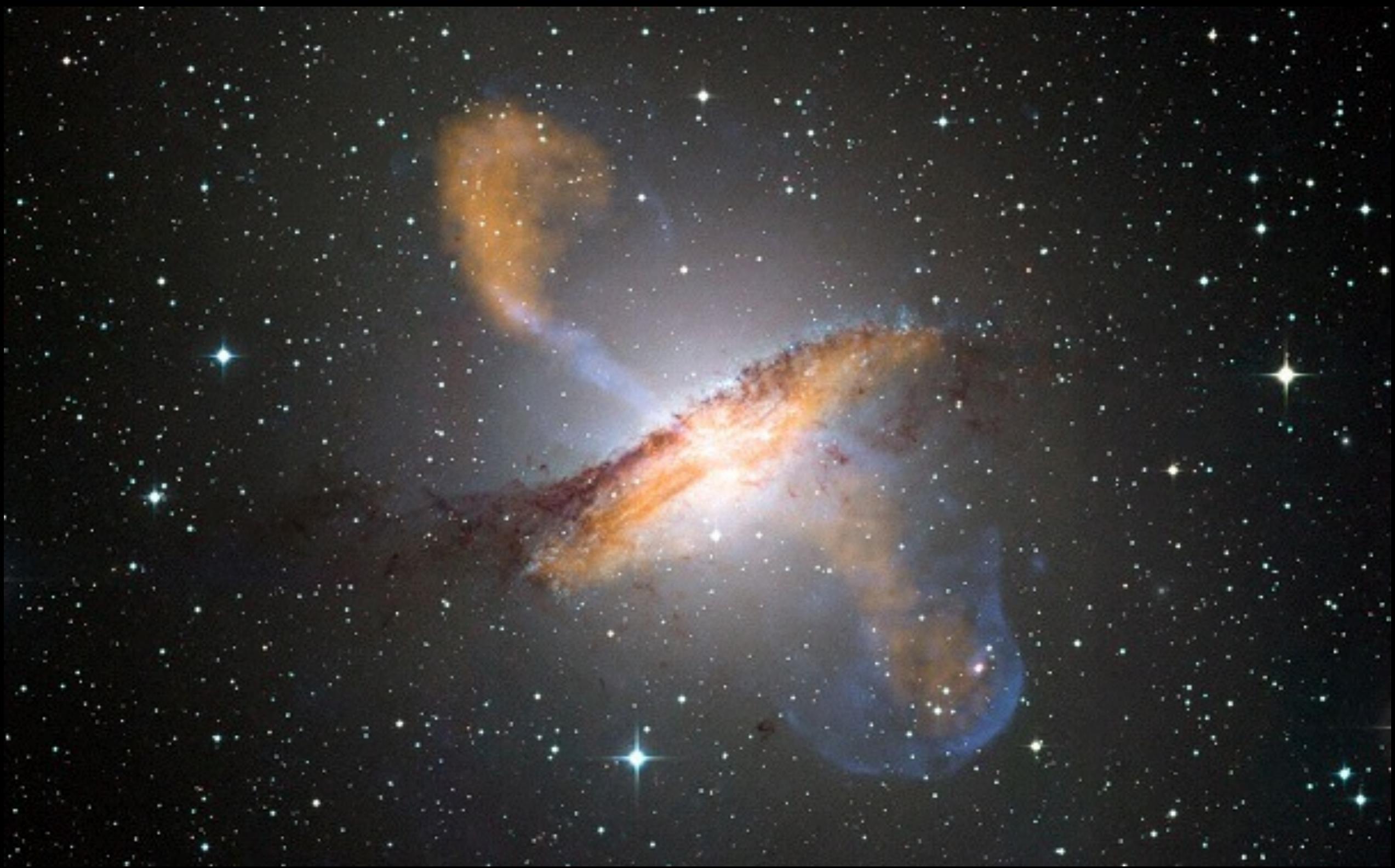
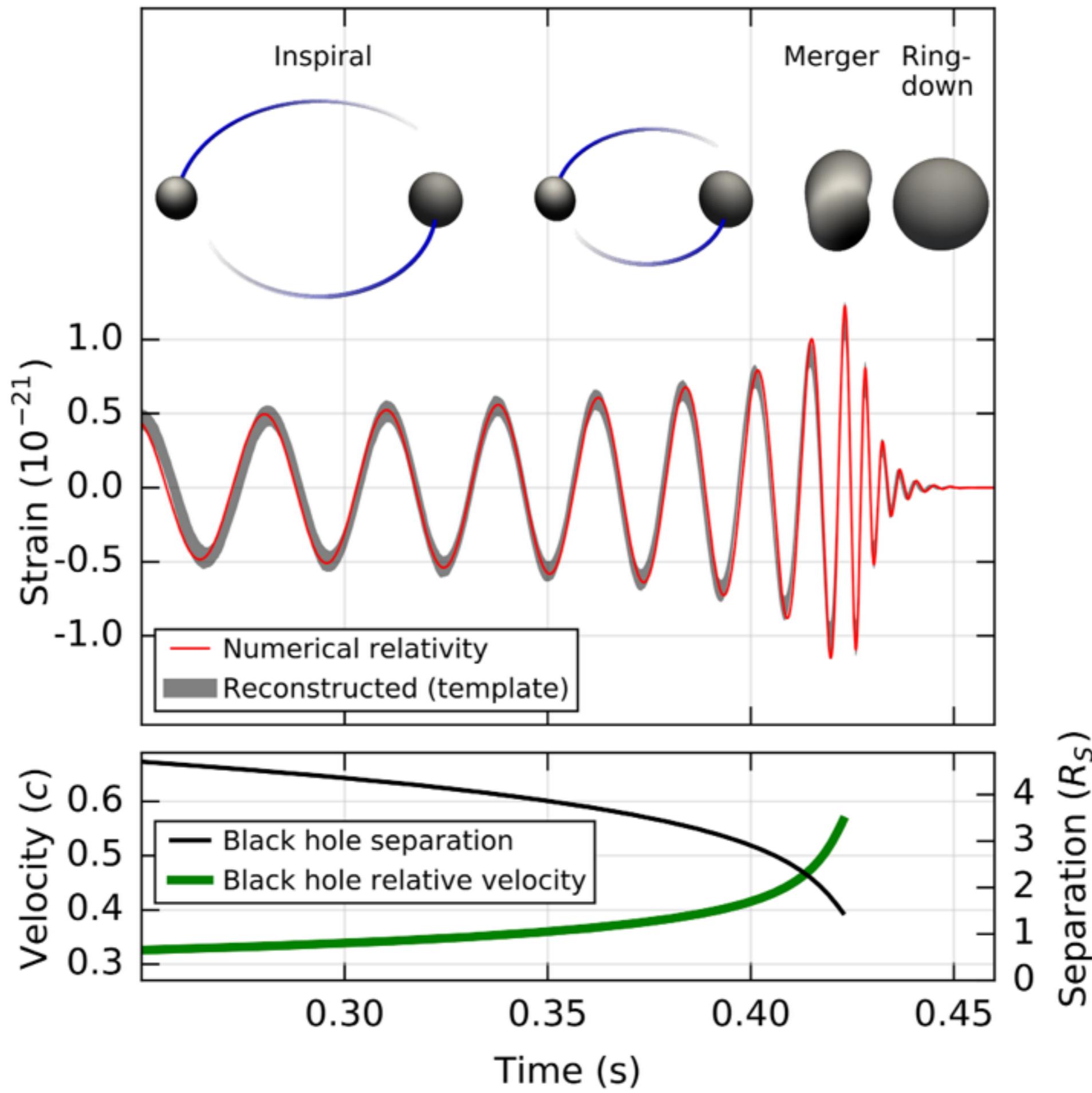


Black holes,
quantum information, entanglement,
... and all that

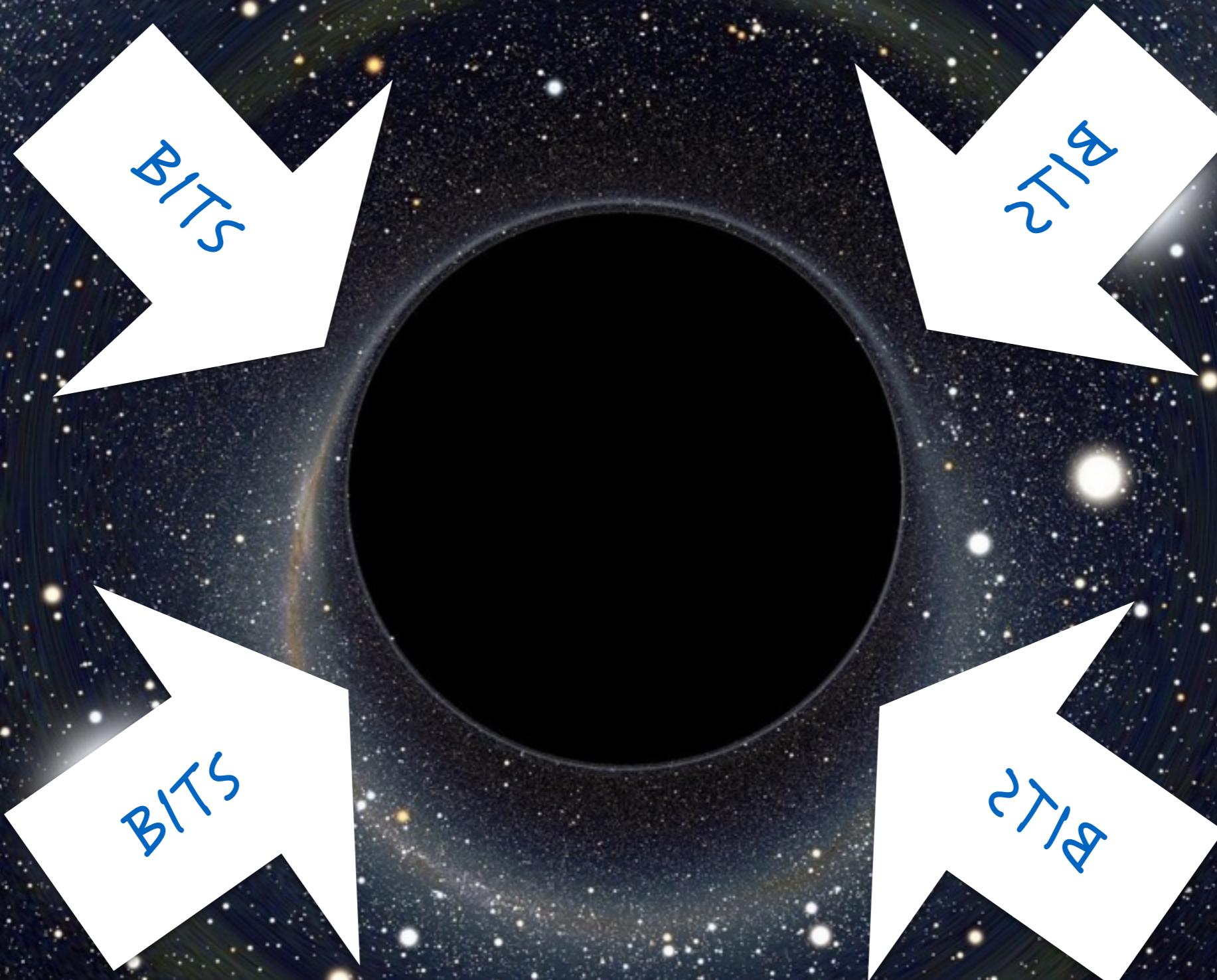
J.L.F. Barbón





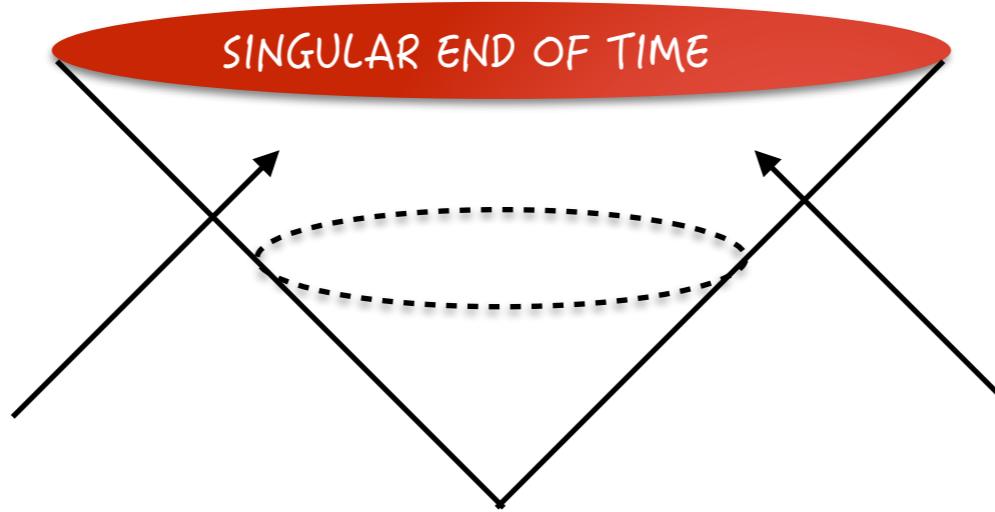


BLACK HOLES AND INFORMATION



BLACK HOLES AND CLASSICAL INFORMATION

Classically, the hole eats and destroys information



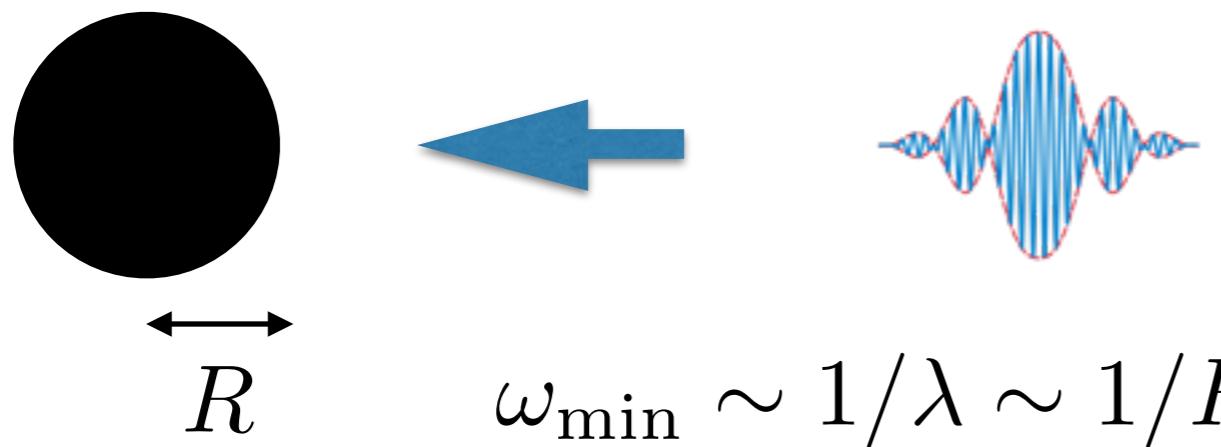
No processing. Singularity an infinite sink, so
the amount of hidden information could be infinite,

i.e.

