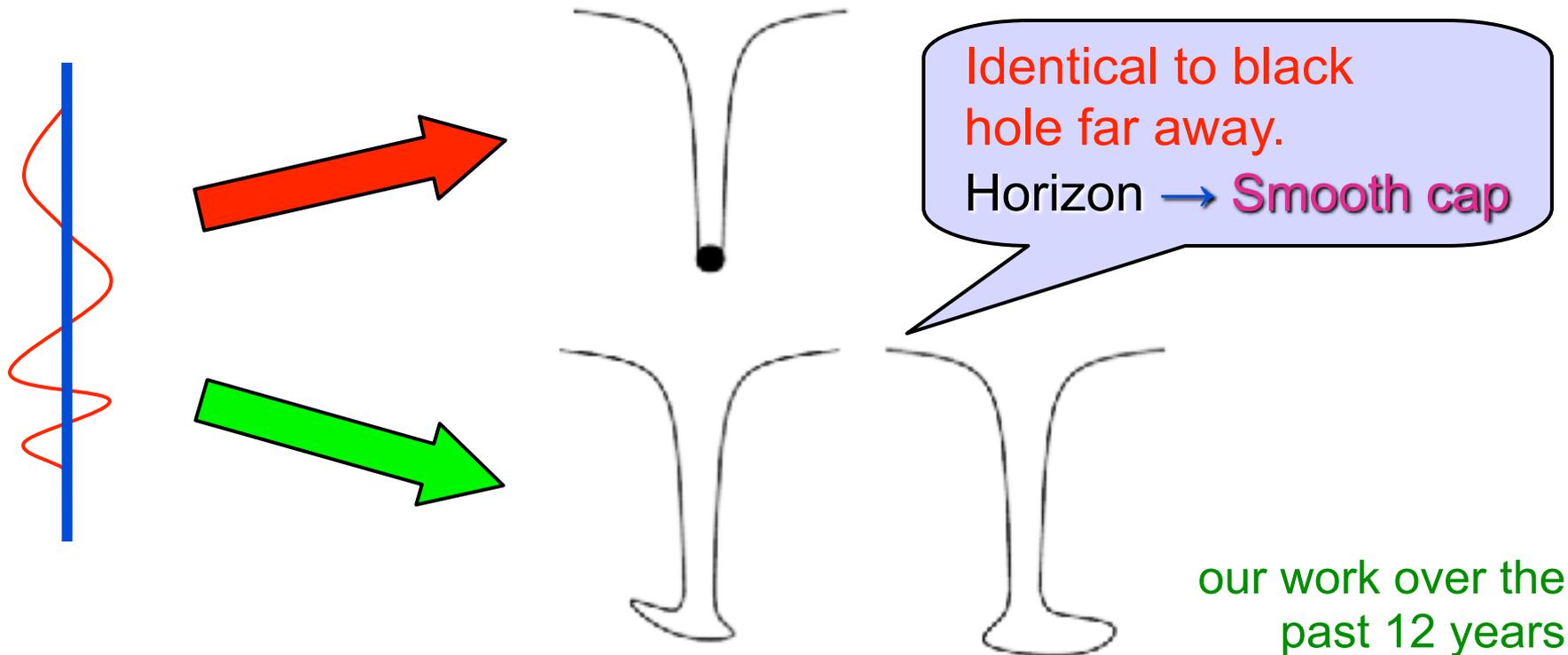
Susskind
Horowitz, Polchinski
Damour, Veneziano

Strominger and Vafa (1996):

*Black Hole Microstates at **Zero Gravity*** (branes + strings)

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One Particular Microstate at **Finite Gravity**:

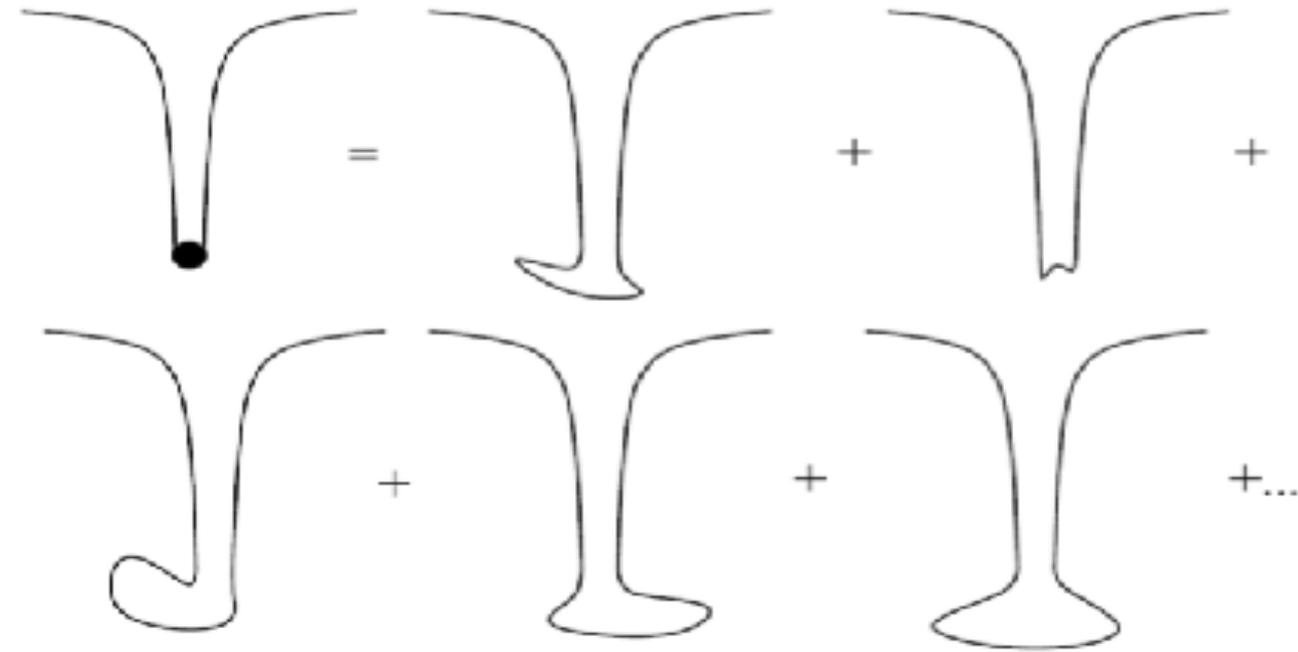


BIG QUESTION: Are *all* black hole microstates becoming geometries with no horizon ?

?

Black hole = ensemble of horizonless microstate configurations

Mathur 2003



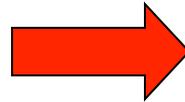
Analogy with ideal gas

Thermodynamics

(Air = ideal gas)

$$P V = n R T$$

$$dE = T dS + P dV$$



Statistical Physics

(Air -- molecules)

e^S microstates

typical

atypical

Thermodynamics

Black Hole Solution



Statistical Physics

Microstate geometries

Long distance physics

Gravitational lensing

Physics at horizon

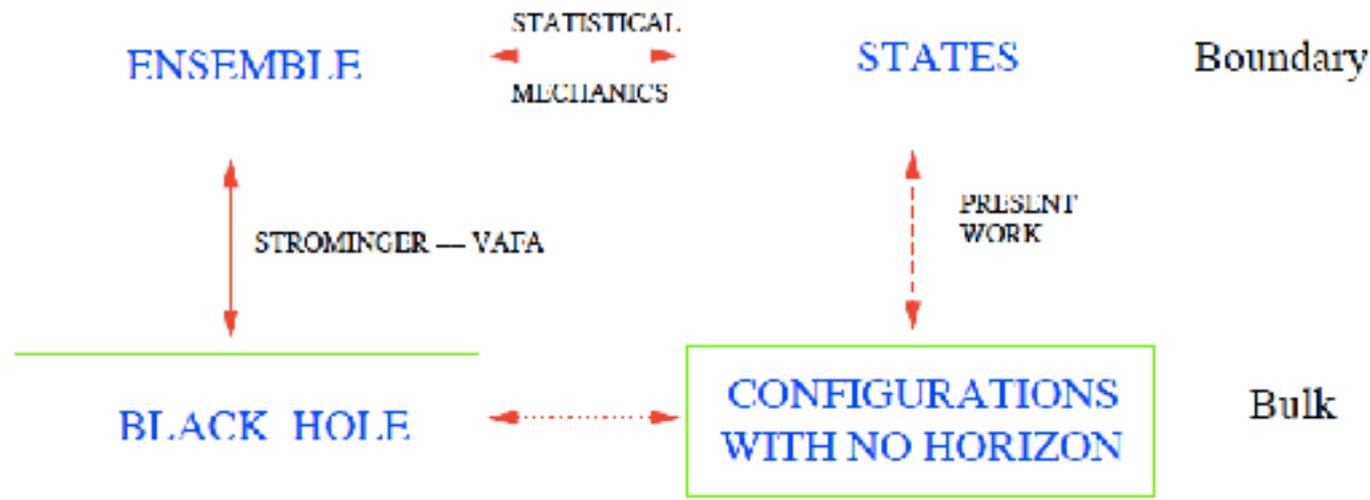
Information loss

Gravity waves ?

