

The Rogue Waves of NLS and Their Universality

Guido Mazzuca
Tulane University

Joint work with A. Gkougkou, K. D. T-R McLaughlin, T. Grava, R. Jenkins, M. Girotti and M. Yattselev

References : Arxiv:2602.05101 & DOI: 10.1088/1361-6544/ae1435 (Nonlinearity)

Rogue Waves



Credits to Ben Hartley
Location: Nazaré, Portugal

Rogue Waves



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Rogue Waves

- Freak or monster waves



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- Freak or monster waves
- Localized in space-time



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Rogue Waves

- Freak or monster waves
- Localized in space-time
- Highly disruptive (usually)



Credits to Ben Hartley
Location: Nazaré, Portugal

Rogue Waves

- Freak or monster waves
- Localized in space-time
- Highly disruptive (usually)
- Difficult to predict



Credits to Ben Hartley
Location: Nazaré, Portugal







Credits: Bright Side (Youtube)



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The NLS equation

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$$i\partial_t\psi = -\frac{1}{2}\partial_{xx}\psi - |\psi|^2\psi$$

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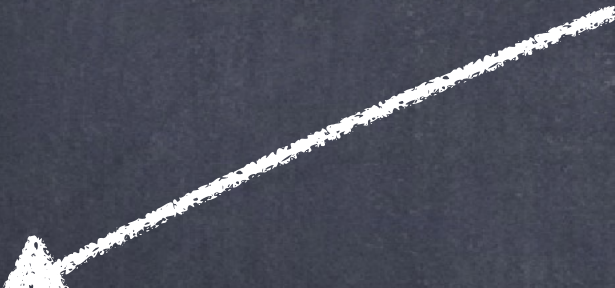
Integrable



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Integrable



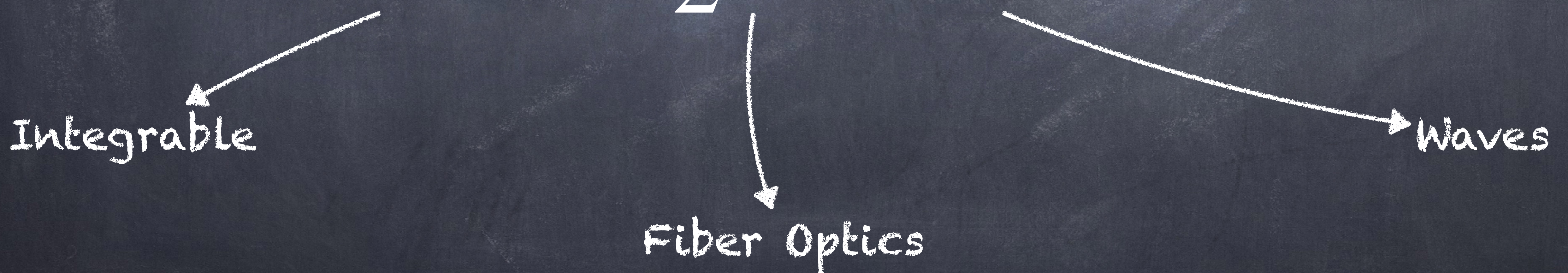
Fiber Optics



The NLS equation

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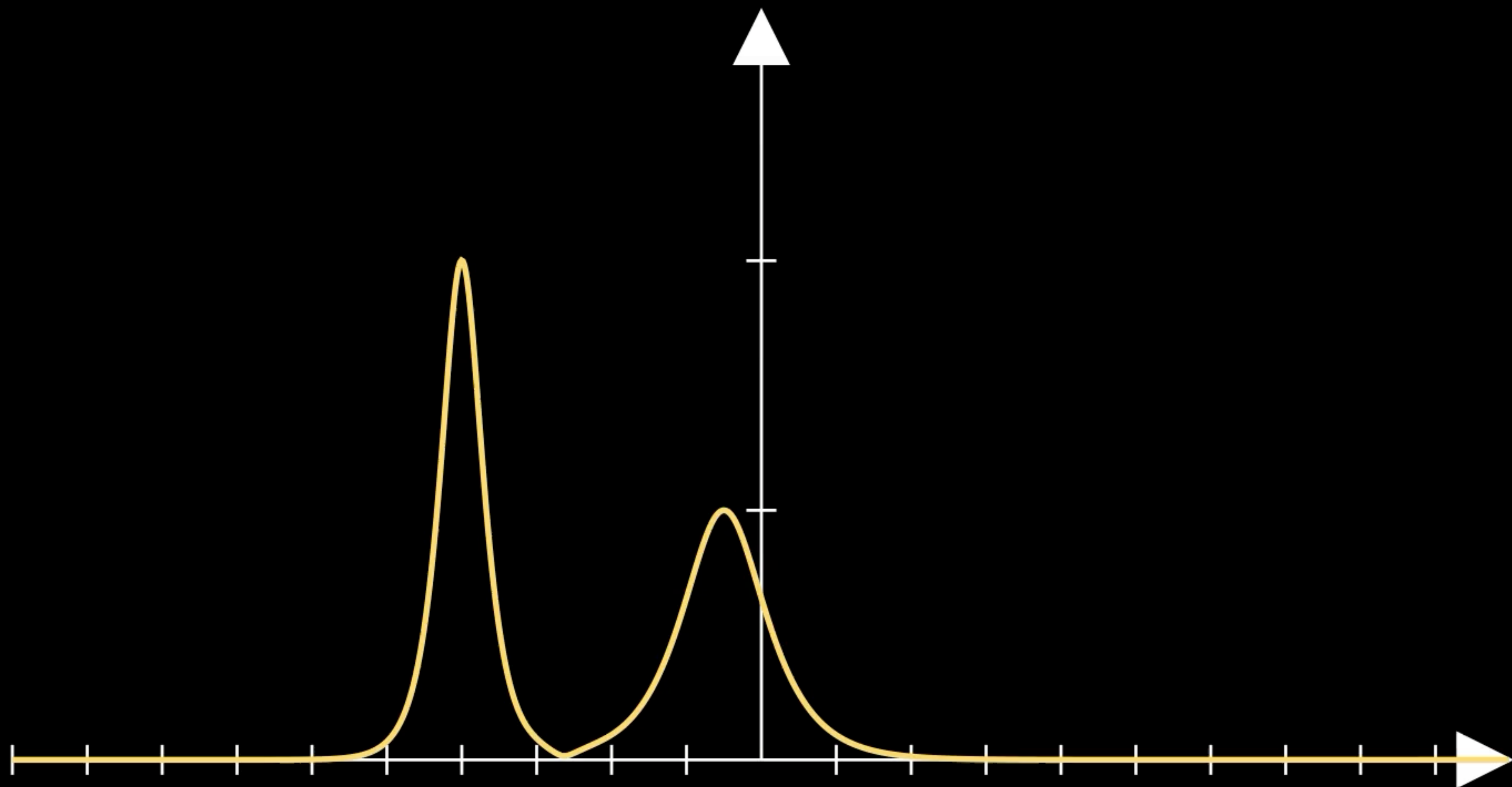