$$S_{\text{classical}} = \infty$$

BLACK HOLES AND QUANTUM INFORMATION

BEKENSTEIN argued for a FINITE capacity because of quantum mechanics



$$\delta S_{\min} \sim \log 2 \sim \delta R^2/G$$

$$\omega_{\min} \sim 1/\lambda \sim 1/R$$

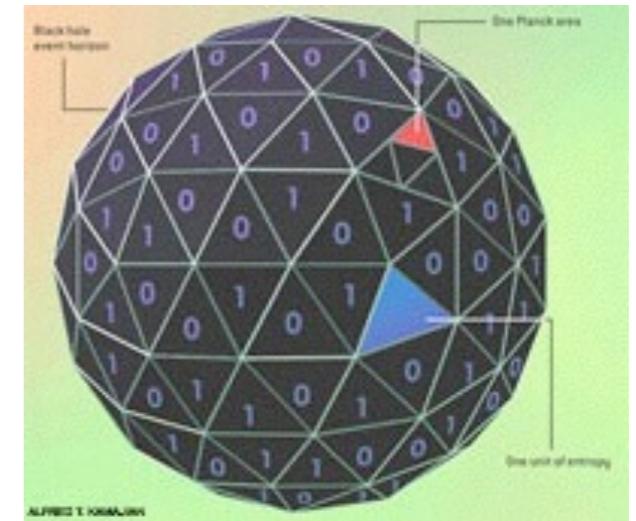
$$S \sim \frac{c^3}{G\hbar} \text{ (area horizon)}$$

STRIKING FACTS

d.o.f. not extensive in 3-space!

THE BEGINNING OF HOLOGRAPHY!

$$\log \dim \mathcal{H}_{\text{bh}} \sim S \sim \frac{\text{Area}}{\ell_P^2}$$



density of states REALLY strong (more than any QFT or string theory)

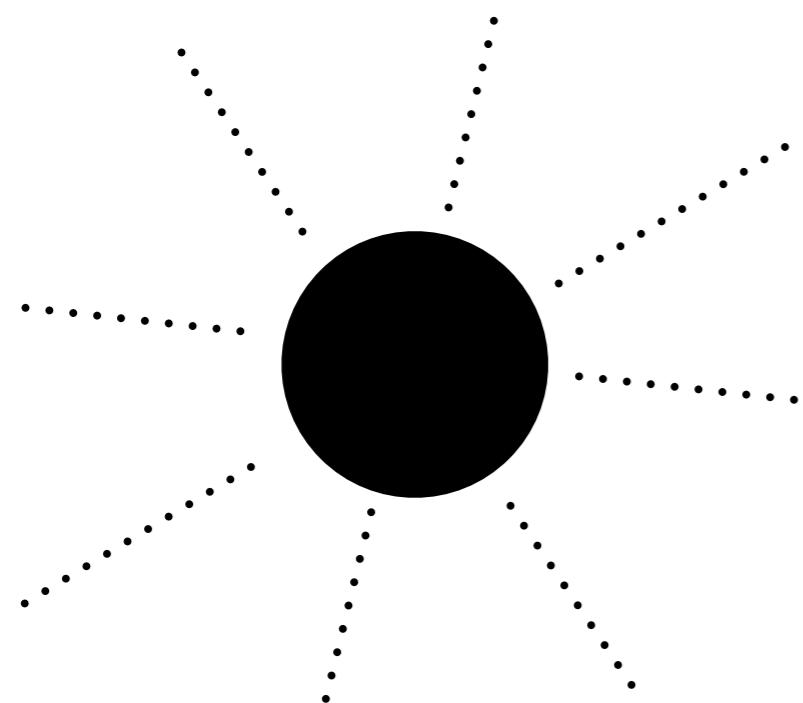
$$A_H \sim R^2 \sim G^2 M^2$$

$$\Omega(M)_{\text{bh band}} \sim \exp(CGM^2)$$

But ... having an entropy, we must have a TEMPERATURE

$$T = \frac{\partial M}{\partial S} \sim \frac{1}{GM} \sim \frac{1}{R}$$

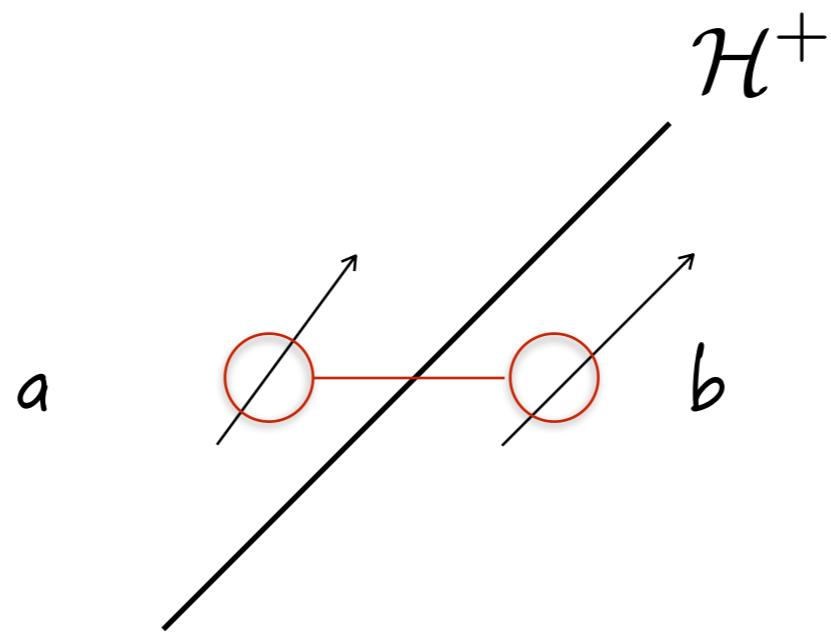
Famously found by Hawking with an entirely different argument



$$T = \frac{1}{4\pi R} = \frac{1}{8\pi GM}$$

$$S = \frac{A_H}{4G}$$

Extra assumption: horizon smoothness \rightarrow entanglement

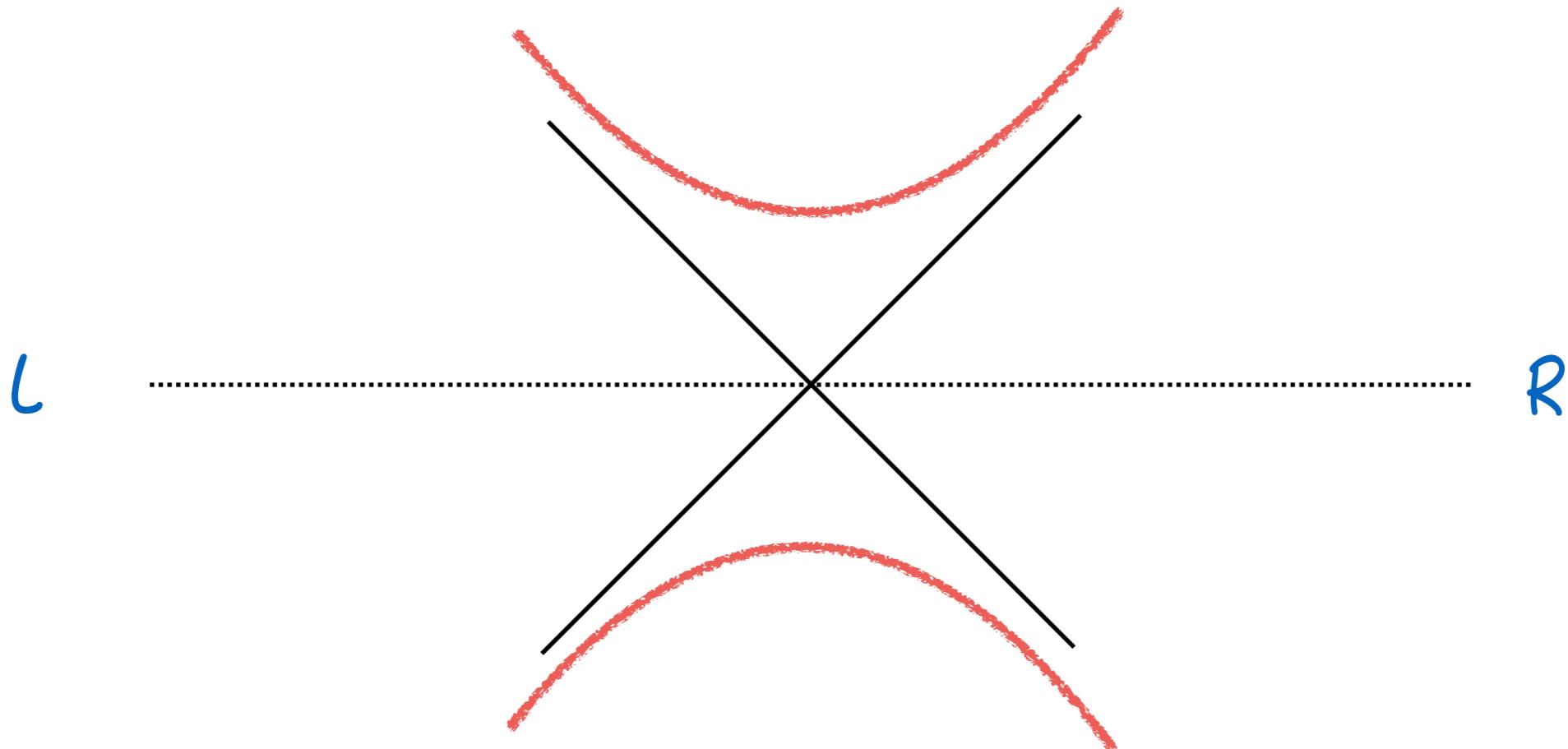


$$|a, b\rangle \propto \sum_{\omega} e^{-\omega/2T} |-\omega\rangle_a \otimes |\omega\rangle_b$$

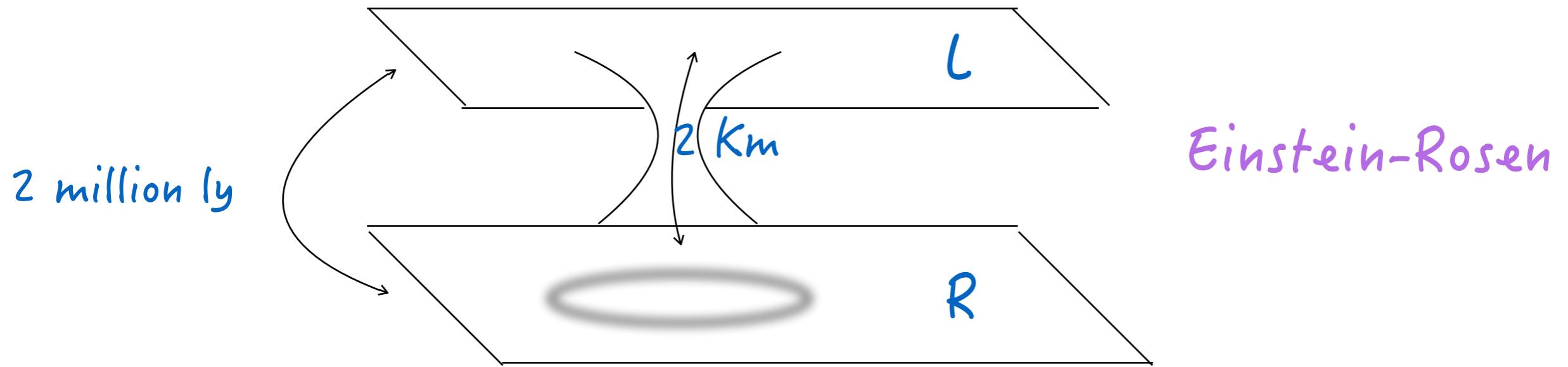
Unruh, ...

$$\rho_b \propto \sum_{\omega} |\omega\rangle_b e^{-\omega/T} {}_b\langle\omega|$$