Other formulations:

e.g. Bena, Warner, 2007

- **Thermodynamics (EFT)** breaks down at horizon. New **low-mass d.o.f.** kick in.
- No spacetime inside black holes. **Quantum superposition** of microstate geometries.



Not some **hand-waving** idea - **provable** by rigorous calculations in String Theory

Word of caution

- To replace classical BH by BH-sized object
 - Gravastar
 - Infinite density firewall hovering above horizon
 - LQG configuration
 - Quark-star, you name it ...satisfy 3 very stringent tests:

1. Same growth with G_N !!!

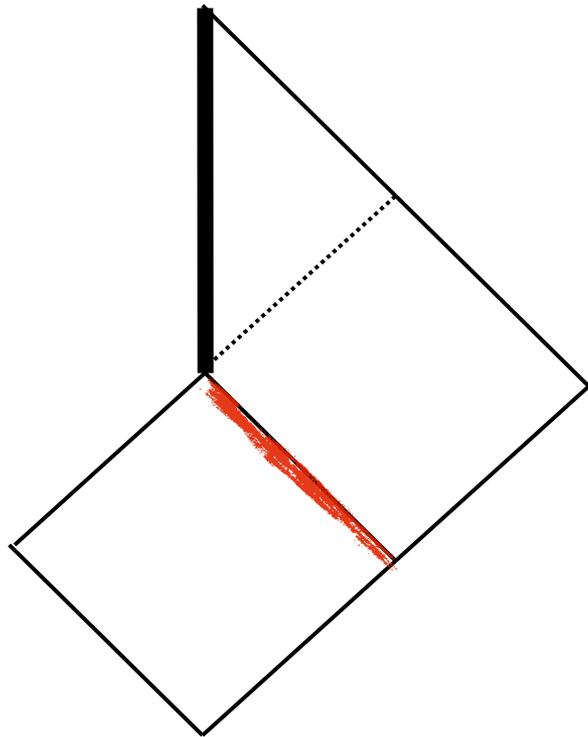
Horowitz

- BH size grows with G_N
- Size of objects in other theories becomes smaller

- BH microstate geometries pass this test
- Highly nontrivial mechanism:
- D-branes = solitons, tension $\sim 1/g_s \rightarrow$ lighter as G_N increases

2. Mechanism not to fall into BH

Very difficult !!!



GR Dogma:

**Thou shalt not put anything
at the horizon !!!**

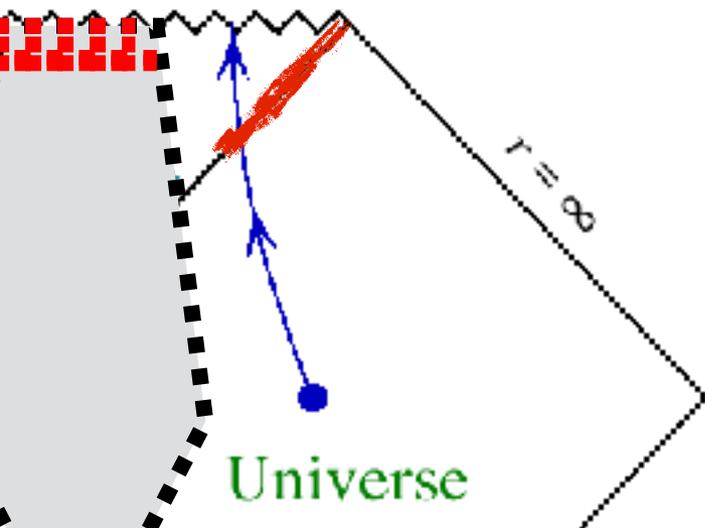
- Null \rightarrow speed of light.
- If massive: ∞ boost \rightarrow ∞ energy
- If massless: dilutes with time
- Nothing can live there !
(or carry degrees of freedom)
- No membrane, no spins
- No (fire)wall

Otherwise b.s.

Must have a support mechanism !

3. Avoid forming a horizon

- Collapsing shell forms horizon Oppenheimer and Snyder (1939)
- If curvature is low, no reason not to trust classical GR
- By the time shell becomes **curved-enough for quantum effects to become important**, horizon in causal past



Go backwards in time !

BH has e^S microstates with no horizon

Small tunneling probability = e^{-S}

Will tunnel with probability **ONE !!!**

Kraus, Mathur; Bena, Mayerson, Puhm, Vercnocke

Only e^S horizon-sized microstates can do it !

BPS Microstates geometries - 11D SUGRA / T⁶

5D 3-charge BH (Strominger-Vafa)

Linear system \mathbb{R}^4 base (4D Hyper Kahler)

4 layers: $*G^I = G^I$

Bena, Warner
Gutowski, Reall

$$d * dZ_1 = G^2 \wedge G^3$$

$$d\vec{k} + *d\vec{k} = G^1 Z_1 + G^2 Z_2 + G^3 Z_3$$

Focus on Gibbons-Hawking (Taub-NUT) base:

$$ds^2 = V (dx_1^2 + dx_2^2 + dx_3^2) + V^{-1} (d\psi + \vec{A})^2$$

$$\nabla \times \vec{A} = \nabla V$$

$$V = \frac{1}{r}$$

\mathbb{R}^4

$$V = 1 + \frac{1}{r}$$

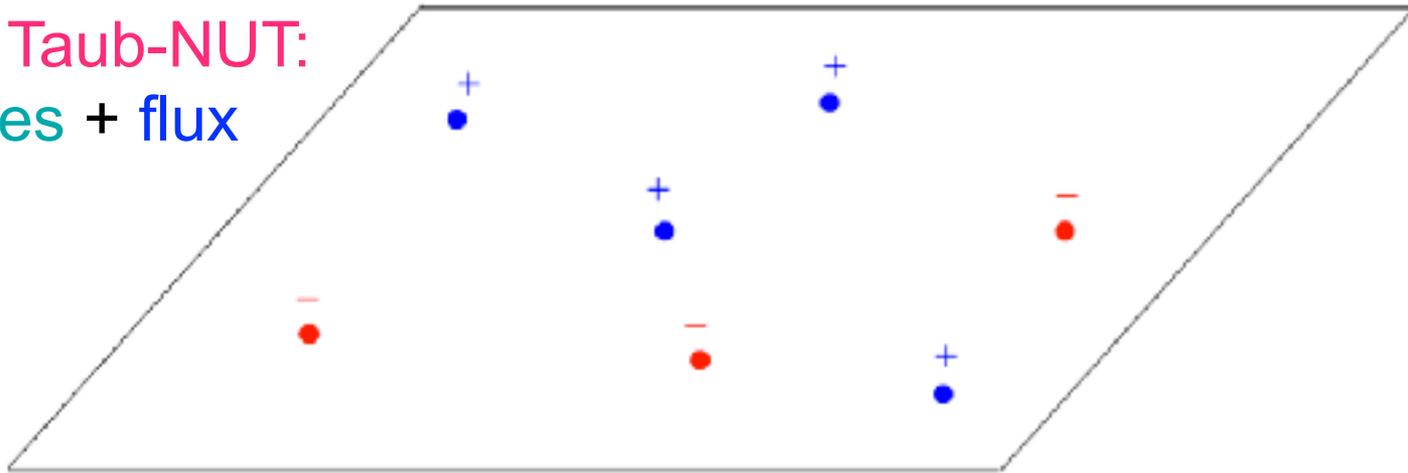
Taub-NUT

8 harmonic functions

Gauntlett, Gutowski,
Bena, Kraus, Warner

Simplest Microstate Geometries

Multi-center Taub-NUT:
many 2-cycles + flux



Base **singular** (signature changing sign)

Full solution **smooth** (@ Taub-NUT centers $\sim \mathbb{R}^4$)

Same **mass, charge, J, size** as BH with large horizon area

Lots of solutions !