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graph TD; A["i\partial_t\psi = -\frac{1}{2}\partial_{xx}\psi - |\psi|^2\psi"] --> B[Integrable]; A --> C[Fiber Optics]; A --> D[Waves];
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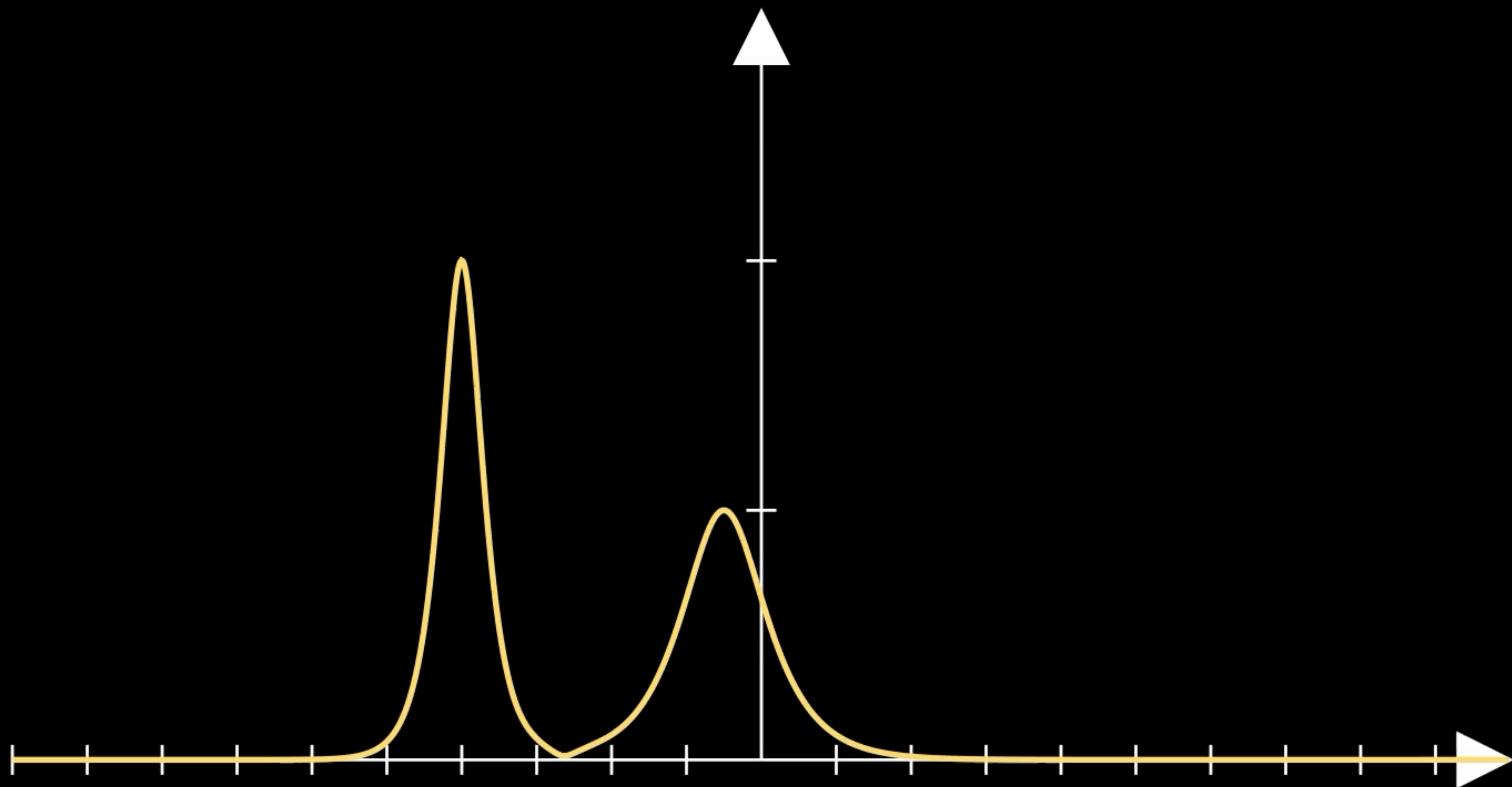
Fiber Optics

Waves

True Nonlinear Evolution
Linear Superposition



True Nonlinear Evolution
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Questions

$$i\partial_t\psi = -\frac{1}{2}\partial_{xx}\psi - |\psi|^2\psi$$

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- How likely is it?