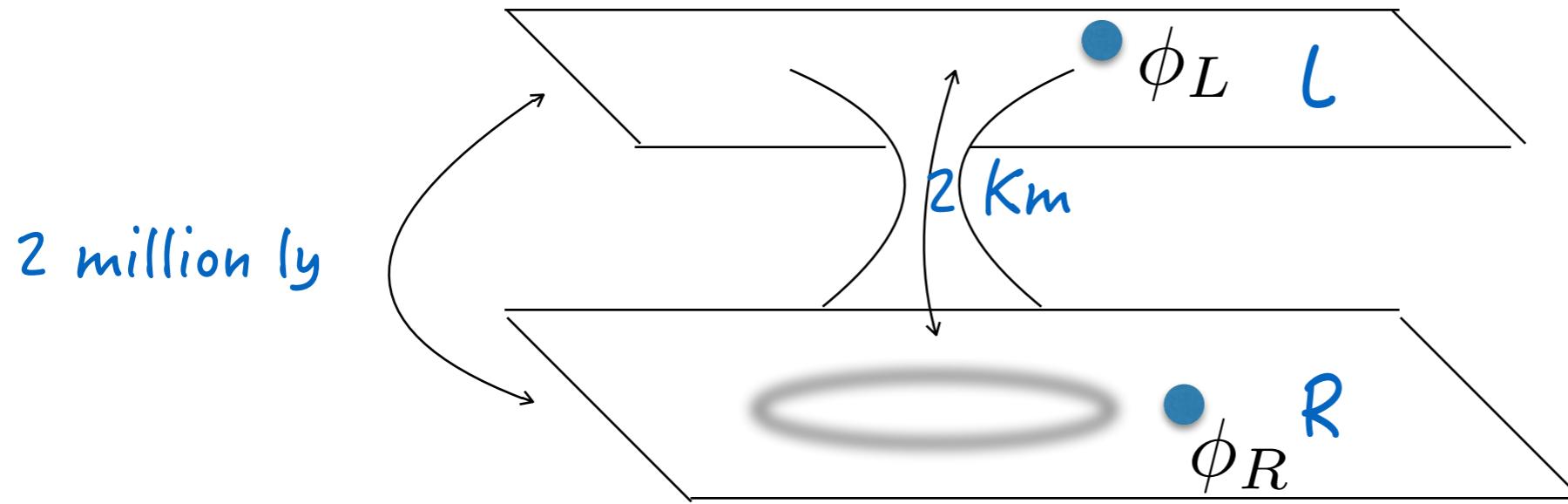
The most symmetrical case: the maximal analytic extension, or two black holes sharing the interior



the time-symmetrical slice contains a wormhole !



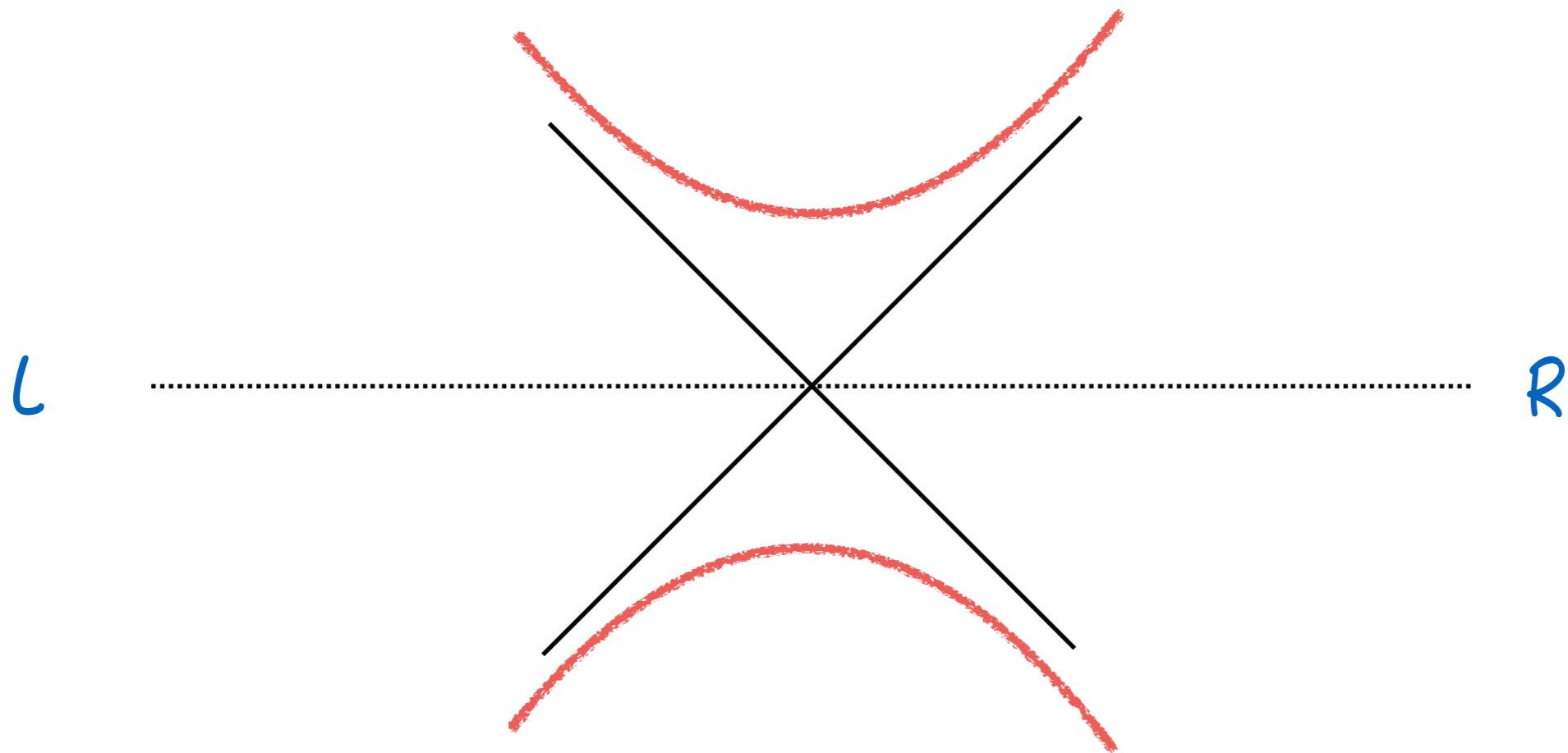
wormhole detection !



$$\langle \phi_L \phi_R \rangle \sim \frac{1}{(2 \text{ km})^2}$$

despite the fact that they are
separated by 2 million l.y. ... or more,
along the "outside route"

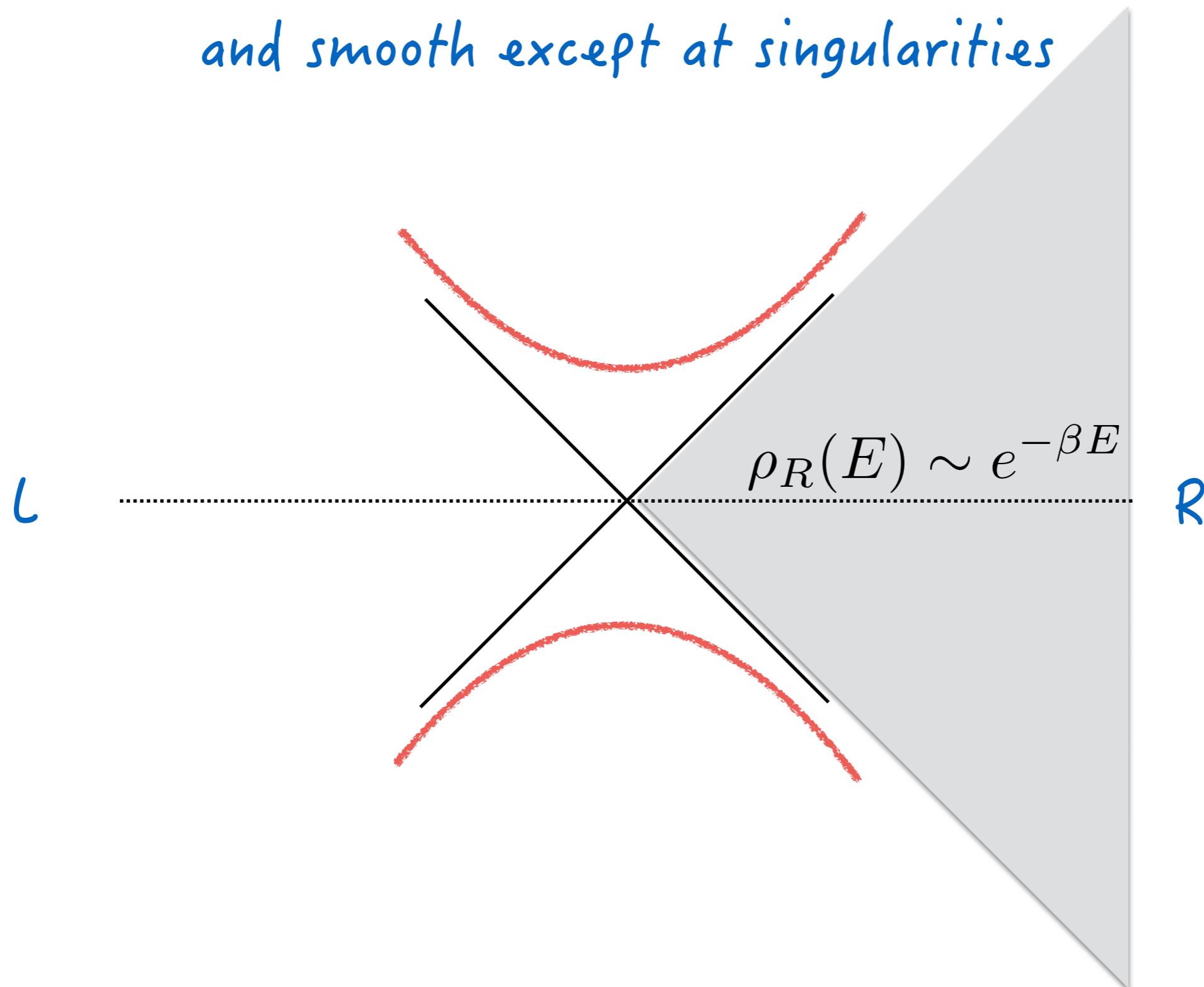
Global state of quantum fields (Hartle-Hawking) is pure and smooth except at singularities



$$|\text{HH}\rangle \sim \int_E e^{-\beta E/2} |E\rangle_L \otimes |E\rangle_R$$

Israel

Global state of quantum fields (Hartle-Hawking) is pure and smooth except at singularities



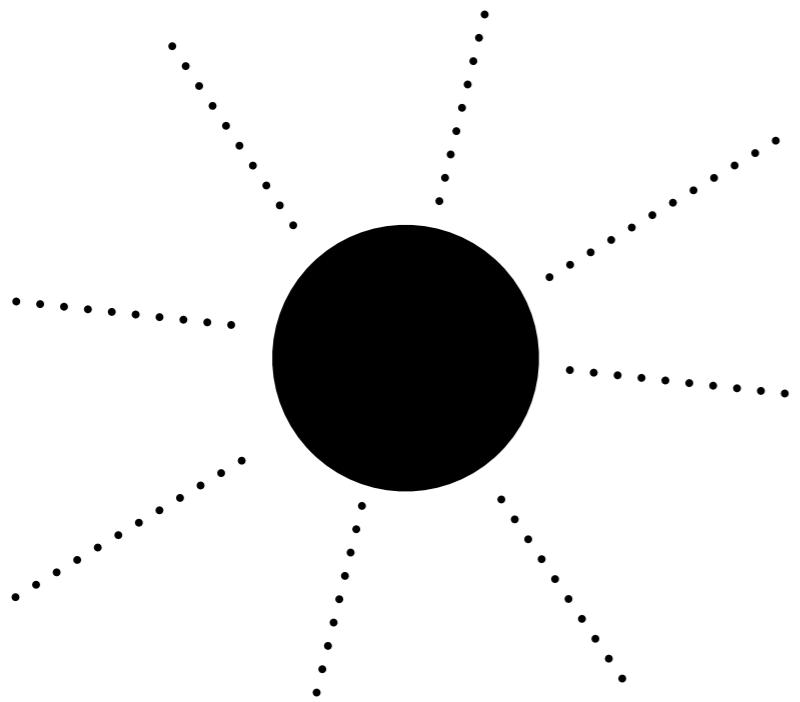
Spectrum of one-sided energies is continuous

BLACK HOLES AND QUANTUM ENTANGLEMENT

Horizon smoothness guarantees large amounts of entanglement in the local state of the near-horizon region

Whether the converse is also true is subject of current heated debate

Evaporation phenomenology



$$\begin{aligned}\omega_q &\sim \frac{1}{R} \\ \lambda_q &\sim R \\ \tau_q &\sim R\end{aligned}$$

$$\#_q \sim \frac{M}{\omega_q} \sim M R \sim S$$

$$\tau_{\text{ev}} \sim \tau_q \cdot \#_q \sim R S$$

So, a black hole is a long-lived “resonance” with a very narrow width and a huge density of states