Microstates geometries

- Where is the **BH charge** ?

$$L = q A_0$$

magnetic

$$L = \dots + A_0 F_{12} F_{34} + \dots$$

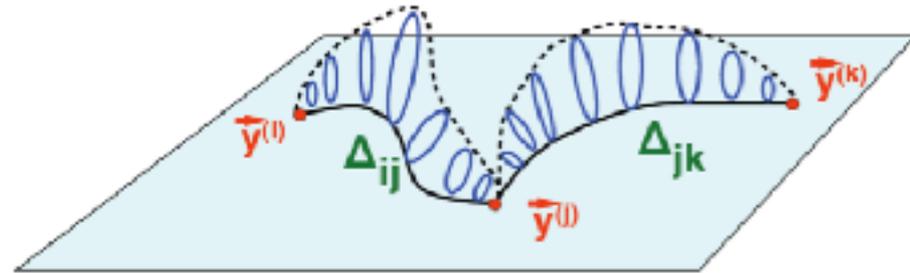
- Where is the **BH mass** ?

$$E = \dots + F_{12} F^{12} + \dots$$

- **BH angular momentum**

$$J = E \times B = \dots + F_{01} F_{12} + \dots$$

2-cycles + magnetic flux



Bubbling Geometries

Black Hole Solitons

beautiful GR story behind

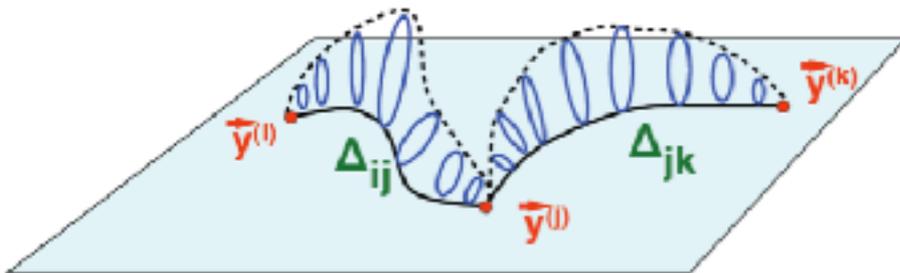
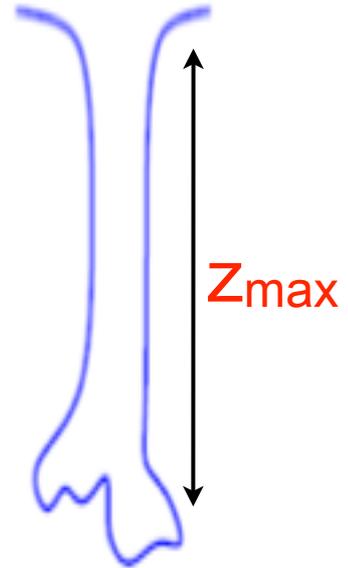
Gibbons, Warner

The charge is dissolved in magnetic fluxes. No singular sources.

Klebanov-Strassler

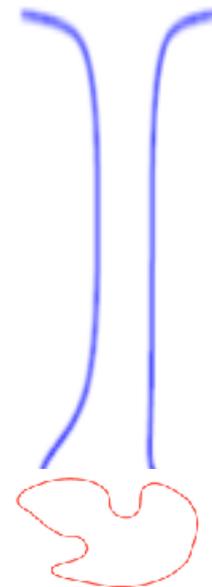
Four Scales

- Classical BH has 2 scales:
 - Mass / Horizon Size
 - Planck Length
- Microstate geometries have 2 more
 - **Redshift** from the bottom of the throat, Z_{\max}
 - **Size of bubbles:** $\lambda_T \sim k \ell_P$



More general bubbling solutions

- Add supertubes
 - supersymmetric brane configs
 - arbitrary shape Mateos, Townsend
- Construct backreacted solution
 - Taub-NUT Page Green's functions (painful)
- Smooth !
 - exactly as in flat space
 - Lunin, Mathur; Emparan, Mateos, Townsend
 - Lunin, Maldacena, Maoz
- Entropy: $S \sim (Q^{5/2})^{1/2}$
- Not yet black-hole-like ($Q^{3/2}$)
- Need more degrees of freedom !

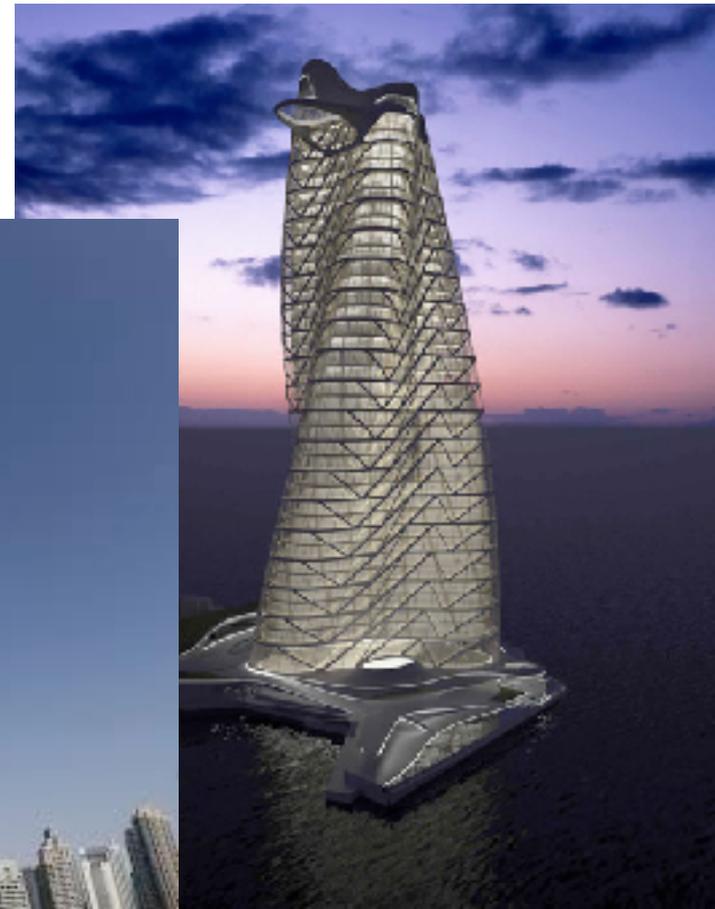
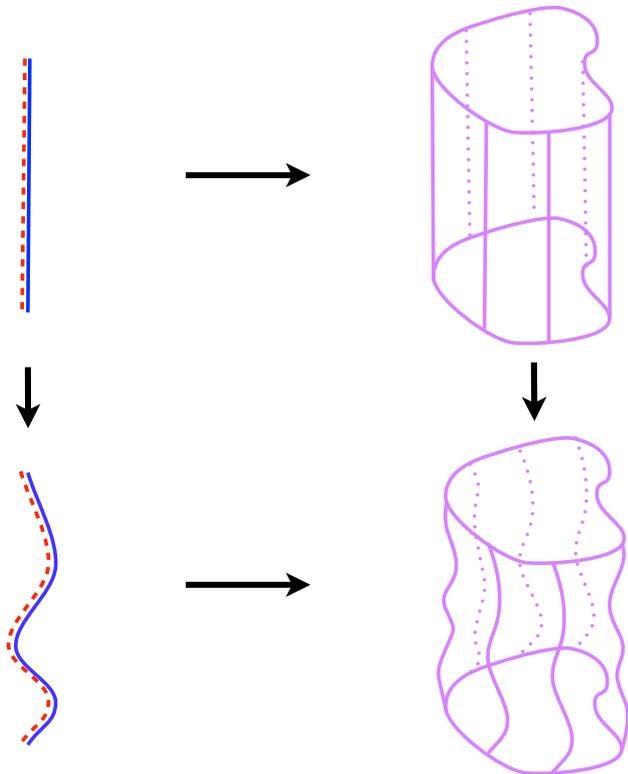