Questions

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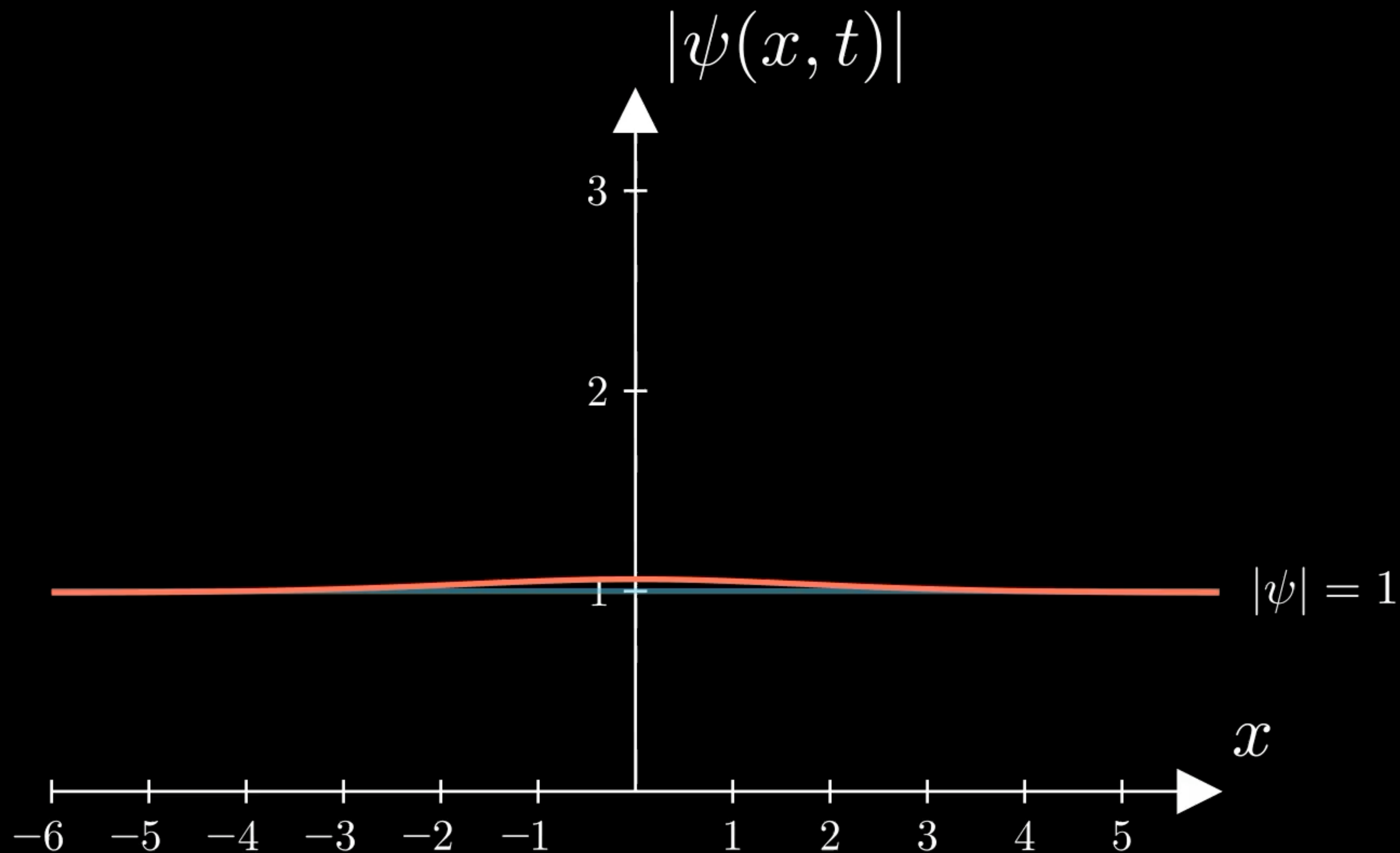
- What is a rogue wave for NLS?
- How likely is it?
- Can we describe it analytically?

Rogue Wave

$$|\psi(x, t)| \gg \mathbb{E}[|\psi(x, t)|]$$

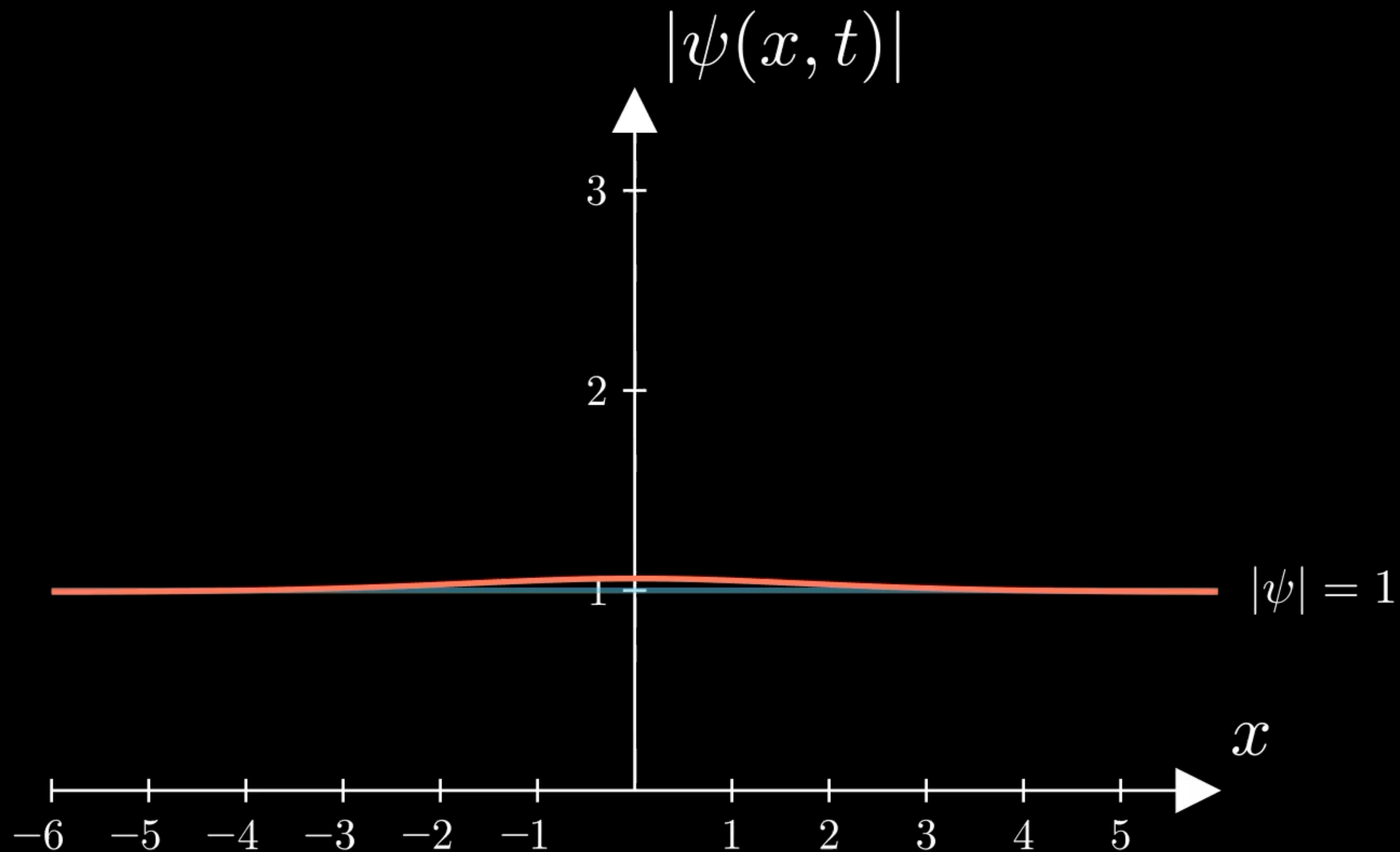
Peregrine Breather Profile

Time (t): -4.00



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Time (t): -4.00



Communications on

PURE AND APPLIED MATHEMATICS

**Universality for the Focusing Nonlinear Schrödinger Equation
at the Gradient Catastrophe Point: Rational Breathers and
Poles of the Tritronquée Solution to Painlevé I**

[Marco Bertola](#) ✉ [Alexander Tovbis](#) ✉

First Universality result of the day

How likely is it?