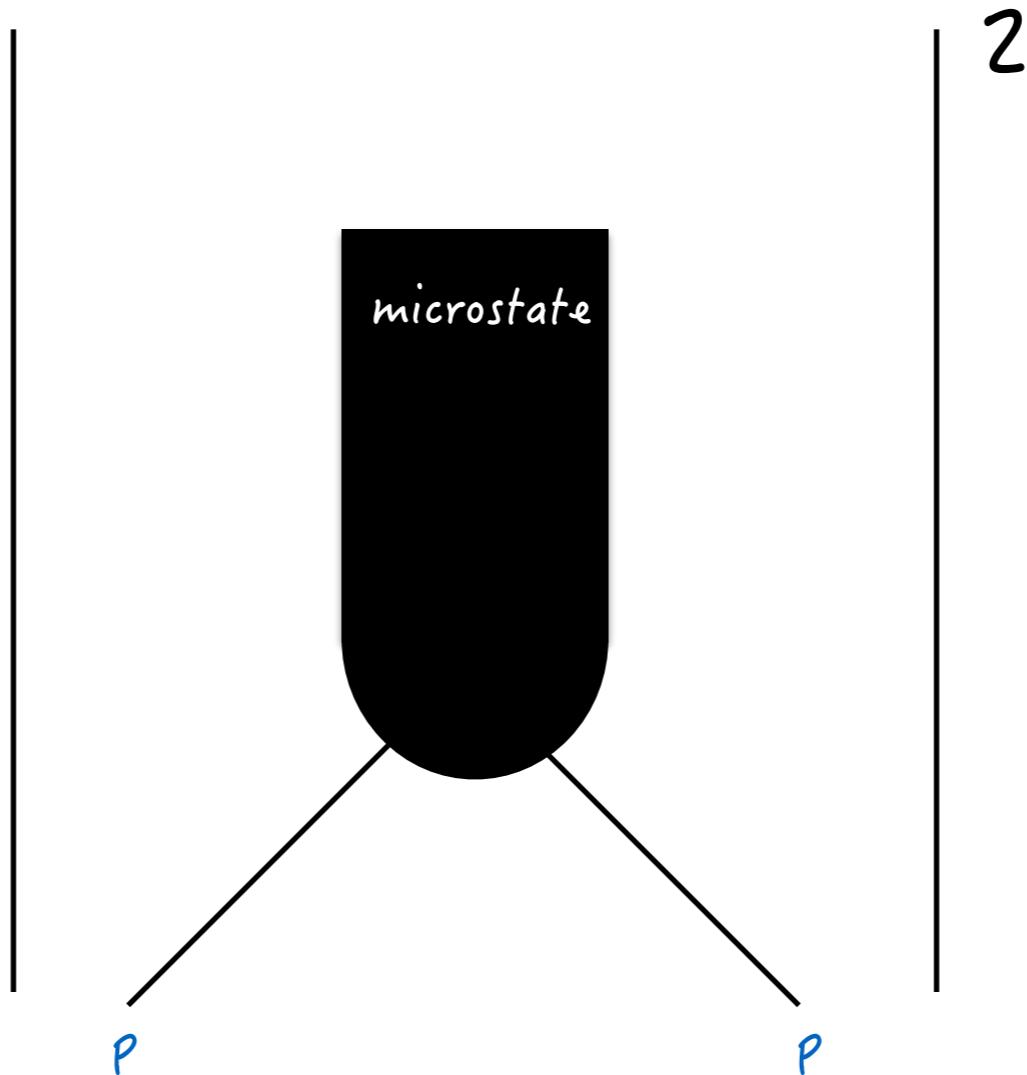
$$\Gamma \sim \frac{1}{\tau_{\text{ev}}} \sim \frac{1}{RS} \ll M$$

$$\Omega(M) \propto e^S \sim \exp(4\pi^2 GM^2)$$

It is natural to treat the formation/evaporation as any other S-matrix process

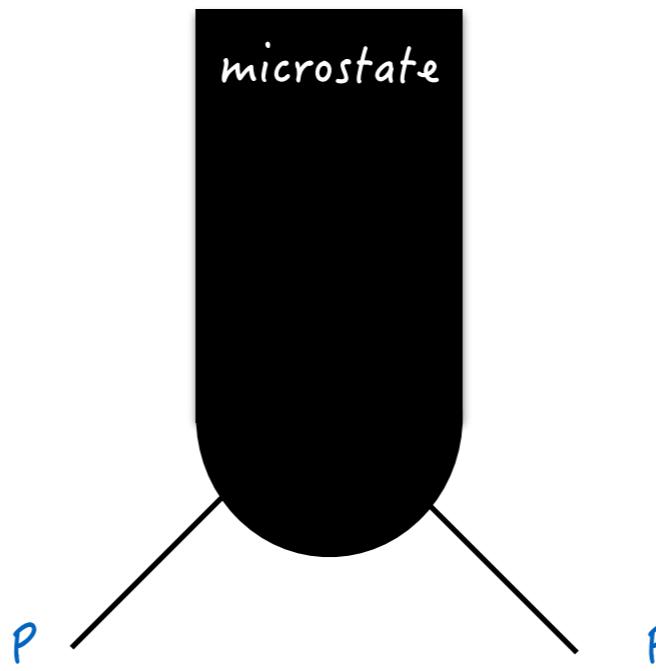
Constraints from UNITARITY

$$\sum e^S \text{ states}$$



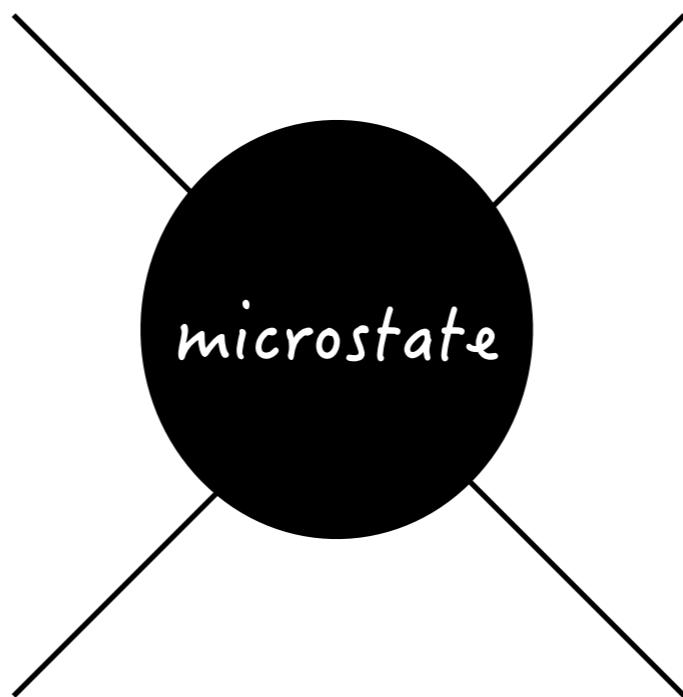
$$= \mathcal{O}(1)$$

proton-proton-bh microstate
vertex



$$\sim e^{-S/2}$$

$2 \rightarrow 2$ scattering
amplitude



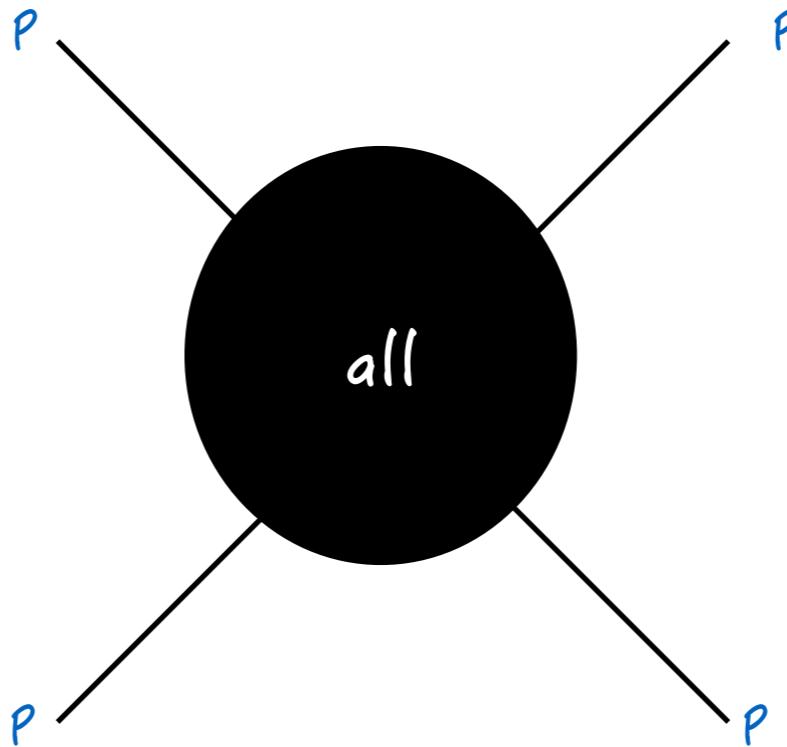
$$\sim e^{-S}$$

proton-proton-bh microstate
vertex



$$\sim e^{-S/2}$$

$2 \rightarrow 2$ scattering
amplitude



$$\sim e^{-S/2}$$

Scattering amplitudes of order e^{-M^2/M_P^2}
are much softer than anything QFT
or even string theory can produce

Both the very strong density of states and
the extremely soft high energy scattering
suggest that the fundamental theory of quantum
gravity is not local ... holographic?

In what follows, we will try to "follow the money"
in the form of these elusive e^{-S} effects

Any effect of order ℓ^{-S} is non-perturbative in gravity and beyond low-energy effective field theory

The effective expansion parameter for graviton loops is

$$\alpha_{\text{eff}} = G \omega_q^2 \sim \frac{G}{R^2} \sim \frac{1}{S}$$

$$e^{-S} \sim e^{-1/\alpha_{\text{eff}}}$$

... but such effects are precisely those we need to control to avoid the infamous "information paradox"

In Hawking's picture, each rad quantum is well approximated by a member of an EPR pair: radiation comes out in a highly mixed state