Even more general solutions

Bena, de Boer, Shigemori, Warner

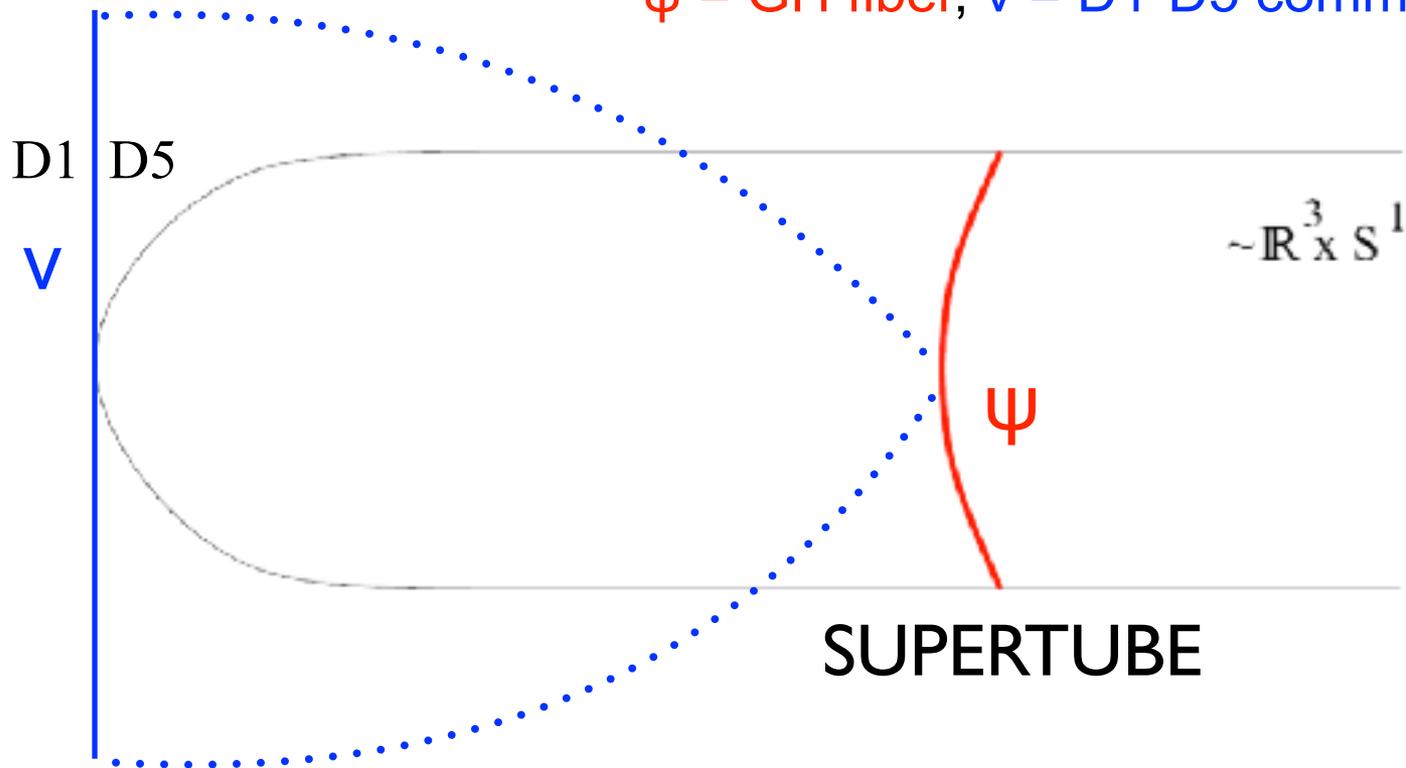
- **Supertubes (locally 16 susy)** - 8 functions of **one** variable ($c = 8$)
- **Superstrata (locally 16 susy)** - 4 functions of **two** variables ($c = \infty$)
- Double supertube transition:

Should be Smooth !!!



Habemus Superstratum

ψ = GH fiber, v = D1-D5 common direction



- ψ -dependent solutions=supertubes Lunin, Mathur; Taylor, Skenderis
interchange fibers: v -dependent solutions
- Constructed smooth solution parameterized by arbitrary
function of 2 variables $F(\psi, v)$
Bena, Giusto, Russo, Shigemori, Warner

String Theory to the rescue

- Superstrata conjectured in 2011
constructed in 2015
- 5D microstates with GH bubbles: $U(1)^3$
- Oscillations \rightarrow singularities
- Precision Holography: Skenderis, Taylor, Kanitscheider
- Open string emission: Giusto, Russo, Turton
- *There is another Skywalker !*
- At least $U(1)^4$
- Metric depends on $Z_1 Z_2 - Z_4^2$ Coiffuring
Bena, Ross, Warner
- Singularities cancel - solution smooth !!!



Largest family of solutions known to mankind

- Functions of **two** variables: $\infty \times \infty$ parameters

$$ds_{10}^2 = \frac{1}{\sqrt{\alpha}} ds_8^2 + \sqrt{\frac{Z_1}{Z_2}} ds_4^2,$$

$$ds_8^2 = -\frac{2}{\sqrt{\mathcal{P}}} (dv + \beta) \left[du + \omega + \frac{\mathcal{J}}{2} (dv + \beta) \right] + \sqrt{\mathcal{P}} ds_4^2,$$

$$e^{2\sigma} = \frac{Z_1^2}{\mathcal{P}},$$

$$B = \frac{Z_4}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_4 \wedge (dv + \beta) + \delta_2,$$

$$C_0 = \frac{Z_4}{Z_1},$$

$$C_2 = \frac{Z_2}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_1 \wedge (dv + \beta) + \gamma_2,$$

$$C_4 = \frac{Z_4}{Z_2} \widehat{\text{vol}}_4 - \frac{Z_4}{\mathcal{P}} \gamma_2 \wedge (du + \omega) \wedge (dv + \beta) + x_3 \wedge (dv + \beta) + \mathcal{C}$$

$$C_6 = \widehat{\text{vol}}_4 \wedge \left[-\frac{Z_1}{\mathcal{P}} (du + \omega) \wedge (dv + \beta) + a_2 \wedge (dv + \beta) + \gamma_1 \right]$$

$$\frac{Z_4}{\mathcal{P}} \mathcal{C} \wedge (du + \omega) \wedge (dv + \beta),$$

$$\alpha \equiv \frac{Z_1 Z_2}{Z_1 Z_2 - Z_4^2} \quad \mathcal{P} \equiv Z_1 Z_2 - Z_4^2.$$

$$\omega_r^{(2)} = -\frac{Rr}{\sqrt{2} k_2 (m_1^2 - 1)} \frac{m_1 (k_2 + m_1 + 1) \Delta_{k_2 + m_1 - 1, m_1 - 1} + (k_2 + m_1 - 1) \Delta_{k_2 + m_1}}{(r^2 + a^2)^2}$$

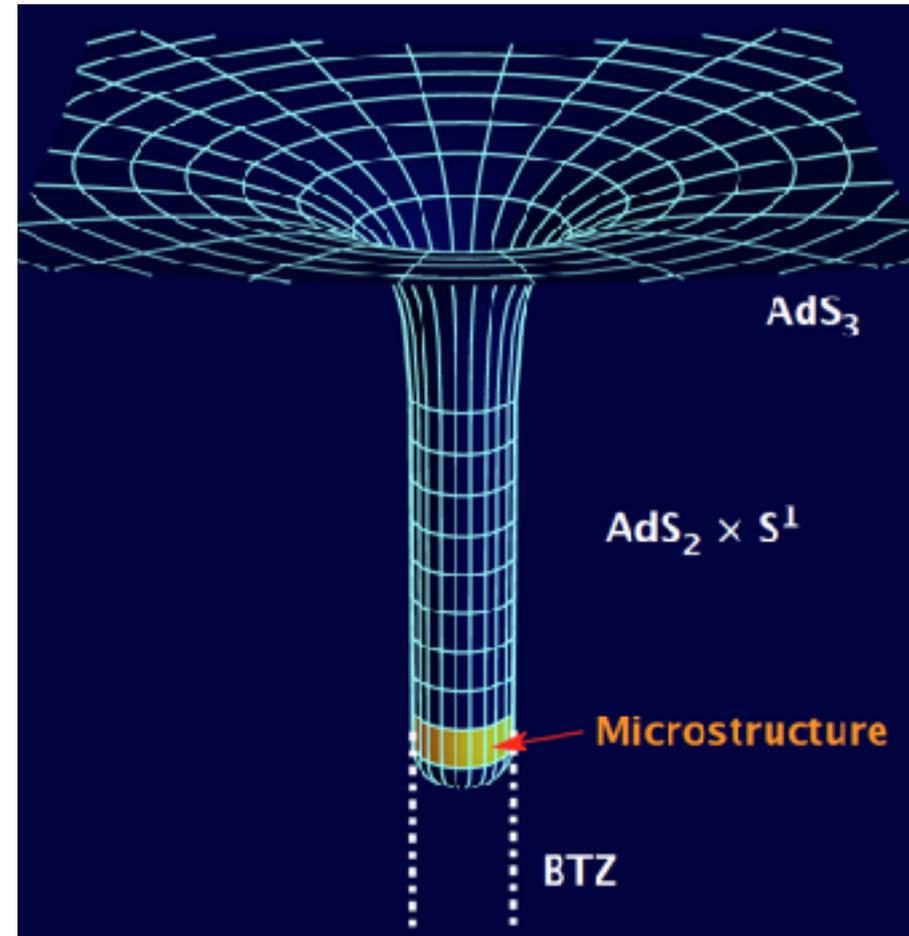
$$\omega_\theta^{(2)} = \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2 \sin \theta \cos \theta} \left[2(m_1 - 1) \Delta_{k_2 + m_1 - 3, m_1 - 1} \right. \\ \left. + (m_1 - 1)(m_1 - 2) \Delta_{k_2 + m_1 - 1, m_1 - 1} + m_1 (k_2 - 2) \Delta_{k_2 + m_1 - 1, m_1 + 1} \right. \\ \left. - m_1 (m_1 - 1) \Delta_{k_2 + m_1 + 1, m_1 - 1} + (m_1^2 (k_2 - 1) + 1) \Delta_{k_2 + m_1 + 1, m_1 + 1} \right],$$