How likely is it?

$$F(\delta, T) = \mathbb{P}\left(\max_{x \in (-L, L)} |\psi(x, T)| > \delta\right)$$

How Likely is it?

$$F(\delta, T) = \mathbb{P}\left(\max_{x \in (-L, L)} |\psi(x, T)| > \delta\right)$$

Rogue waves and large deviations in deep sea

Giovanni Dematteis^{a,b}, Tobias Grafke^{a,c}, and Eric Vanden-Eijnden^{a,1}

^aCourant Institute of Mathematical Sciences, New York University, New York, NY 10012; ^bDipartimento di Scienze Matematiche, Politecnico di Torino, I-10129 Torino, Italy; and ^cMathematics Institute, University of Warwick, Coventry CV4 7AL, United Kingdom

Edited by David A. Weitz, Harvard University, Cambridge, MA, and approved December 11, 2017 (received for review June 13, 2017)

The appearance of rogue waves in deep sea is investigated by using the modified nonlinear Schrödinger (MNLS) equation in one spatial dimension with random initial conditions that are assumed to be normally distributed, with a spectrum approximating real-

random initial data drawn from a Gaussian distribution (30). The spectrum of this field is chosen to have a width comparable to that of the Joint North Sea Wave Project (JONSWAP) spectrum (31, 32) obtained from observations in the North Sea. We cal-

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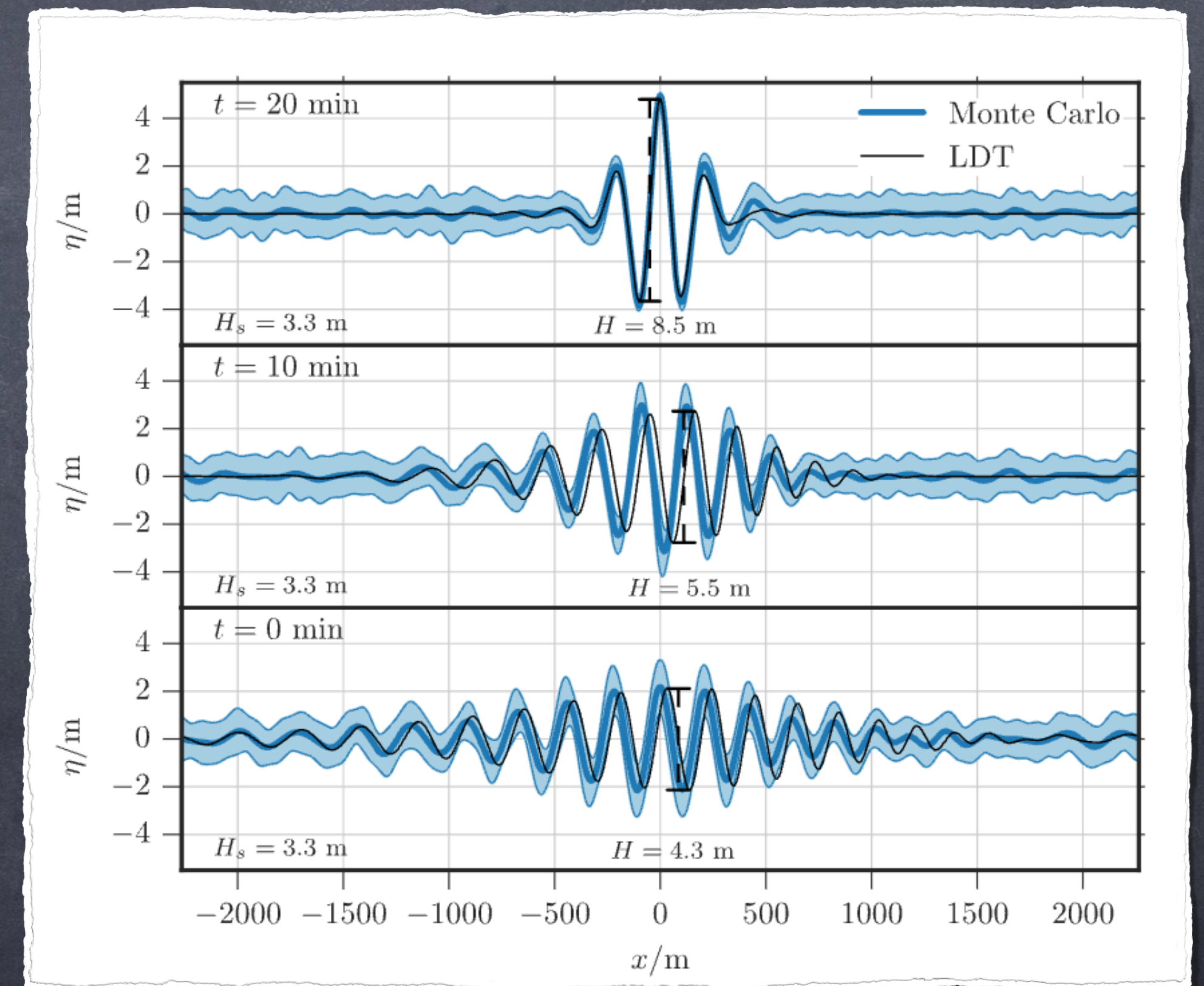
- Obtain a Large Deviation Principle;

$$F(\delta, T) = \mathbb{P}\left(\max_{x \in (-L, L)} |\psi(x, T)| > \delta\right)$$

- Obtain a Large Deviation Principle;
- Characterize the most likely precursor to the rogue wave

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- Obtain a Large Deviation Principle;
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$$F(\delta, T) = \mathbb{P}\left(\sup_{x \in (0, 2\pi)} |\psi(x, T)| > \delta\right)$$