$$\Delta S_{\text{rad}} \Big|_{\text{emission}} \sim \log 2 = O(1)$$

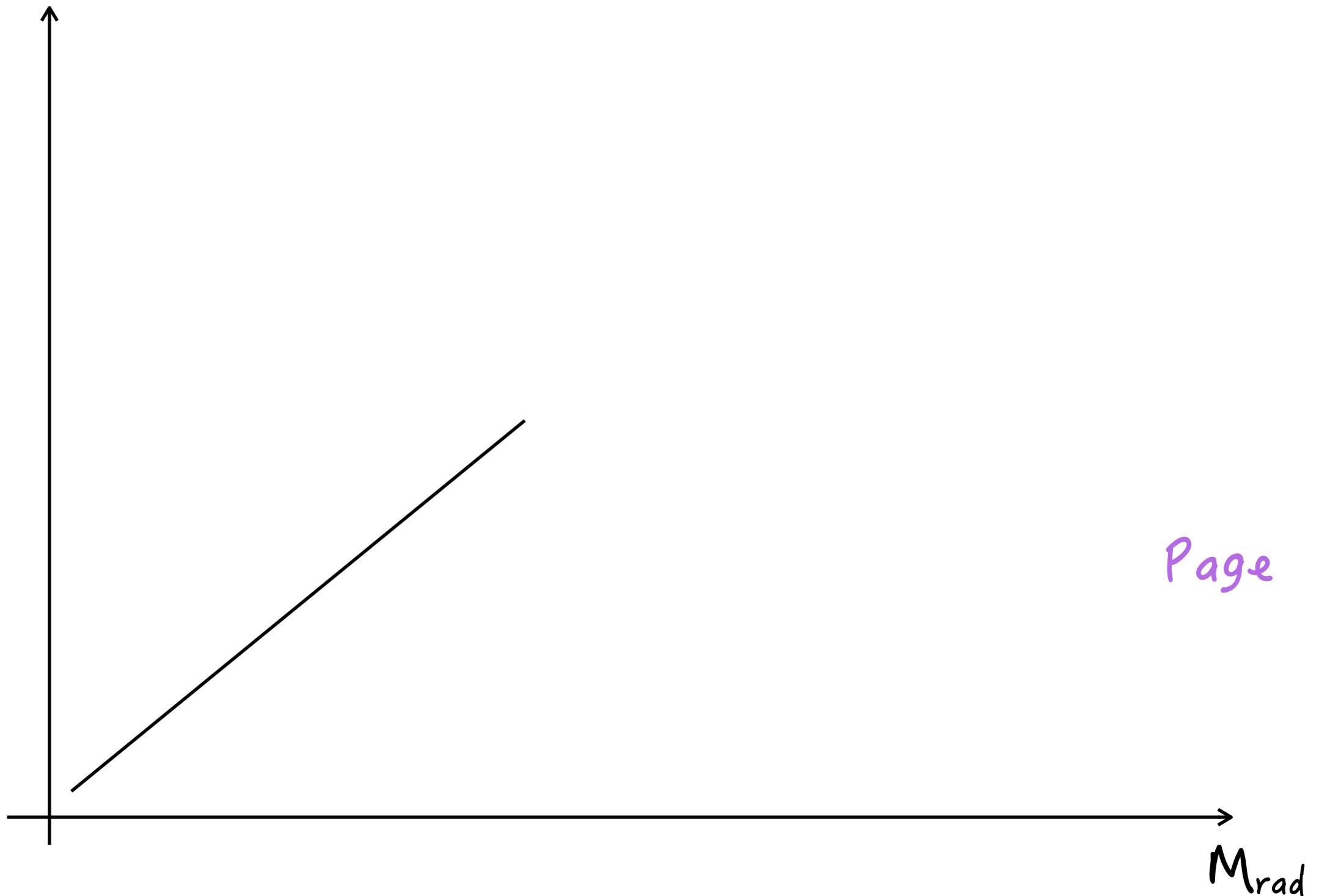
entropy production is roughly maximal

$$S_{\text{rad}} \sim \#_q \sim R M_{\text{rad}}$$

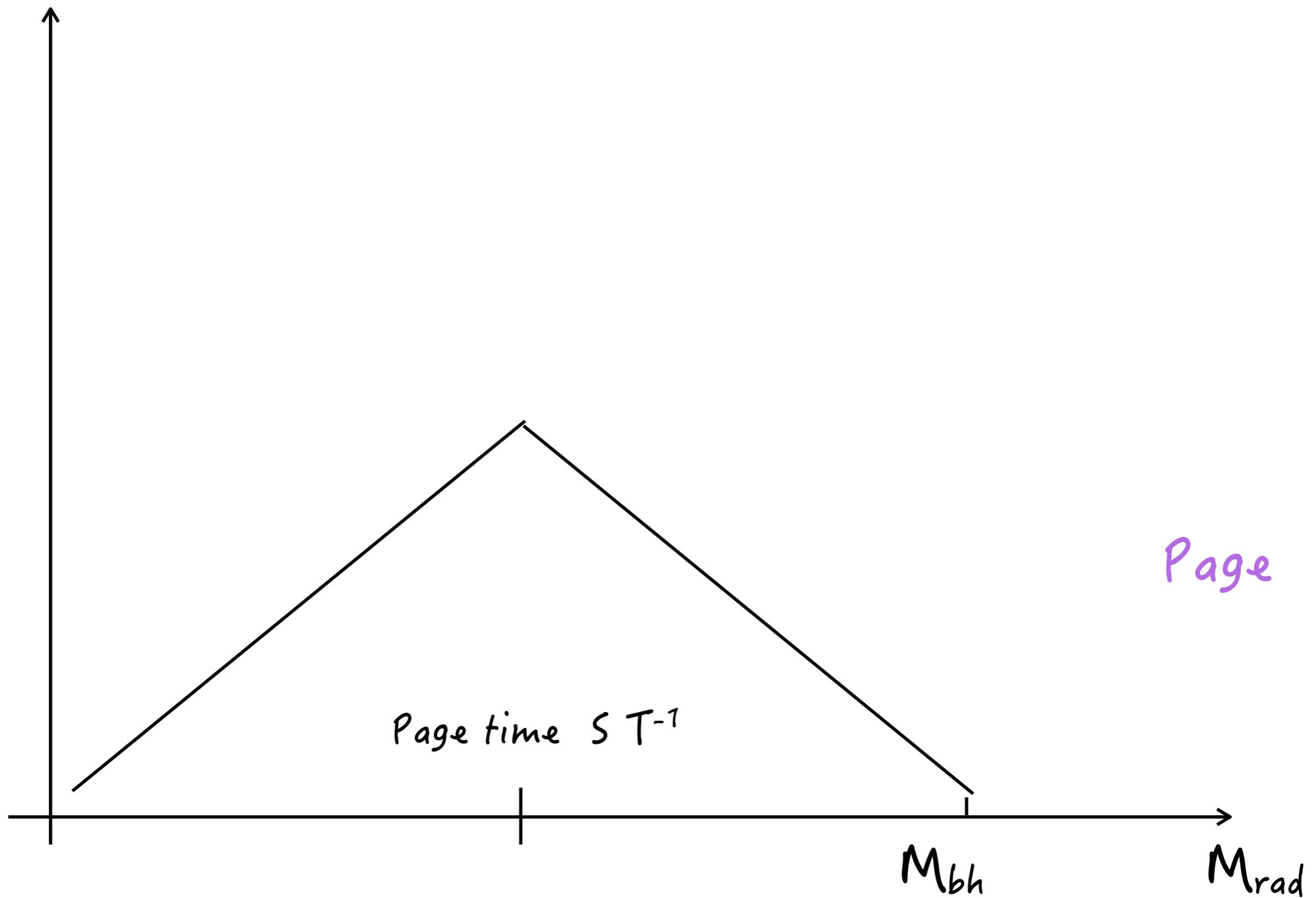
$$S_{\text{rad}} = -\text{Tr} \rho_{\text{rad}} \log \rho_{\text{rad}}$$



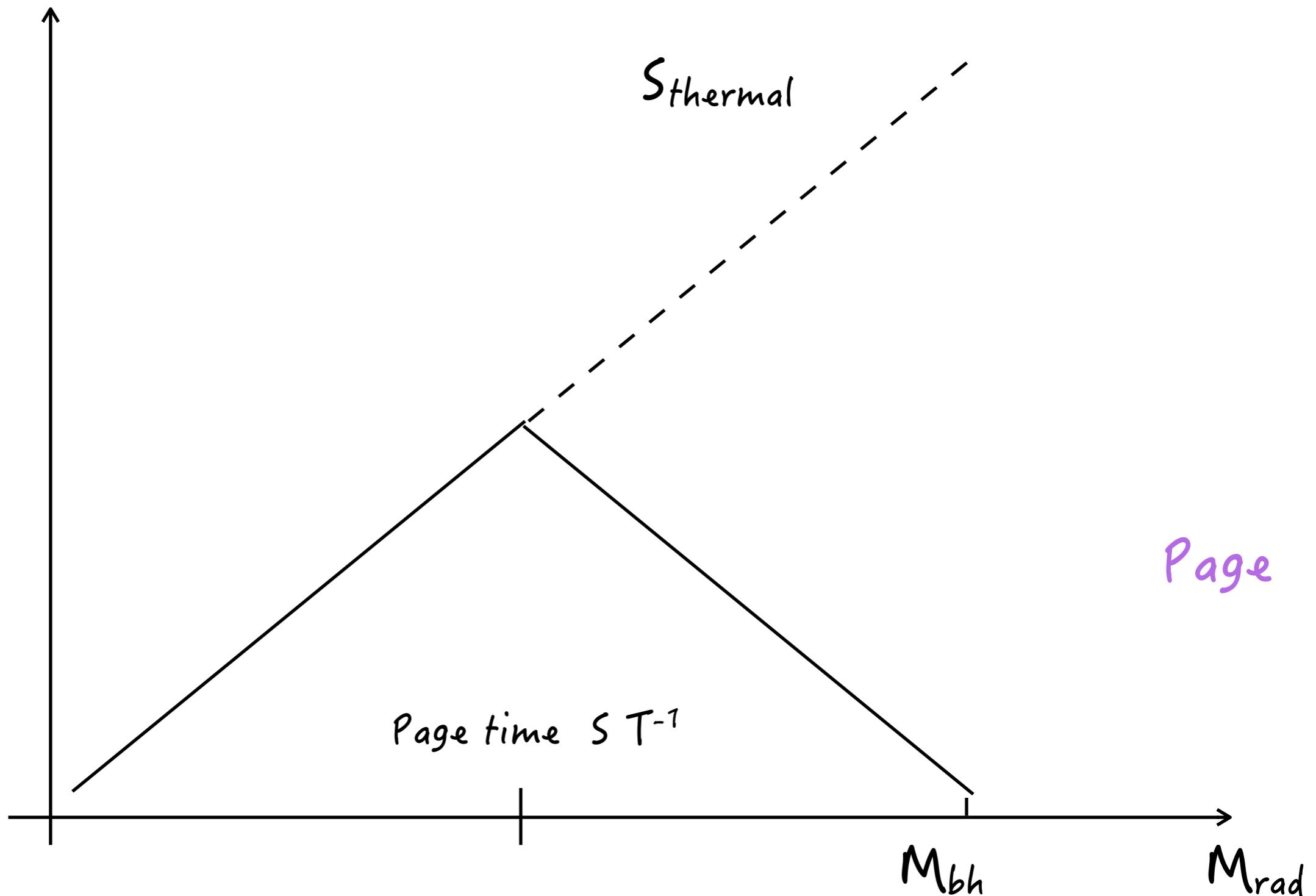
$$S_{\text{rad}} = -\text{Tr} \rho_{\text{rad}} \log \rho_{\text{rad}}$$



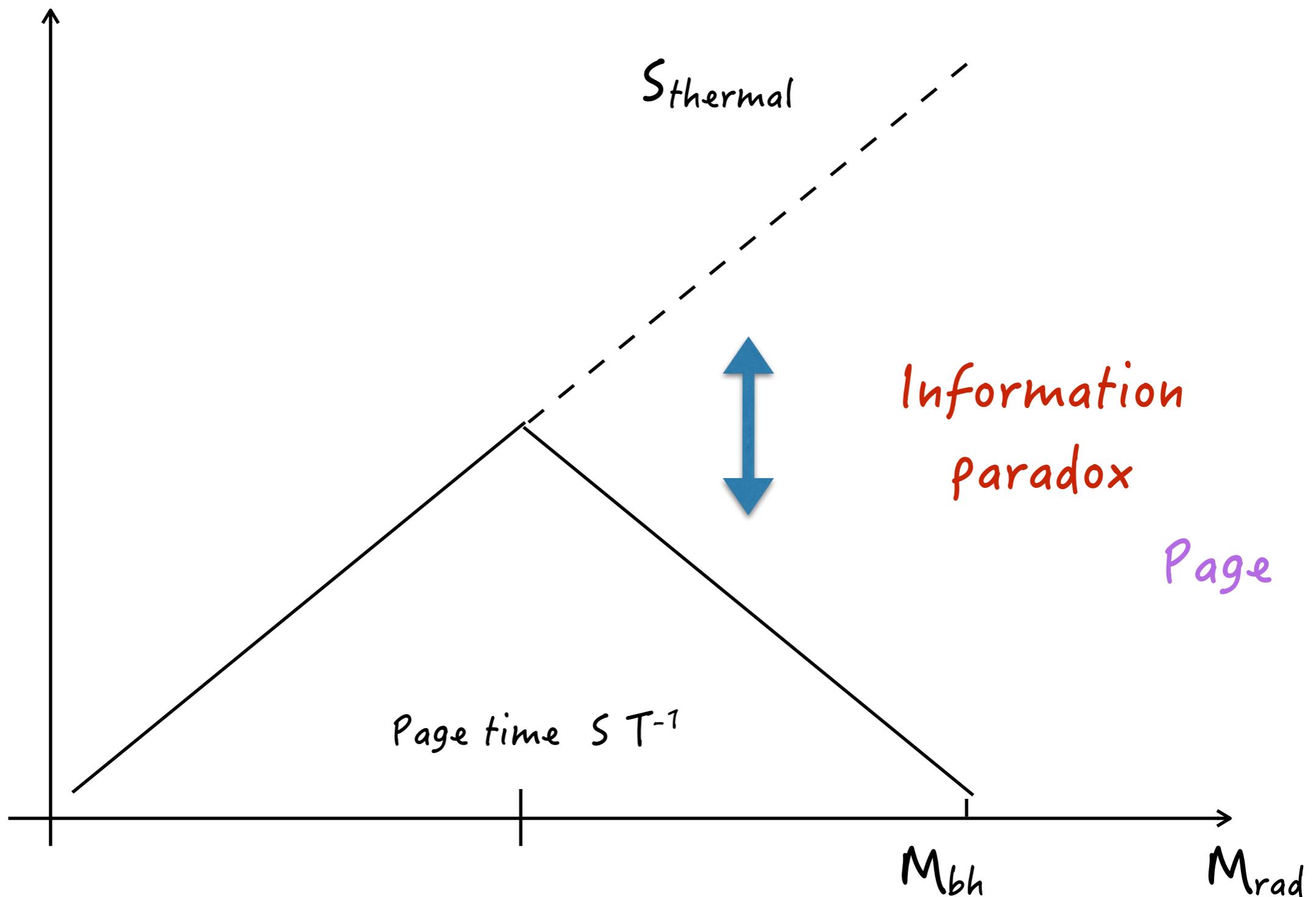
$$S_{\text{rad}} = -\text{Tr} \rho_{\text{rad}} \log \rho_{\text{rad}}$$



$$S_{\text{rad}} = -\text{Tr} \rho_{\text{rad}} \log \rho_{\text{rad}}$$



$$S_{\text{rad}} = -\text{Tr} \rho_{\text{rad}} \log \rho_{\text{rad}}$$



At the Page time, memory effects from initial state on each rad quantum are of order

$$e^{-Tt_{\text{Page}}} \sim e^{-S}$$

But the fine-grained von Neumann entropy of the radiation involves e^S measurements of $O(S)$ degrees of freedom, so we can get an $O(1)$ deviation for S_{rad}

This is actually "natural" ...

$$\begin{pmatrix} 1 & & & & \\ 0 & 0 & & & 0 \\ 0 & & \ddots & & \\ & \ddots & & & \\ 0 & & & 0 & \end{pmatrix}$$

$\epsilon^S \times \epsilon^S$ density
matrix of pure
state

This is actually "natural" ...

$$U^+ \begin{pmatrix} 1 & & & & \\ 0 & 0 & & & \\ 0 & & \ddots & & \\ 0 & & & 0 & \\ & & & & 0 \end{pmatrix} U$$

$\epsilon^S \times \epsilon^S$ density
matrix of pure
state

random $\epsilon^S \times \epsilon^S$
unitary

This is actually "natural" ...

$$U^+ \begin{pmatrix} 1 & & & & \\ 0 & 0 & & & \\ 0 & & \ddots & & \\ 0 & & & 0 & \\ & & & & 0 \end{pmatrix} U = e^{-s} \begin{pmatrix} & & & \\ & & & \\ & & \text{Random} & \\ & & 0(1)_{ij} & \\ & & & \end{pmatrix}$$

$e^s \times e^s$ density
matrix of pure
state

random $e^s \times e^s$
unitary

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$e^S \times e^S$ density
matrix of pure
state

random $e^S \times e^S$
unitary

$$\rho_{ij} = e^{-S} \delta_{ij} + O(e^{-S})_{ij}$$

This is actually "natural" ...

$$U^+ \begin{pmatrix} 1 & & & & \\ 0 & 0 & & & \\ 0 & & \ddots & & \\ 0 & & & 0 & \\ & & & & 0 \end{pmatrix} U = e^{-S} \begin{pmatrix} & & & \\ & & & \\ & & \text{Random} & \\ & & O(1)_{ij} & \\ & & & \end{pmatrix}$$

$e^S \times e^S$ density matrix of pure state

random $e^S \times e^S$ unitary

$$\rho_{ij} = e^{-S} \delta_{ij} + O(e^{-S})_{ij}$$

Hawking estimate

Nonperturbative correction of order $\exp(-1/\alpha_{\text{eff}})$

This is actually "natural" ...

$$U^+ \begin{pmatrix} 1 & & & & \\ 0 & 0 & & & 0 \\ 0 & 0 & \ddots & & \\ & & & 0 & \\ 0 & & & & 0 \end{pmatrix} U = e^{-S} \begin{pmatrix} & & & \\ & & \text{Random} & \\ & & O(1)_{ij} & \\ & & & \end{pmatrix}$$

$e^S \times e^S$ density matrix of pure state

random $e^S \times e^S$ unitary

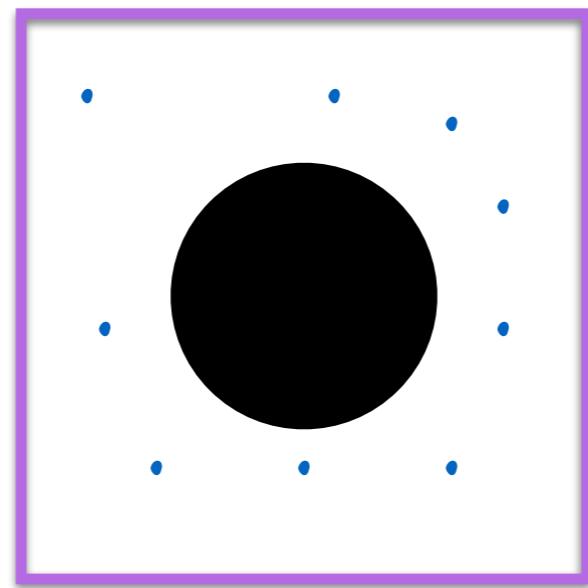
PURIFICATION
AT PAGE TIME
NOT EXPENSIVE!