$$\omega_\phi^{(2)} = -\frac{R}{\sqrt{2}} \frac{\Delta_{k_2 + m_1 + 1, m_1 + 1}}{\Sigma} \sin^2 \theta - \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2} \left[2(m_1 - 1) \Delta_{k_2 + m_1 - 3, m_1 - 1} \right. \\ \left. + (m_1^2 - 2m_1 + k_2 - 1) \Delta_{k_2 + m_1 - 1, m_1 - 1} + m_1 (k_2 - 2) \Delta_{k_2 + m_1 - 1, m_1 + 1} \right. \\ \left. + m_1 (k_2 - m_1 - 1) \Delta_{k_2 + m_1 + 1, m_1 - 1} + (k_2 (m_1^2 + m_1 - 1) - m_1 (m_1 + 1)) \Delta_{k_2 + m_1 + 1, m_1 + 1} \right]$$

$$\omega_\delta^{(2)} = -\frac{R}{\sqrt{2}} \frac{\Delta_{k_2 + m_1 + 1, m_1 + 1}}{\Sigma} \cos^2 \theta - \frac{R}{\sqrt{2} k_2 (m_1^2 - 1) a^2} \left[(k_2 - 1)(m_1 - 1) \Delta_{k_2 + m_1 - 3, m_1 - 1} \right. \\ \left. - 2(m_1 - 1) \Delta_{k_2 + m_1 - 3, m_1 - 1} - (m_1 - 1)(m_1 - 2) \Delta_{k_2 + m_1 - 1, m_1 - 1} \right. \\ \left. + (m_1 - 1)(k_2 - 3) \Delta_{k_2 + m_1 - 1, m_1 + 1} + m_1 (m_1 - 1) \Delta_{k_2 + m_1 + 1, m_1 - 1} \right. \\ \left. + (m_1 - 1)(m_1 (k_2 - 1) + 1) \Delta_{k_2 + m_1 + 1, m_1 + 1} \right].$$

Deep superstrata

- BH microstates with GH bubbles - **very large J**
- Typically **~ 99%** Heidmann
- **85%** of maximal value Bena, Wang, Warner
- Impossible to lower by playing with GH bubbles
- Build deep superstrata: J can be arbitrarily small Bena, Giusto, Martinec Russo, Shigemori, Turton, Warner (PRL editor's selection)
- First BTZ microstates



Superstrata

Entropy:

- D1-D5 supertube - dimension of moduli space
 - classically: dimension = ∞
 - quantize: dimension = $4 N_1 N_5$ = number of momentum carriers
- Counting (+ fermions) (à la Maldacena Strominger Witten)
 $S = 2 \pi (N_1 N_5 N_p)^{1/2} !!!$ Bena, Shigemori, Warner

It remains to dot the i's and cross the t's :

- We have AdS-CFT duals. Solutions more and more messy as one approaches typical states (long strings). Recursive construction
- D1-D5 CFT - fractional momentum carriers. Have some, not all.
- Fluxes + warping: Small & Crumply \longrightarrow Big & Fluffy & Smooth
- Are typical microstates spanned by smooth solutions ?

MSW Superstrata

Bena, Martinec, Turton, Warner

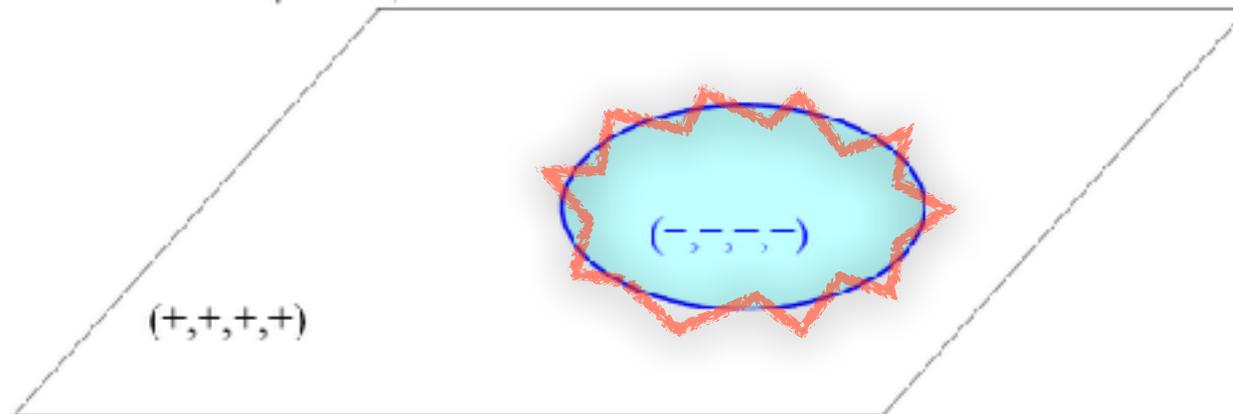
- D1-D5 solution: $AdS_3 \times S^3 \times T^4$
 - T-dualize on the Hopf fiber of S^3 + few more times
 - $AdS_3 \times S^2 \times T^6$: NS vacuum of the MSW CFT
- Central charges match
- subsector of MSW CFT \Leftrightarrow subsector of D1-D5 CFT !!!
- One arbitrary function worth of smooth solutions to $U(1)^4$ 5D ungauged supergravity

Why did we miss them solutions for past 12 years ?!?

Singular 4D ambipolar bases have one function worth of **singular fluxes** that gives rise to **smooth 5D solutions**

$$ds^2 = V (dx_1^2 + dx_2^2 + dx_3^2) + V^{-1}(d\psi + \vec{A})^2$$

$$V = \frac{1}{r} - \frac{Q}{|\vec{r} - \vec{a}|} + \frac{Q}{|\vec{r} - \vec{b}|}$$



- Signature of base changes from $(+, +, +, +)$ to $(-, -, -, -)$
- Z_i blow up and change sign at interface:

$$d * d Z_1 = G^2 \wedge G^3 \quad \Rightarrow \quad Z_i \sim \frac{1}{V}(\dots)$$

- Full 11D metric is smooth:

$$ds^2 = -Z^{-2}(dt + \vec{k})^2 + Z \left[V (dx_1^2 + dx_2^2 + dx_3^2) + V^{-1}(d\psi + \vec{A})^2 \right] + ds_{1,6}^2$$

Extra singular wiggly G^i sourced at the interface

SUSY microstates – the story:

- We have a huge number of them
 - Arbitrary continuous functions of 2 variables
 - Smooth solutions. 4 scales !
 - Superstrata reproduce black hole entropy ☺
Bena, Shigemori, Warner
- Dual to CFT states in typical sector
 - This is where BH states live too ☺
 - AdS-CFT perspective: highly weird if BH microstates had horizon
Bena, Wang, Warner; Taylor, Skenderis
- Two non-backreacted calculations:
 - BH entropy - scaling multicenter config ☺
Denef, Moore; Denef, Gaiotto, Strominger, Van den Bleeken, Yin
 - Higgs-Coulomb map.
Bena, Berkooz, de Boer, El Showk, Van den Bleeken; Lee, Wang, Yi

Quantum Gravity in AdS₂

- Everybody & their brother & SYK
- AdS₂ - no finite-energy excitations
Maldacena, Strominger
- backreaction of particle in AdS₂ either
 - destroys UV
 - singularity in IR
(? ↔ SYK 4-pt. function **not** conformally invariant)
- Singularities in String Theory and AdS-CFT solved by **string and brane dynamics** involving **extra dimensions** 20 years of examples