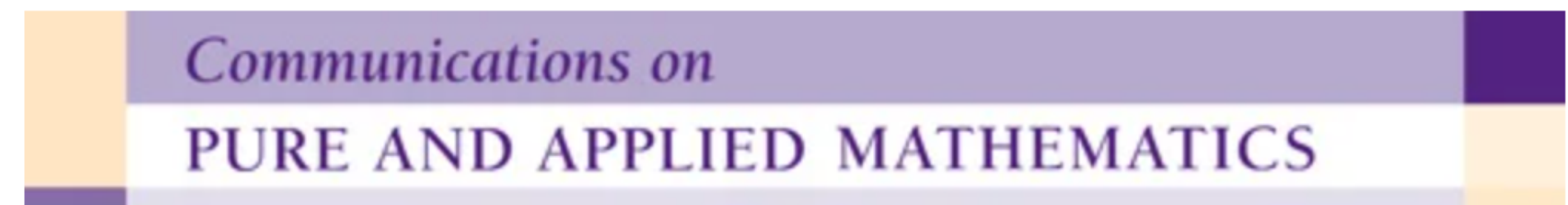
RESEARCH ARTICLE

Large deviations principle for the cubic NLS equation

[Miguel Angel Garrido](#), [Ricardo Grande](#) ✉, [Kristin M. Kurianski](#), [Gigliola Staffilani](#)

$$\lim_{\varepsilon \rightarrow 0^+} \varepsilon \log \left(\mathbb{P} \left(\sup_{x \in (0, 2\pi)} |\psi(x, T)| > \varepsilon^{-\frac{1}{2}} \delta \right) \right) = - \frac{\delta^2}{\sum_k c_k^2}$$

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$$\psi_0(x) = \sum_{k \in \mathbb{Z}} c_k \eta_k e^{ikx}$$

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RESEARCH ARTICLE

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$$\psi_0(x) = \sum_{k \in \mathbb{Z}} c_k \eta_k e^{ikx}$$

1.3.2. *Main ideas in Theorem 1.9.* The goal of this result is to show that the most likely way in which rogue wave arise in this weakly nonlinear context is due to phase synchronization⁵. The construction of the one-sided neighborhood $\mathcal{U}(\varepsilon)$ in (1.13) is based on two-scales: we control the moduli of $\mathcal{O}(|\log \varepsilon|)$ Fourier modes, but we force $\mathcal{O}(\varepsilon^{-1/2})$ phases to almost synchronize at a certain point in space and time.

What is the shape of the NLS
rogue wave?



RHP for the NLS

Pure soliton!



RHP for the NLS

Pure soliton!

$$M(z) \in \text{Mat}(\mathbb{C}, 2)$$

$M(z)$ is analytic in $\mathbb{C} \setminus \bigcup_{j=1}^N \{z_j, \bar{z}_j\}$



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$$\theta(z, x, t) = xz + tz^2$$



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$$M(z) = \mathbb{I} + \frac{1}{2iz} \begin{pmatrix} -m(x, t) & \psi(x, t) \\ \overline{\psi(x, t)} & m(x, t) \end{pmatrix} + \mathcal{O}(z^{-2})$$



RHP for the NLS

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Spectral parameters

$$M(z) = \mathbb{I} + \frac{1}{2iz} \begin{pmatrix} -m(x, t) & \psi(x, t) \\ \overline{\psi(x, t)} & m(x, t) \end{pmatrix} + \mathcal{O}(z^{-2})$$

Theoretical Max

$$\max_{x,t} |\psi(x, t)| \leq 2 \sum_{j=1}^N \mathfrak{I}(z_j)$$

Theoretical Max

$$\max_{x,t} |\psi(x,t)| \leq 2 \sum_{j=1}^N \Im(z_j)$$

$$a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j} \quad \bigoplus \quad c_j = \frac{1}{a'(z_j)} \quad \longrightarrow$$

$$|\psi(0,0)| = 2 \sum_{j=1}^N \Im(z_j)$$

*See NN and RW papers for a proof

Solitons Synchronization

$$z_j = v_j + i\mu \quad v_j \in [-j\Delta, -(j-1)\Delta] \quad c_j = 2i\mu e^{-2iz_j^2}$$

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$$c_j = \frac{1}{a'(z_j)} \approx 2i\mu$$

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$$\frac{1}{N\mu} \psi \left(\frac{2X}{N\Delta}, 1 + \frac{T}{(N\Delta)^2} \right) = 2\psi_0(X, T) + \mathcal{O} \left(\max \left\{ \frac{1}{N}, \frac{1}{\Delta} \right\} \right),$$

$$\psi_0(X, T) = - \int_0^1 e^{i \left(2Xs - \frac{T}{2}s^2 \right)} ds.$$

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$$\frac{1}{N\mu} \psi_N \left(\frac{2X}{NV}, 1 \right) = -2e^{iX} \frac{\sin(X)}{X} + \mathcal{O} \left(\max \left\{ \frac{1}{N}, \frac{1}{\Delta} \right\} \right)$$

Soliton Synchronization

$$z_j = v_j + i\mu_j \quad v_j = -j\Delta = -\alpha j N^\gamma, \gamma > \frac{1}{2} \quad c_j = 2i\mu_j e^{-2iz_j^2}$$