Hawking estimate

$$\rho_{ij} = e^{-S} \delta_{ij} + O(e^{-S})_{ij}$$

Nonperturbative correction of order $\exp(-1/\alpha_{\text{eff}})$

A sharper way of getting the same insight is to box the black hole and have it equilibrate at fixed box temperature



Gibbons-Hawking

What is a good box for gravity?

According to Hawking & Page ... AdS!

MOREOVER ...

In this case there is a candidate for the exact Hamiltonian
with a black hole spectrum

$$S_{\text{AdSbh}_{d+1}} \sim N_*^{1/d} (M\ell)^{\frac{d-1}{d}} \sim S_{\text{CFT}_d}$$

$$N_* \sim \frac{\ell^{d-1}}{G}$$

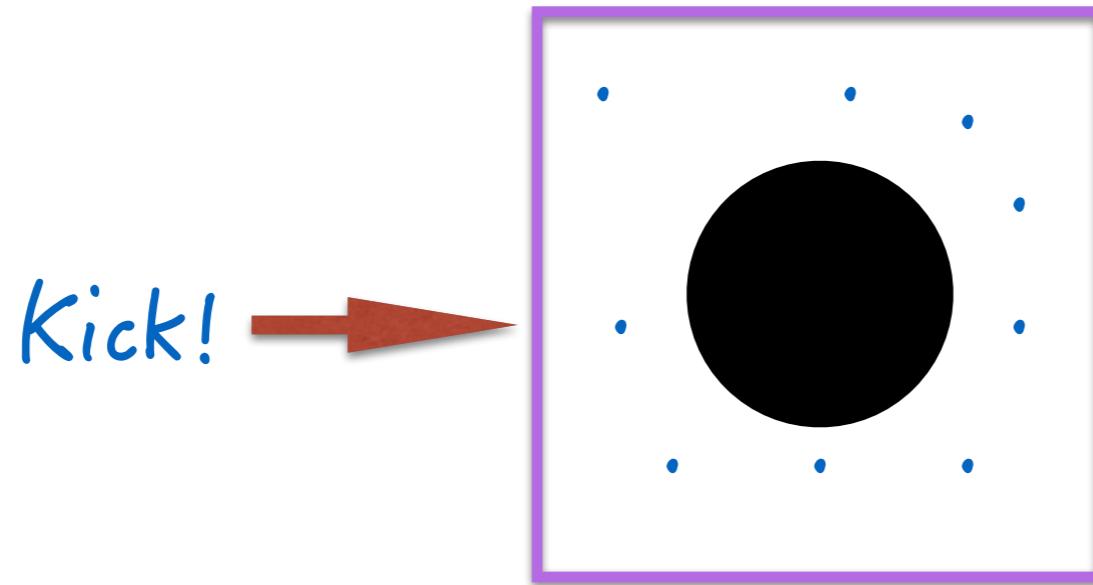
ℓ = curvature radius of AdS

Enter AdS/CFT !

Maldacena

BOX and everything inside = CFT

The decay of local box kicks reveals the absorption of information by the black hole



$$\langle B(t) B(0) \rangle \sim e^{-\Gamma t}$$

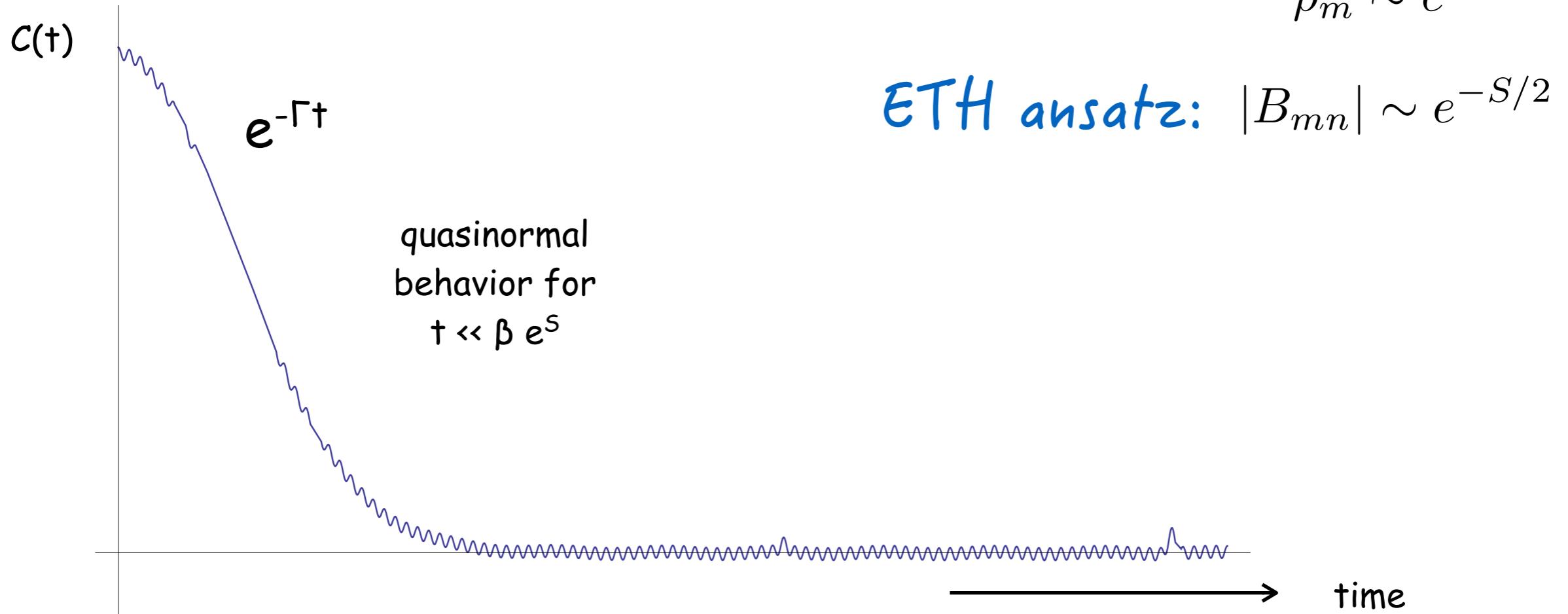
If the spectrum of bh microstates is approximated as continuous (neglecting spacings of order e^{-S}) then the correlation will continue to decay forever

But the long-time average must be bounded-below
in any system with discrete spectrum

$$C(t) = \text{Tr} [\rho B(t) B(0)] = \text{Tr} [\rho B e^{itH} B e^{-itH}]$$

$$C(t) = \sum_{mn} \rho_m |B_{mn}|^2 e^{i(E_m - E_n)t} \quad C(0) = 1$$

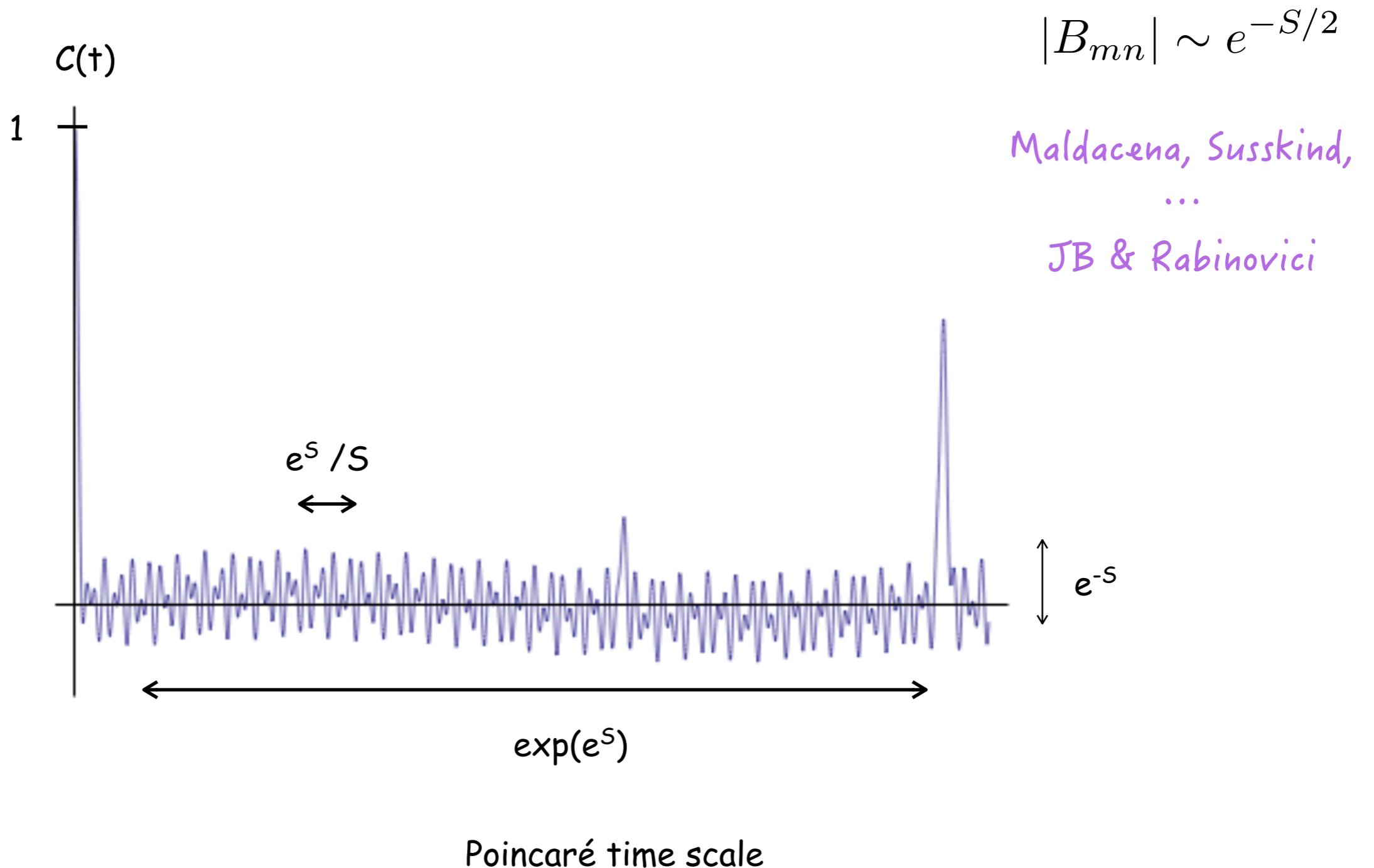
$$\rho_m \sim e^{-S}$$



$$C(t) = \sum_{mn} \rho_m |B_{mn}|^2 e^{i(E_m - E_n)t}$$

$$C(0) = 1$$

$$\rho_m \sim e^{-S}$$



The church of the doubled Hilbert space

Any thermal box can be obtained by tracing over a second identical copy, if appropriately entangled into a global pure state

$$\rho_R = \text{Tr}_L \sum_n C_n \Psi_n^L \otimes \Psi_n^R$$

$$(C_n)_{\text{thermal}} = \left[\frac{e^{-\beta E_n}}{\sum_m e^{-\beta E_m}} \right]^{1/2}$$

BUT!!

If the entanglement basis is taken to be the high-energy band of two "entangled" CFTs ...