Quantum Gravity in AdS₂

- Typical microstate geometries have **long AdS₂ throat**

- Limit when length $\rightarrow \infty$

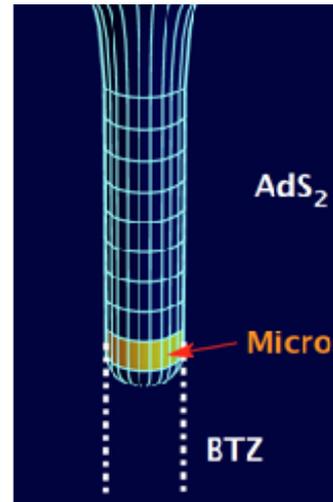
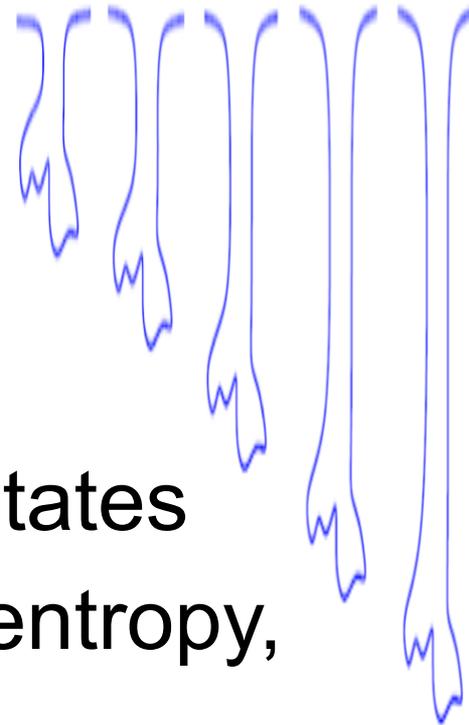
- Solutions above \rightarrow **asymptotically-AdS₂**

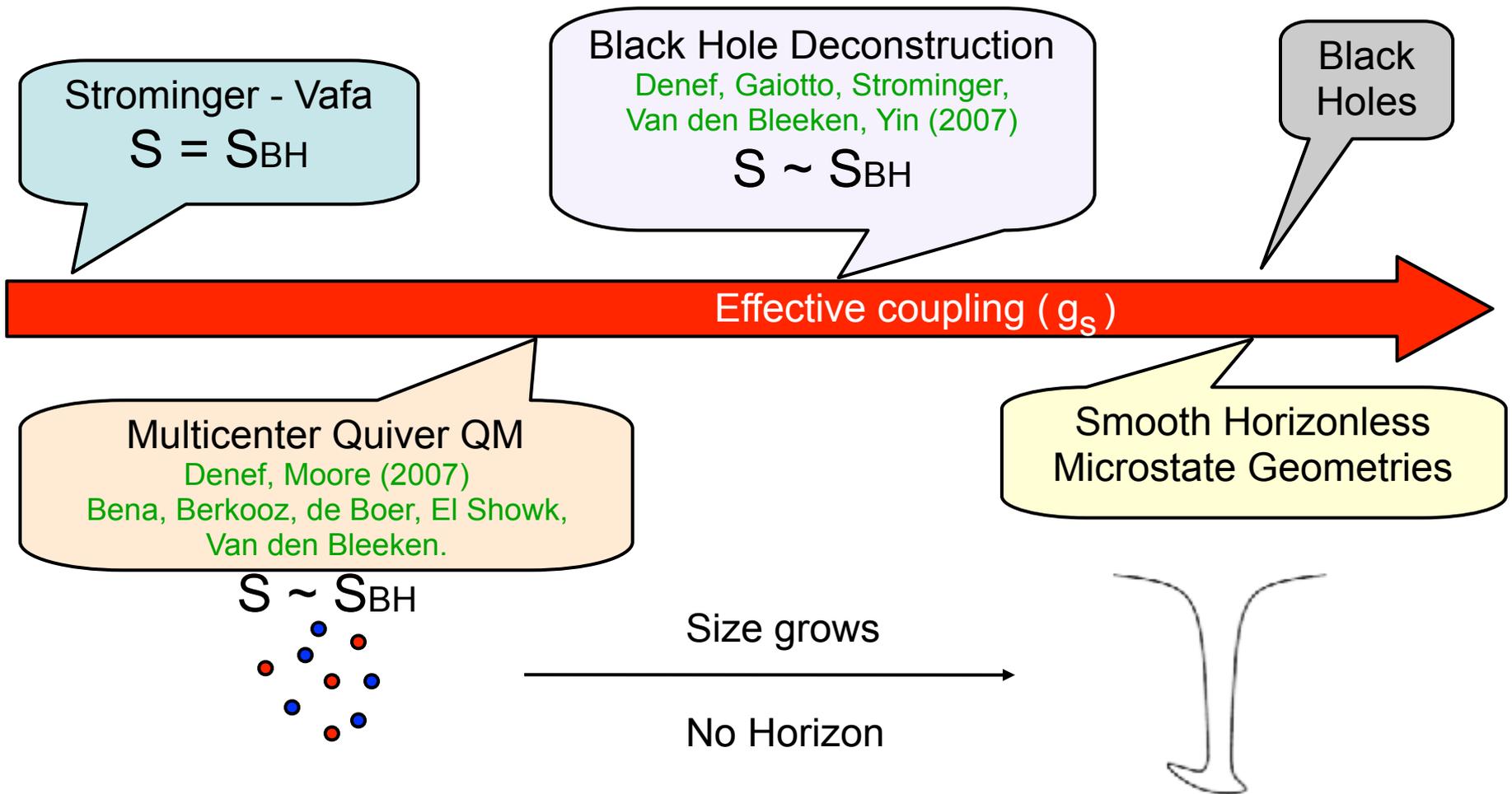
Bena, Heidmann, Turton

- **Same entropy** as microstates

- If superstrata count BH entropy, so do these solutions !

- Ground states of QM dual to AdS₂ Sen



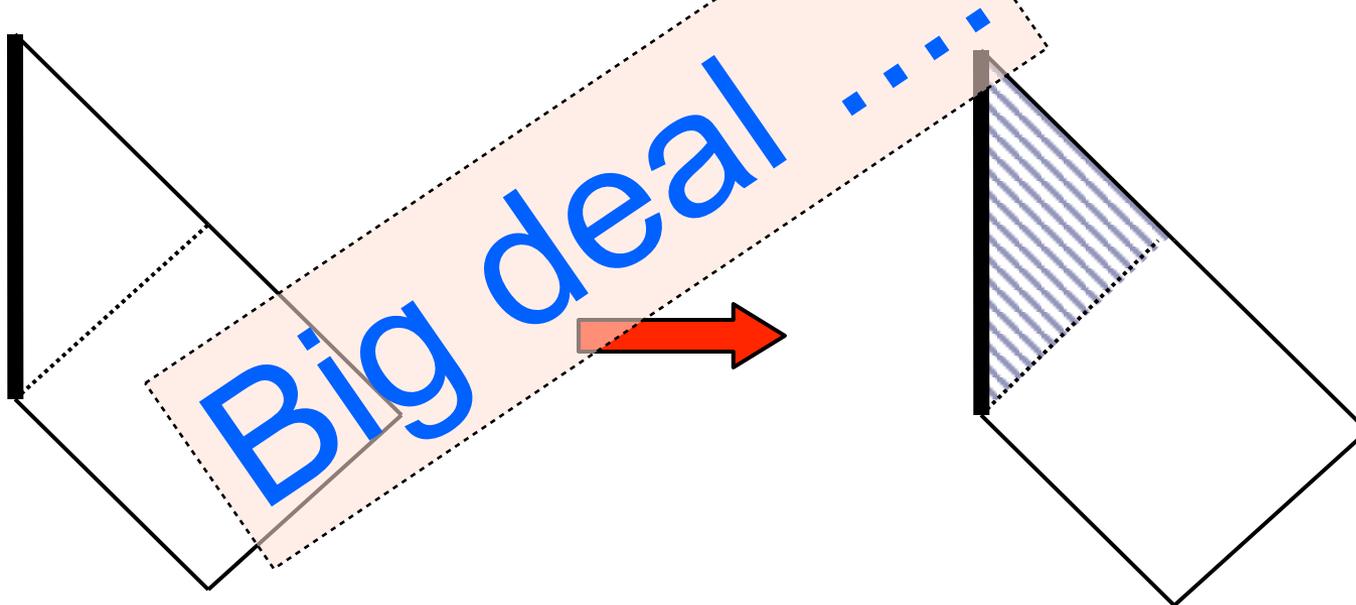


Punchline: Typical states **grow** as G_N increases.
 Horizon never forms.
 Quantum effects from singularity **extend to horizon**