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Iid sub-exp



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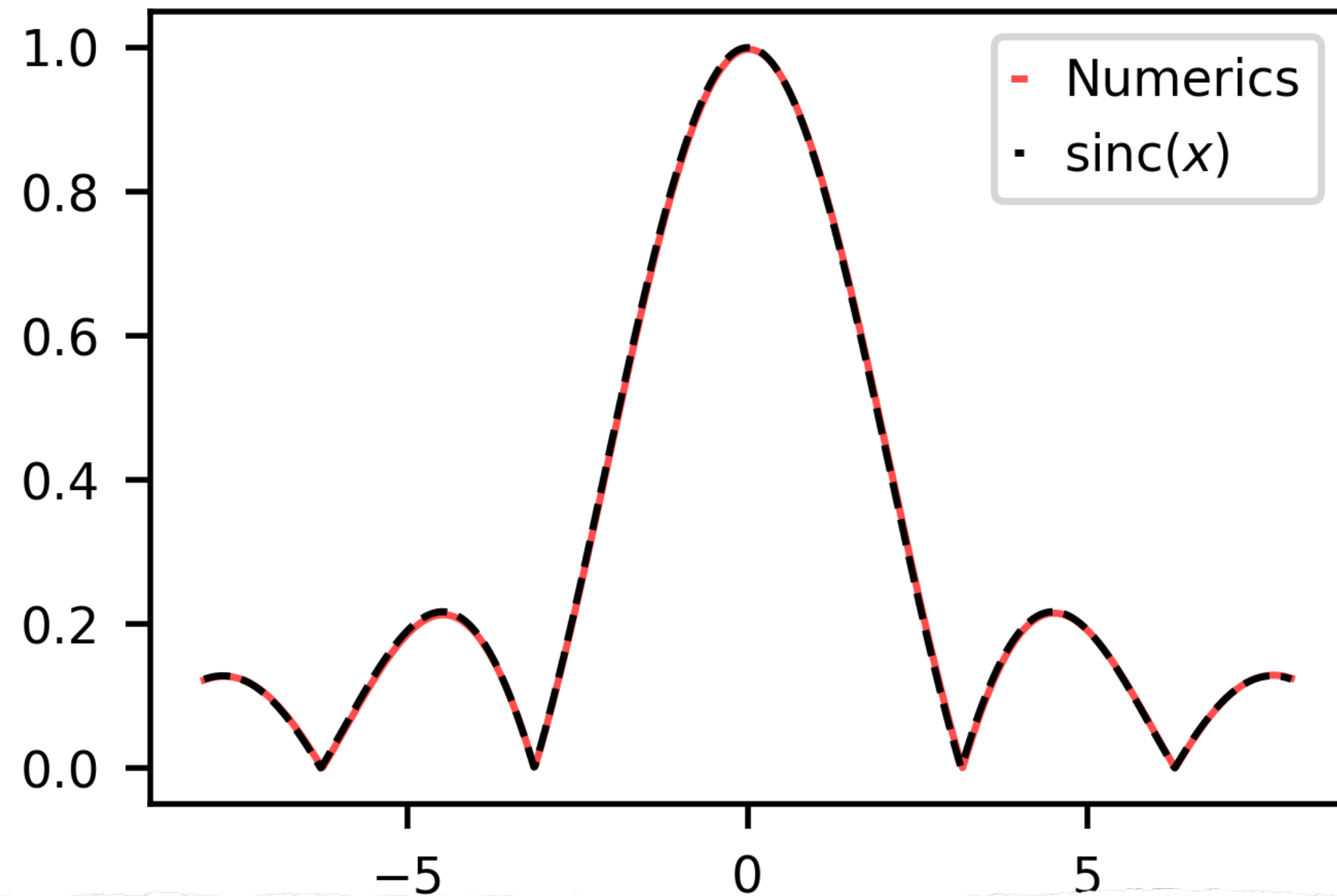
Iid sub-exp



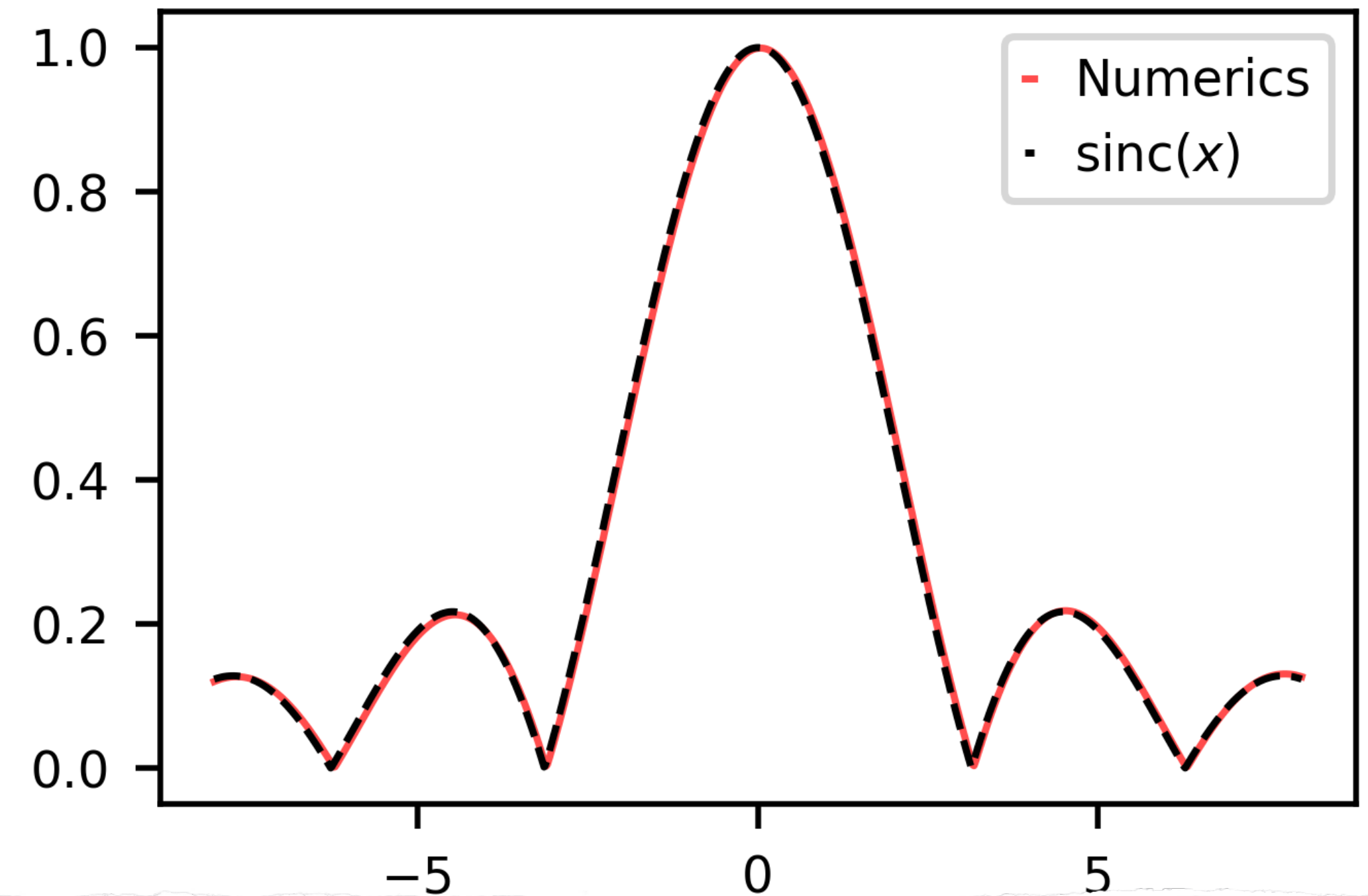
$$\frac{1}{2N\mu_{\mathcal{D}}} \psi \left(\frac{2X}{N\Delta}, 1 + \frac{T}{(N\Delta)^2} \right) \xrightarrow{N \rightarrow \infty} \psi_0(X, T)$$