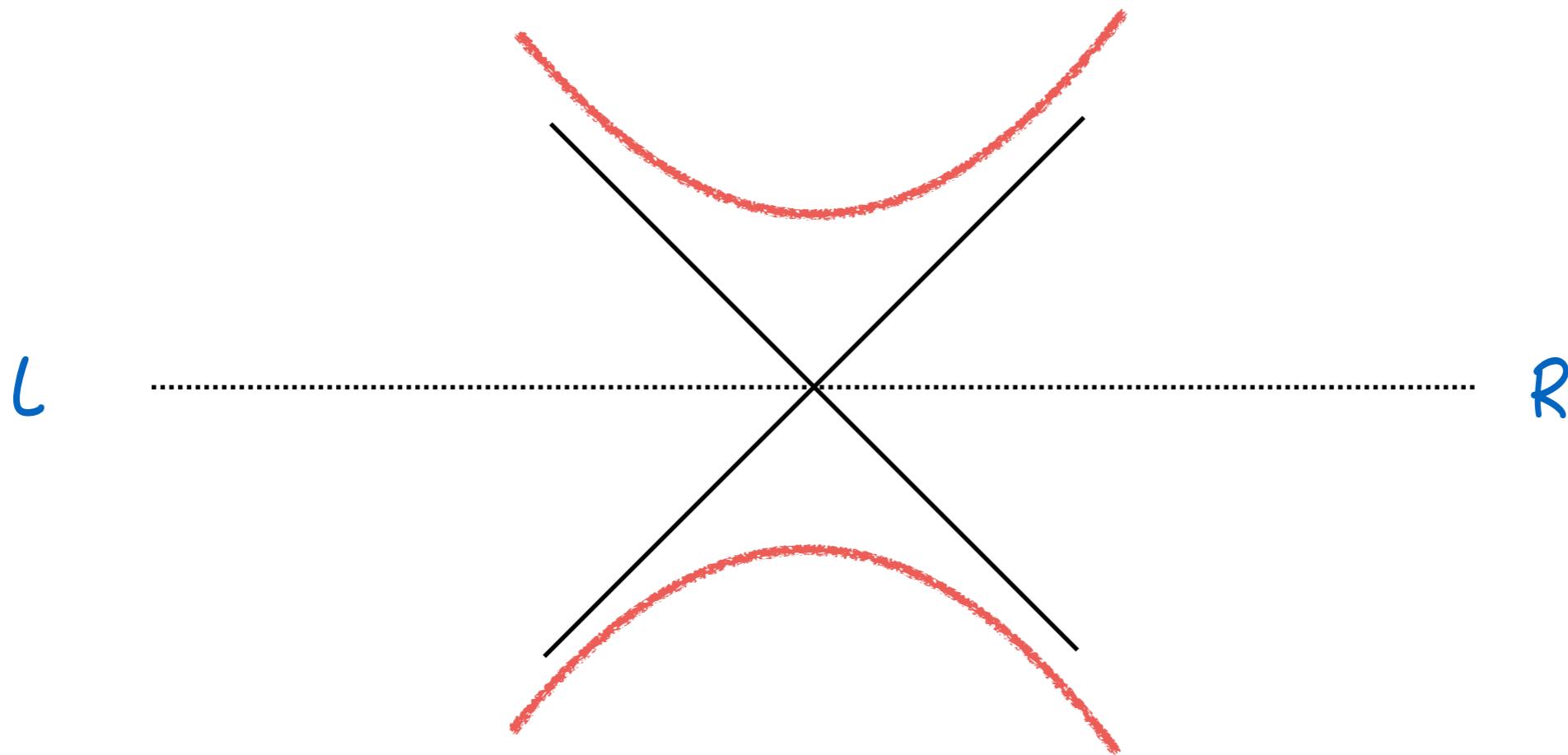
$$|TFD\rangle \sim \sum_{E_n} e^{-\beta E_n/2} |E_n\rangle_L \otimes |E_n\rangle_R$$

neglecting the tiny e^{-S} spacings,
we can approximate by continuous spectrum of fields in
the background of an AdS black hole, to get ...

$$\int_E e^{-\beta E/2} |E\rangle_L \otimes |E\rangle_R$$

The HH state of the bulk fields !

Two black holes sharing an interior



$$|\text{HH}\rangle \sim \int_E e^{-\beta E/2} |E\rangle_L \otimes |E\rangle_R$$

So, by bringing in the ℓ^{-S} details, AdS/CFT implies that the background geometry itself is a result of entanglement of the two disconnected CFTs

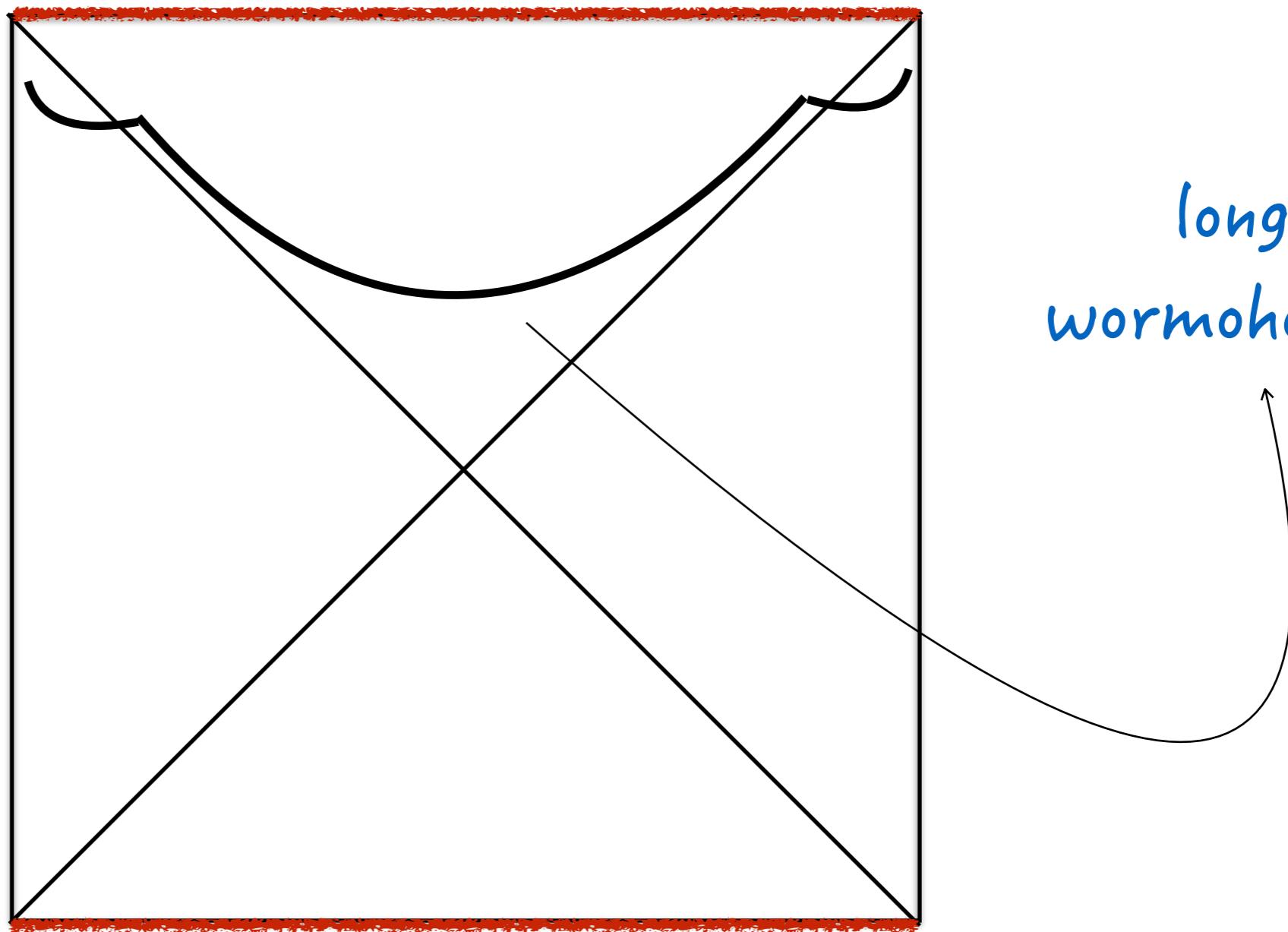
SLOGAN: $EPR = ER$

Maldacena-Susskind

Accumulating a density of entanglement
of $S \gg 1$ well-separated Bell pairs within a
transversal size of order $(G S)^{1/2}$ seems to
generate a geometrical bridge of area $G S$

What about dynamics?

$$|TFD\rangle_t \sim \sum_n e^{-\beta E_n/2 + 2iE_n t} |E_n\rangle_L \otimes |E_n\rangle_R$$



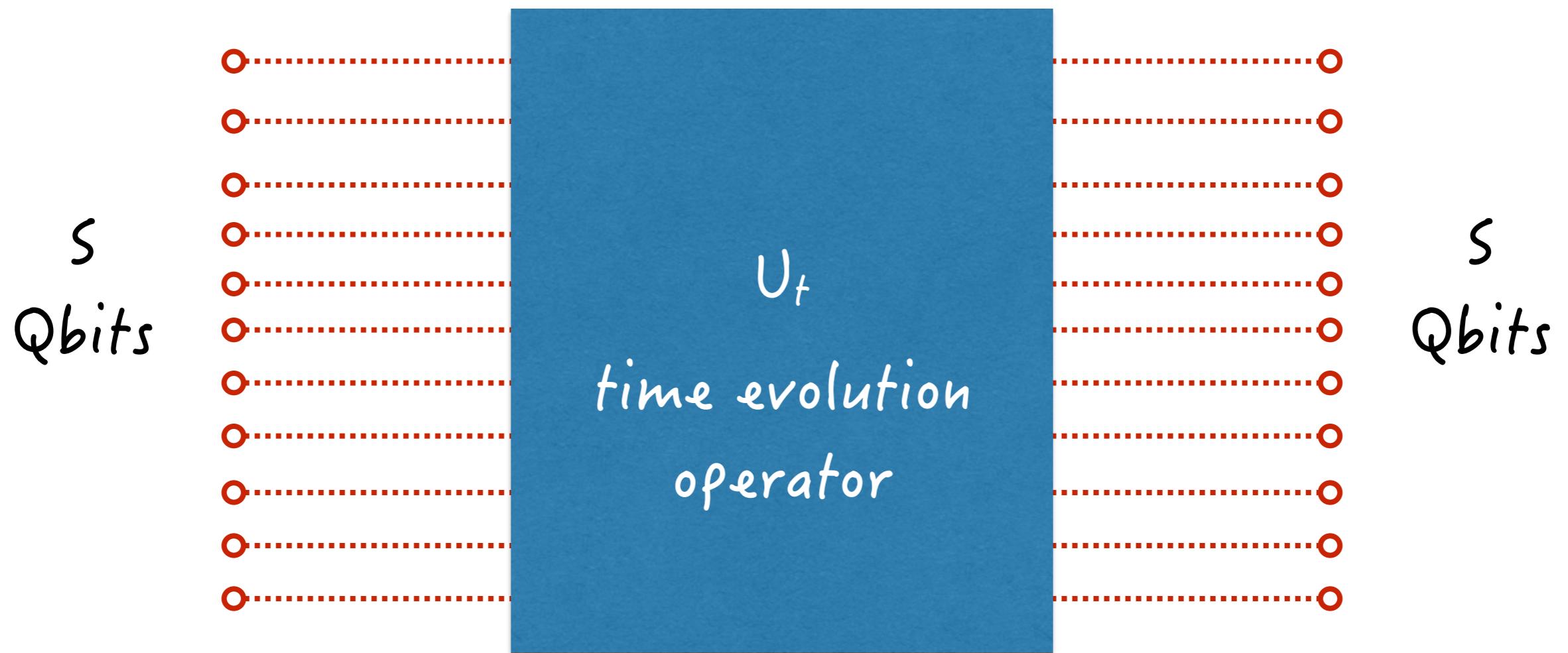
What is the entanglement interpretation of the growing wormhole?

A hint is obtained by looking at a system of 25 entangled Qbits



What is the entanglement interpretation of the growing wormhole?

A hint is obtained by looking at a system of 25 entangled Qbits



A notion of "size" for U is that of quantum complexity theory



Complexity (U) = minimal size of quantum circuit

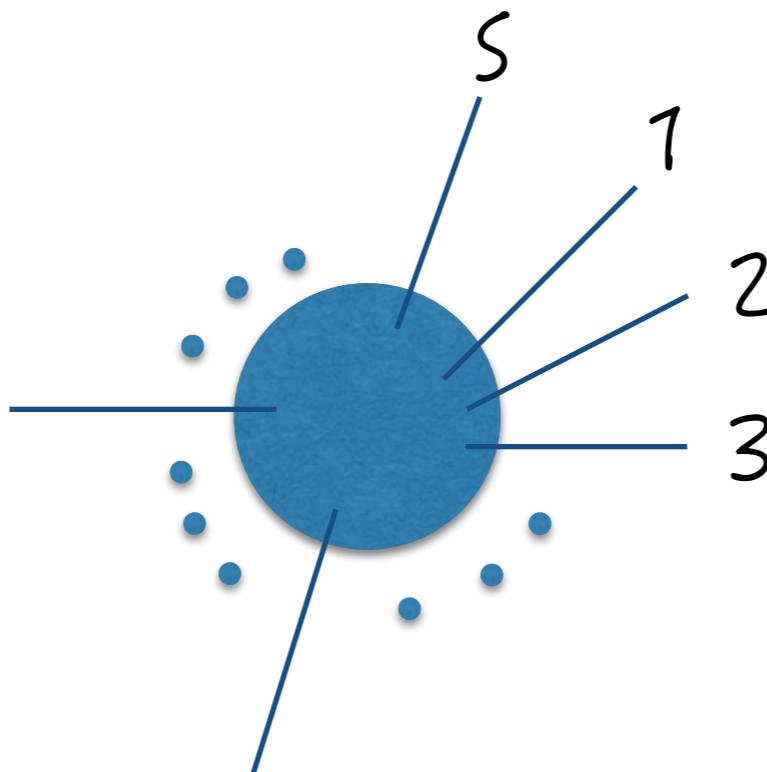
Hartman-Maldacena, Susskind, ...