Similar story for **non-SUSY extremal black holes**

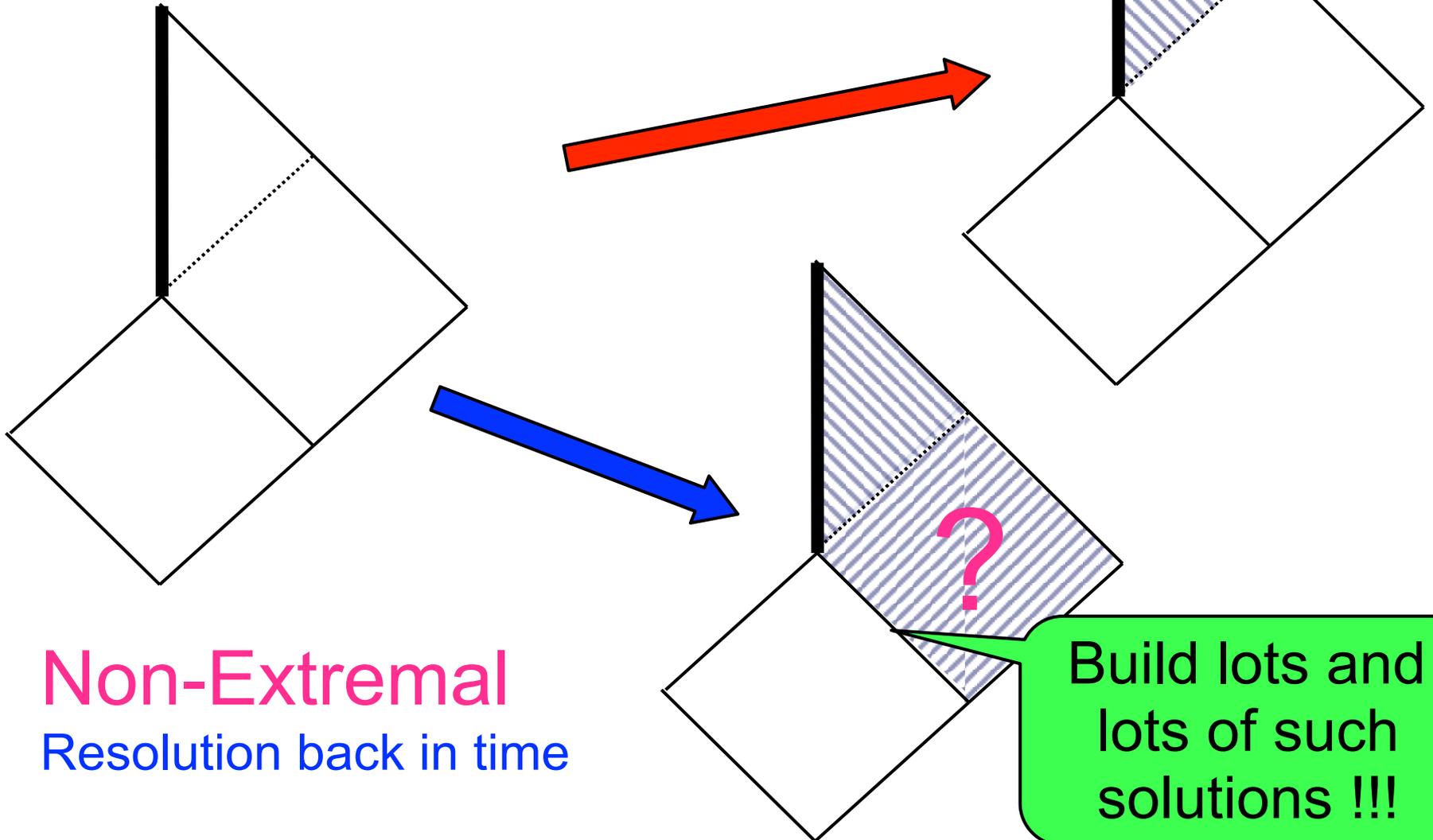
BPS Black Hole = Extremal

- This is **not so strange**
- Horizon **in causal future** of singularity
- **Time-like singularity** resolved by (stringy) low-mass modes extending to horizon



Penrose
Poisson, Israel
Dafermos
Marolf

The really big deal fuzzball, firewall



Non-Extremal
Resolution back in time

Build lots and
lots of such
solutions !!!

Very few known. Extremely hard to build...

– Coupled nonlinear 2'nd order PDE's do not factorize

Do not pray to the saint who
does not help you !

Romanian proverb

- Idea: perturbative construction - near-BPS
- Add antibranes to BPS bubbling sols.
Kachru, Pearson, Verlinde
- Metastable minima
Bena, Puhm, Vernocke
- Decay to susy minima:
brane-flux annihilation - Hawking radiation
- Microstates of near-extremal BH

Very few known. Extremely hard to build...

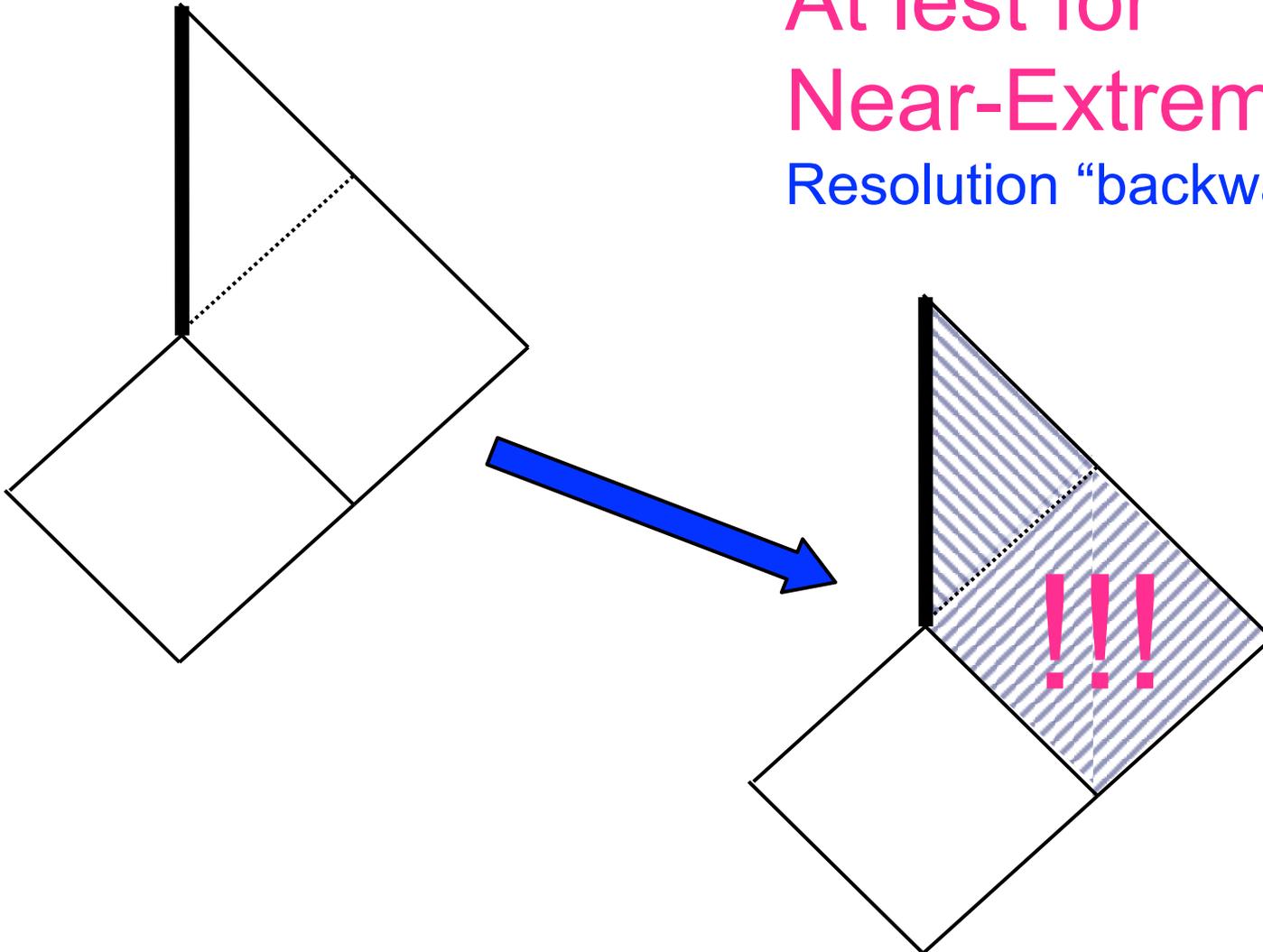
– Coupled nonlinear 2nd order PDE's do not factorize

When a bird is blind, God sometimes makes its nest ! another Romanian proverb

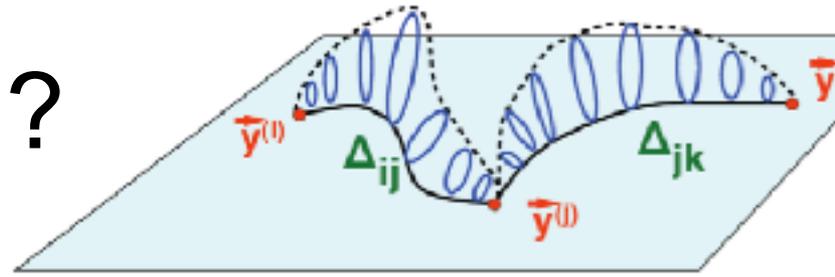
- For some solutions the 2nd order PDE's do factorize !!! Bossard, Katmadas
- We can build analytically certain classes of non-extremal solutions Bena, Bossard, Katmadas, Turton
- Add extra cycles to JMART
- Method can get us far from extremality.
- How far ? How generic ? Antibranes ?