Solitons Synchronization

beta distribution, $N = 400, \Delta = 200$



χ -distribution, $N = 400, \Delta = 200$



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RW and Painlevé

$$z_j = v_j + i\mu_j \qquad c_j = \frac{1}{a'(z_j)} \qquad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

RW and Painlevé

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Iid sub-exp

The diagram illustrates the decomposition of the complex number z_j into its real part v_j and imaginary part $i\mu_j$. Both v_j and $i\mu_j$ are circled in red. Two red arrows originate from a red rectangular box labeled 'Iid sub-exp' and point to the circled terms v_j and $i\mu_j$, indicating that both components are influenced by independent and identically distributed sub-exponential variables.

RW and Painlevé

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Iid sub-exp

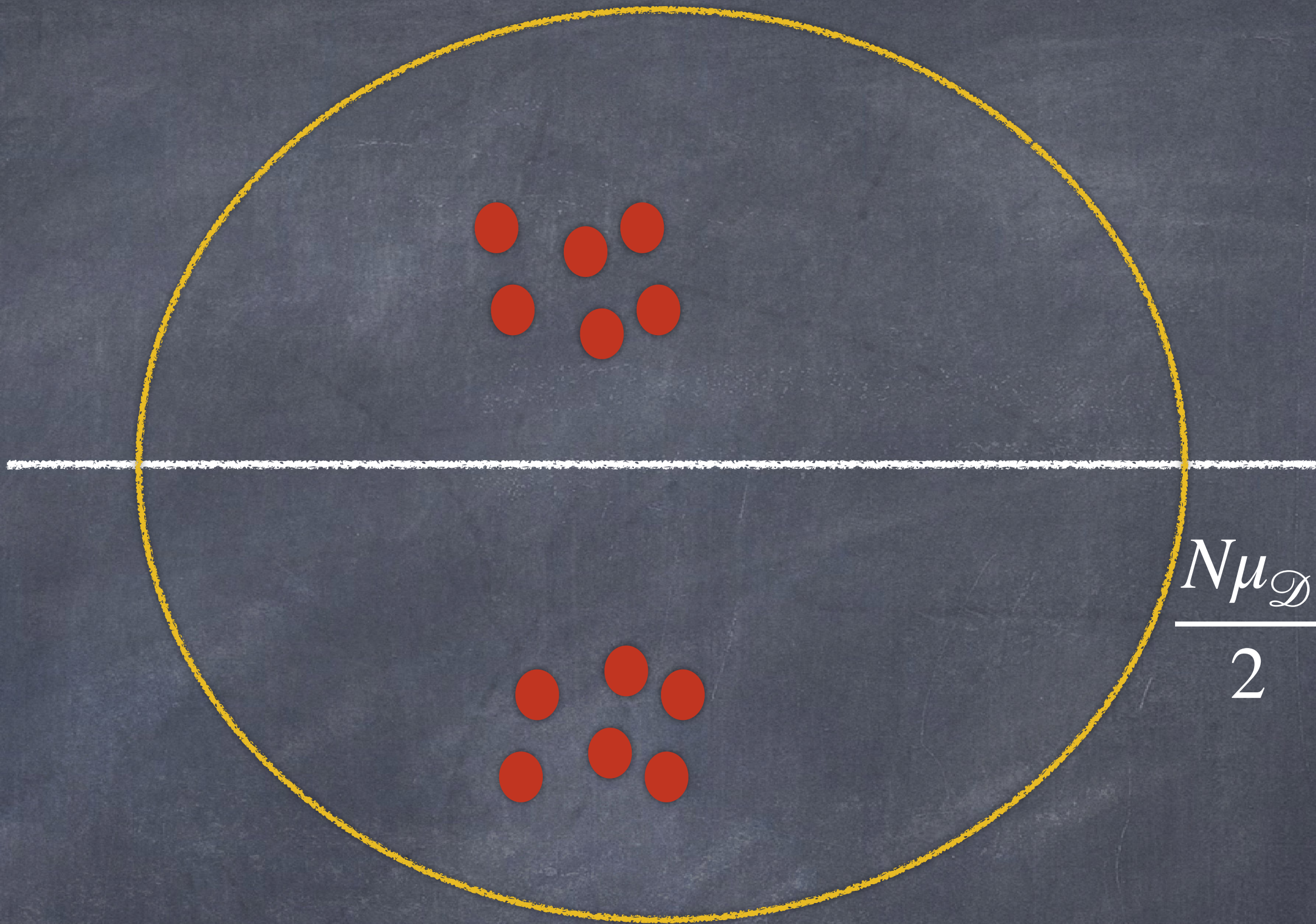
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Sync at the origin