Parametrizing complexity of entanglement

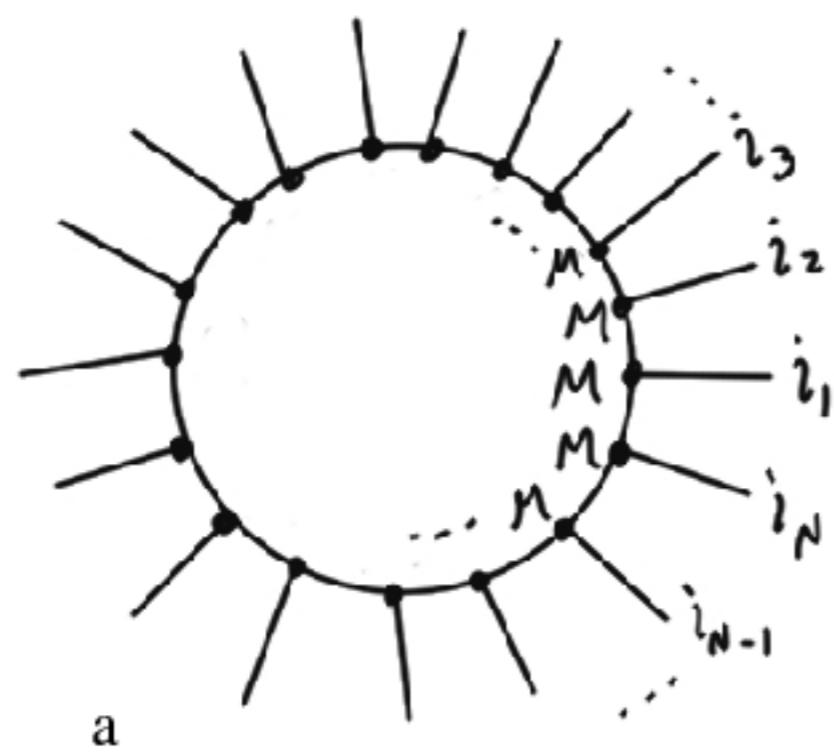
Pick a tensor decomposition of Hilbert space of dimension $\exp(S)$ into S factors of $O(1)$ dimension

$$\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2 \otimes \cdots \otimes \mathcal{H}_S$$

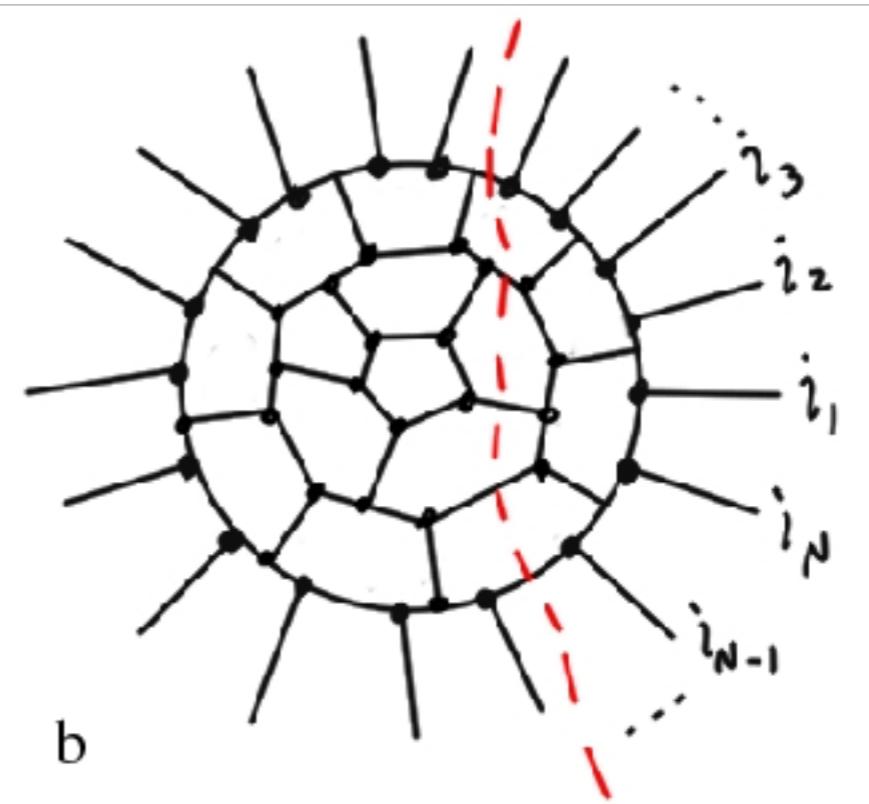
A tensor of S indices gives a generic state



The decomposition of the big tensor in small building blocks gives a notion of "complexity of entanglement"



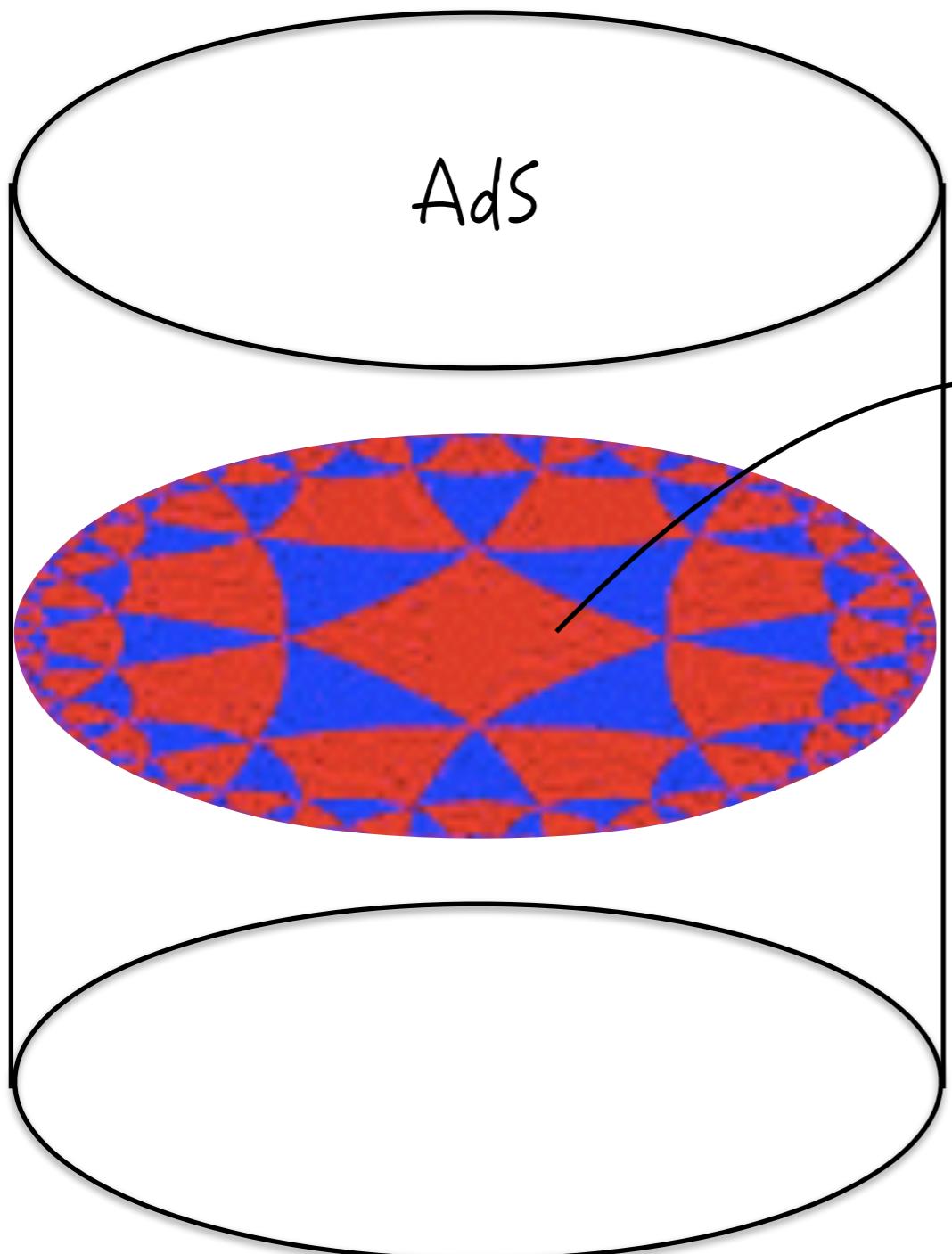
rather simple
entanglement pattern



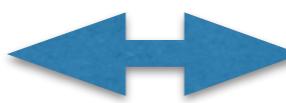
somewhat more complex
entanglement pattern

picture from M van Raamsdonk's review

The key example



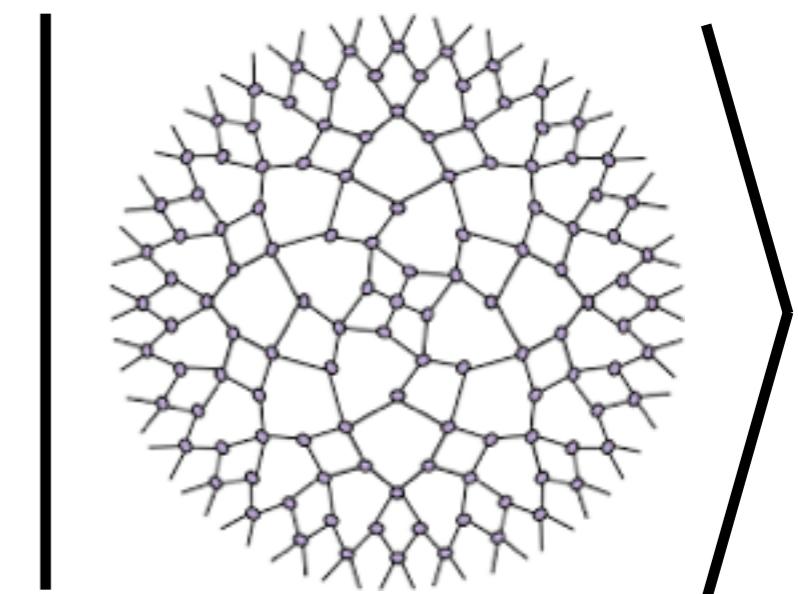
Size of TN



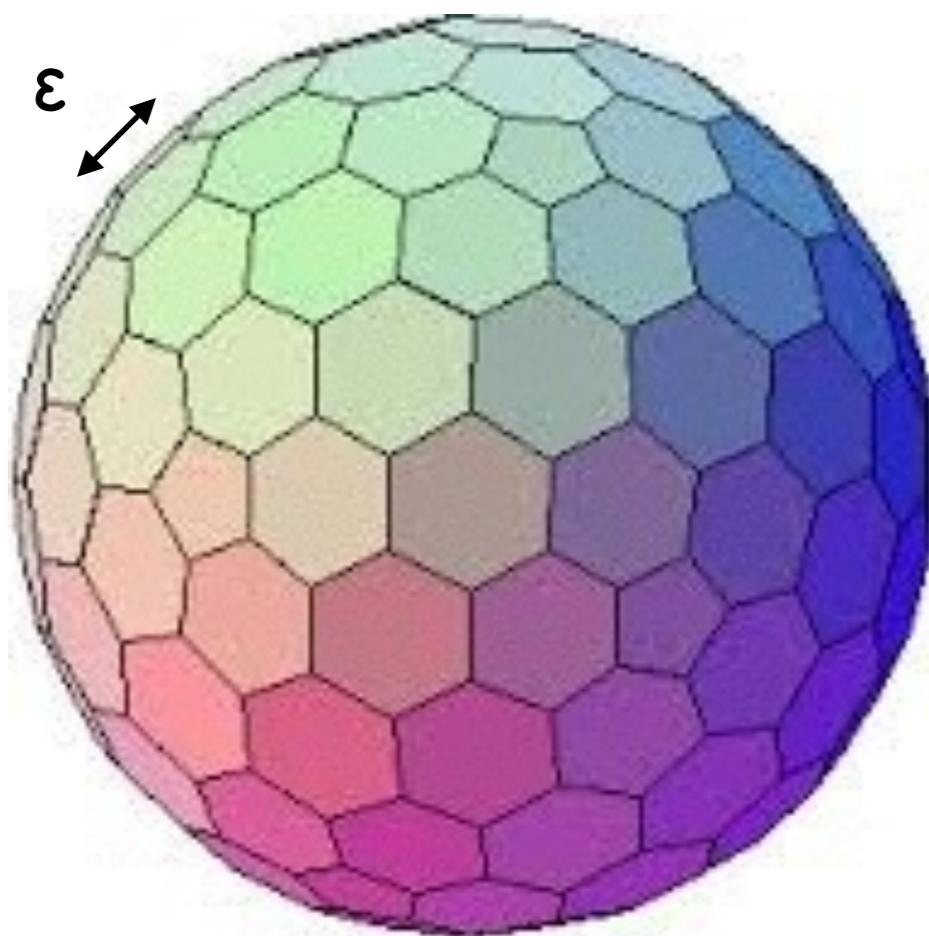
Volume of spatial slice

Vidal,
Swingle

Hyperbolic TN ansatze for
conformal ground states



In quantum complexity theory, we try to map out the quantum Hilbert space as if it was a classical phase space



$$\frac{U(e^S)}{U(e^S - 1) \times U(1)}$$

$$\# \text{ cells} \sim \exp(e^S \log(1/\epsilon))$$

Just like classical entropy is bounded by the coarse-grained volume of phase space, quantum complexity is bounded by the number of ϵ -cells

Being serious about e^{-S} effects ...

we came to the astonishing suggestion that the very fabric of space could be a "condensate" of pure quantum entanglement

Plenty of circumstantial evidence from AdS/CFT ...

ENTANGLEMENT \leftrightarrow BULK CONNECTIVITY

COMPLEXITY \leftrightarrow BULK VOLUME

A list of open questions & problems

- Need exactly calculable toy models of AdS/CFT along the lines of SYK model
- Give a “renormalized” definition of quantum complexity for continuum CFTs
- Can tensor networks describe bulk gravitons?
- What is the space-time meaning of quantum complexity saturation?
- Can we define approximate local observables for black hole interiors?
- Are there obstructions related to firewalls and/or fuzzballs?

THANK YOU