The really big deal

At least for
Near-Extremal
Resolution “backwards in time!”

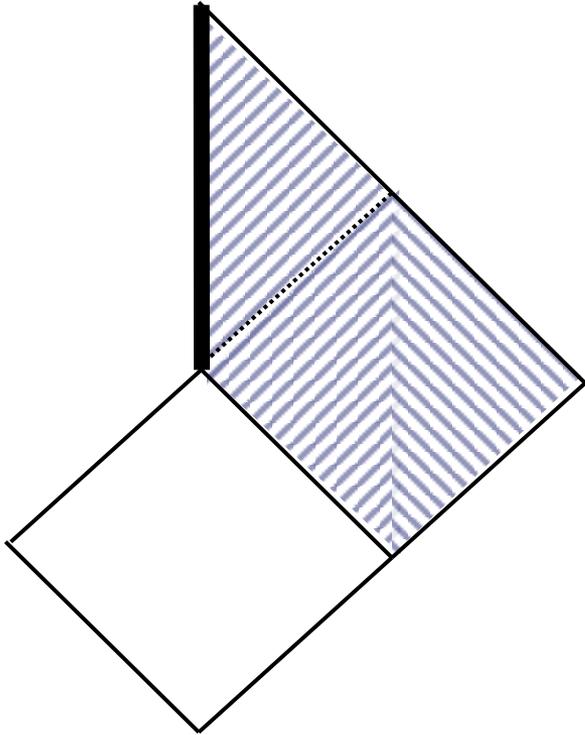


Why not collapsing ?



- 5(+6)d : smooth solutions + **quantized** magnetic flux on topologically-nontrivial **2-cycles**
 - cycles smaller \rightarrow increases energy
 - bubbling = **only** mechanism to avoid collapse in semiclassical limit Gibbons, Warner
 - If **any** state in the e^S -dimensional BH Hilbert space has a semiclassical limit, it **must** be a microstate geometry !
- 4(+6)d : multicenter solutions Denef
 - smooth GH centers with negative charge \rightarrow centers with **negative D6 charge** and **negative mass**
 - common in String Theory (e.g. orientifolds)
 - **Highly unusual** matter from a 4d perspective
 - Usual matter does not hang around, just falls in BH

What about other black holes?



- Near Extremal ?
- Schwarzschild + 1 electron ?

Take electron away

Same Penrose diagram !

String theory **can** resolve BH singularities

“backwards in time.” **Why stop at near-extremal?**

Same Mechanism ?

Pure BH states have no horizon - 4 approaches:

(1) Information-theory arguments

Mathur 2009, AMPS, etc

- secondary question: firewall ? burn or spit through ?

(2) Generic AdS-CFT

Agnostic about theory

No mechanism for Hair !

Taylor, AMPS2 (Papadodimas Raju against)

- nontrivial \Rightarrow no spherical symmetry \Rightarrow no horizon

(3) Follow microstates from weak to strong coupling

- BH deconstruction, String emission, Higgs-Coulomb map

Denef, Gaiotto, Strominger, Van den Bleeken, Yin, Giusto, Russo, Turton
Bena, Berkooz, de Boer, El Showk, Van den Bleeken; Lee, Wang, Yi

(4) Lots of BH microstate geometries = Hair !!!

- One mechanism in three hypostases:

Bubbling \Leftrightarrow Brane polarization \Leftrightarrow NonAbelian

- Can get BH entropy; 2 new scales, E_{gap} , λ_T

A few questions

- **Would all microstates be classical ?**

- No, but classical solutions are the only things one can construct
- **Hovering mechanism extrapolates** \Rightarrow brane polarization, non-Abelian
- Typical states: many small bubbles ($\lambda_T \sim \ell_P$), or just a few ($\lambda_T > \ell_P$)
- Larger bubbles - more entropy Denef, Moore; Bena, Shigemori, Warner

- **Don't people in Saclay say antibranes are bad?**

- **Tachyonic !** Bad for cosmology, **but not for BH !**
- Instabilities in fact **expected** for non-extremal black hole microstates; **JMaRT (+ bubbles)** has them Myers&al, Santos&al
- D1-D5: **BPS left-movers** + **right movers** Mathur

- **What about non-linear instabilities ?**

- first-order backreaction of non-BPS perturbation; D1-D5 **right movers** \Rightarrow Closed string emission
- **Moduli space of classical solutions.** non-BPS \Rightarrow Motion Bena, Pasini Marolf, Michel, Puhm

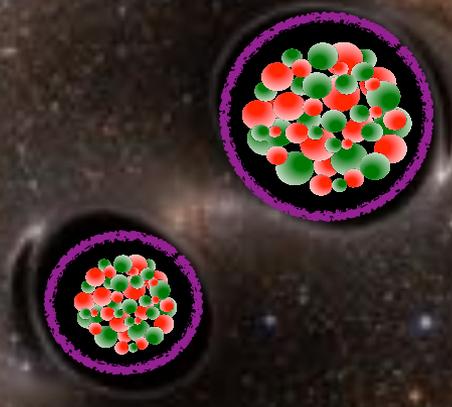
A few questions

- Can you fall through horizon **drinking your coffee** ? (as GR textbooks say)
- Do you rather go **splat** at the horizon scale?
- 3 options:
 - Analyze ∞ **density shells** / membranes / stuff carrying d.o.f. @ horizon (kept from collapsing by the **Tooth Fairy**)
 - Modify gravity by **weird terms** and analyze horizon
 - Use actual solutions of String Theory
- Answer likely depends on E_{gap} , λ_T
- **Known bubbling solutions** or **polarized branes** have no intention to let you fall through unharmed

How can we observe this ?

Universal feature:

- Low-mass degrees of freedom at horizon.



Collision of two black holes:

Gravitational waves emitted close to horizon:

LIGO, eLISA

Summary and Future Directions

- String theory configurations that **hover above horizon**.
Topology + fluxes \Leftrightarrow **brane polarization** \Leftrightarrow **nonabelian d.o.f.**
- **BPS black hole microstates** = horizonless solitons
 - **low-mass modes** affect **large (horizon) scales**
 - Convergence of many research directions
 - BPS **superstrata** - 2 variables - **Black Hole Entropy !**
- Extensive extremal non-BPS story
- Extend to **non-extremal** black holes
 - **Near**-extremal
 - Metastable supertubes Bena, Puhm, Vernocke
 - **Far** from extremality — 2'nd order nonlinear coupled PDE
 - **Systematic construction** Bena, Bossard, Katmadas, Turton
 - Others: numerics? inverse scattering? blackfolds?
 - Maybe start thinking about **experimental** consequences ?
 - Gravity waves
 - Supermassive BH formation easier

Some speculative connections

- **A.** 10-yr old question: what is the dual of **pure Higgs states** ?
- Martinec: W-branes - pure Higgs entropy from condensing M2 branes wrapping 2-cycles in GH space (F1 between fluxed D6 in 10D)
- Similar to D0-D4: bi-fundamentals come from F1 between D0 and D4
 - F1's source fields in hypermultiplets of sugra.
 - Long time belief: need sugra solutions with hypermultiplets
Ortin, Raymaekers, Van den Bleeken
 - Think deeper: hypermultiplets = **red herring**
- String emission calculations - first order in operators that correspond to going on the Higgs branch
- Going on the Higgs branch turns on **(1,1)** metric components on the T^6 . Same from four-charge system **Bianchi, Morales, Pieri**
- Makes sense - condensation of F1 between 2 D2's bend them into each other. Source extra **(1,1)** components

Some speculative connections

- **B.** MSW entropy counting:
- N_1, N_2, N_3 M5 wrapping three T^4 's inside T^6 . Singular ample divisor.
- Smooth ample divisor = deformation into single M5 brane of length $N_1 \times N_2 \times N_3$; sources $(1,1)$ metric components. Expects them to be present in generic microstate
- **C.** String emission - extra field $(1,1)$ metric on T^6 Giusto, Russo, Turton
- **D.** Smoothness of superstrata - coiffuring - same field
- **E.** Function worth of MSW microstate solutions - same field
- **Five** different indications we are converging on the right ingredient.