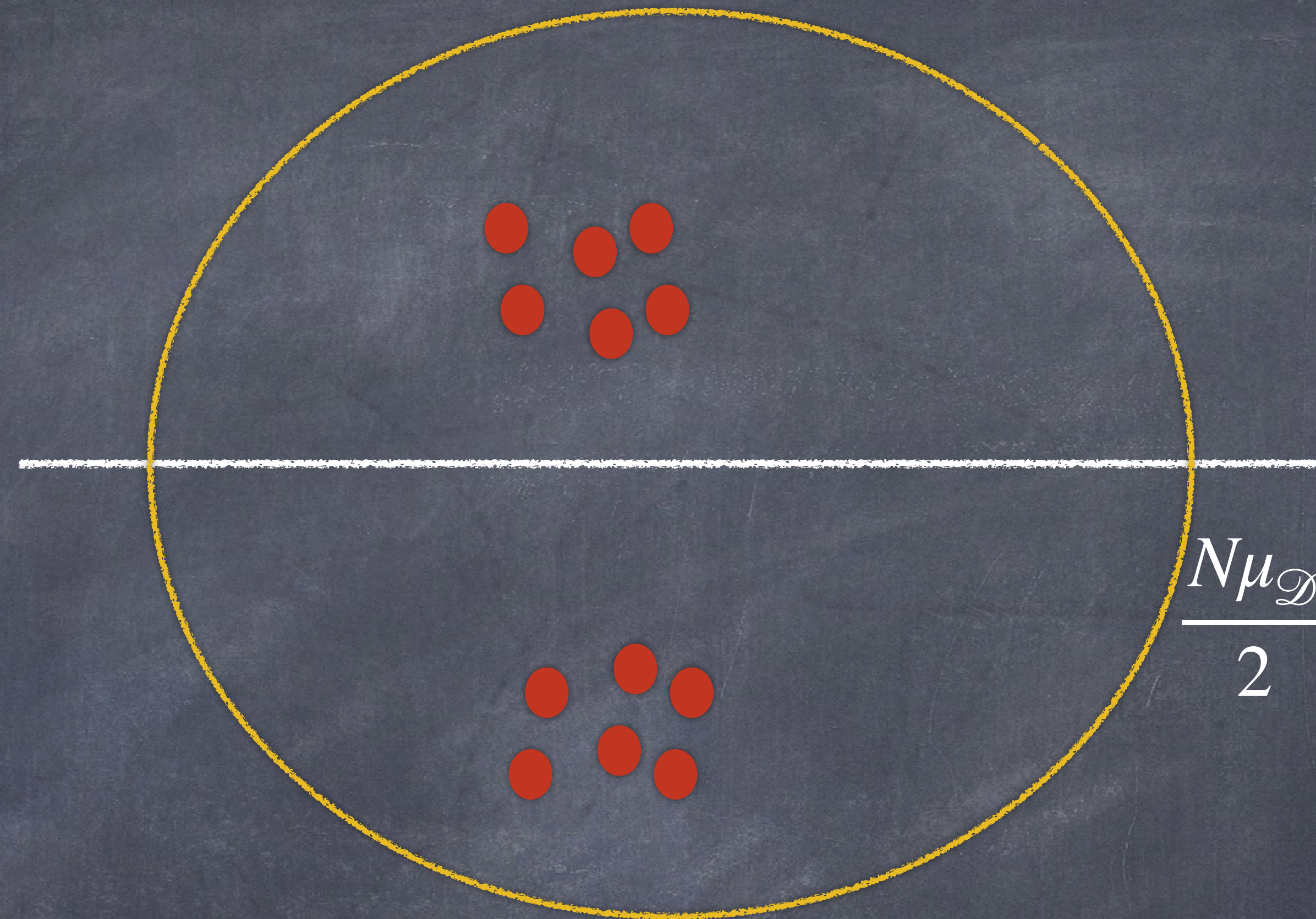
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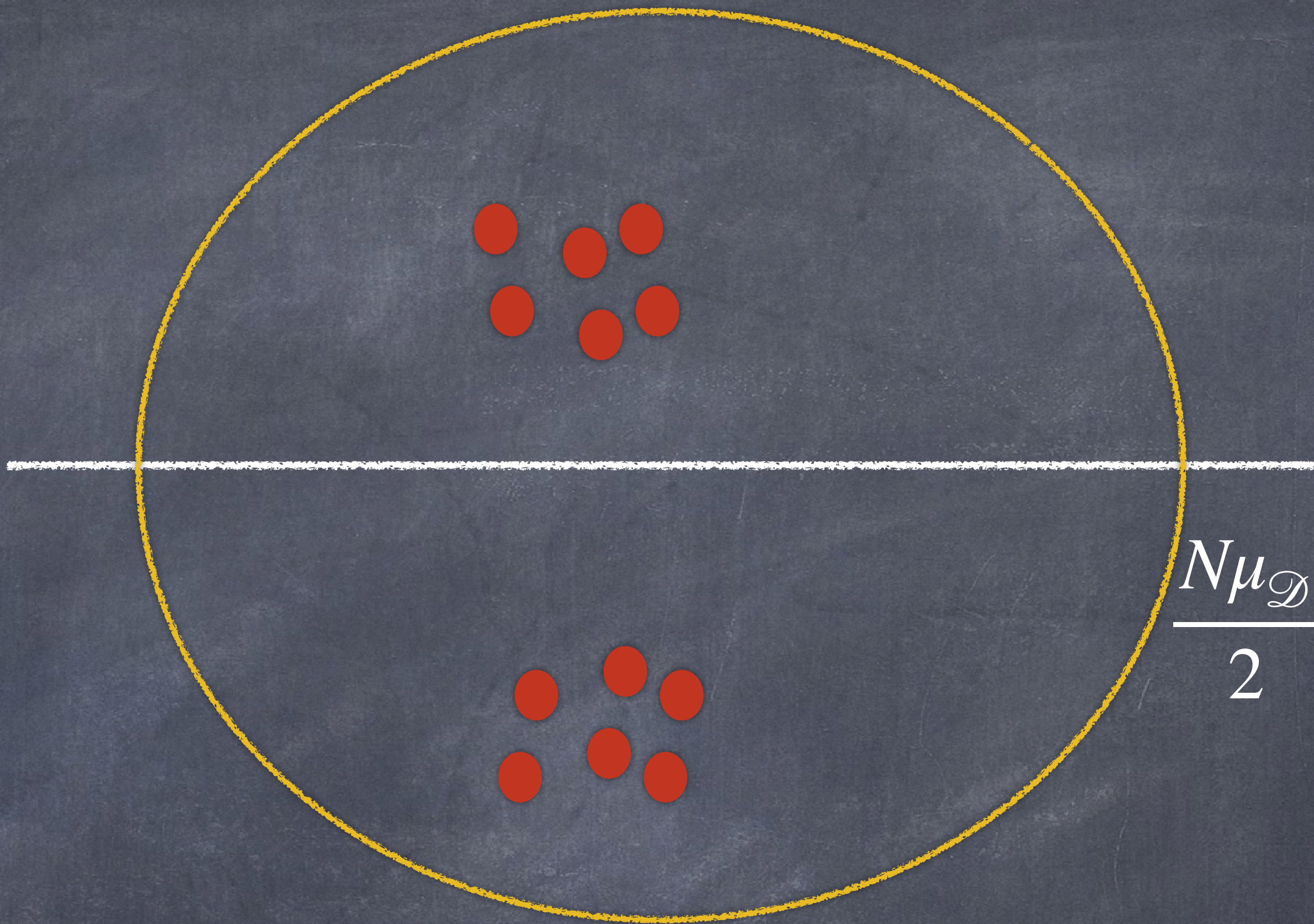
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$$A = \left\{ |\mu_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \left\{ |v_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \{\text{too technical}\}$$

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$$\mathbb{P}(A^c) < e^{-dN}$$

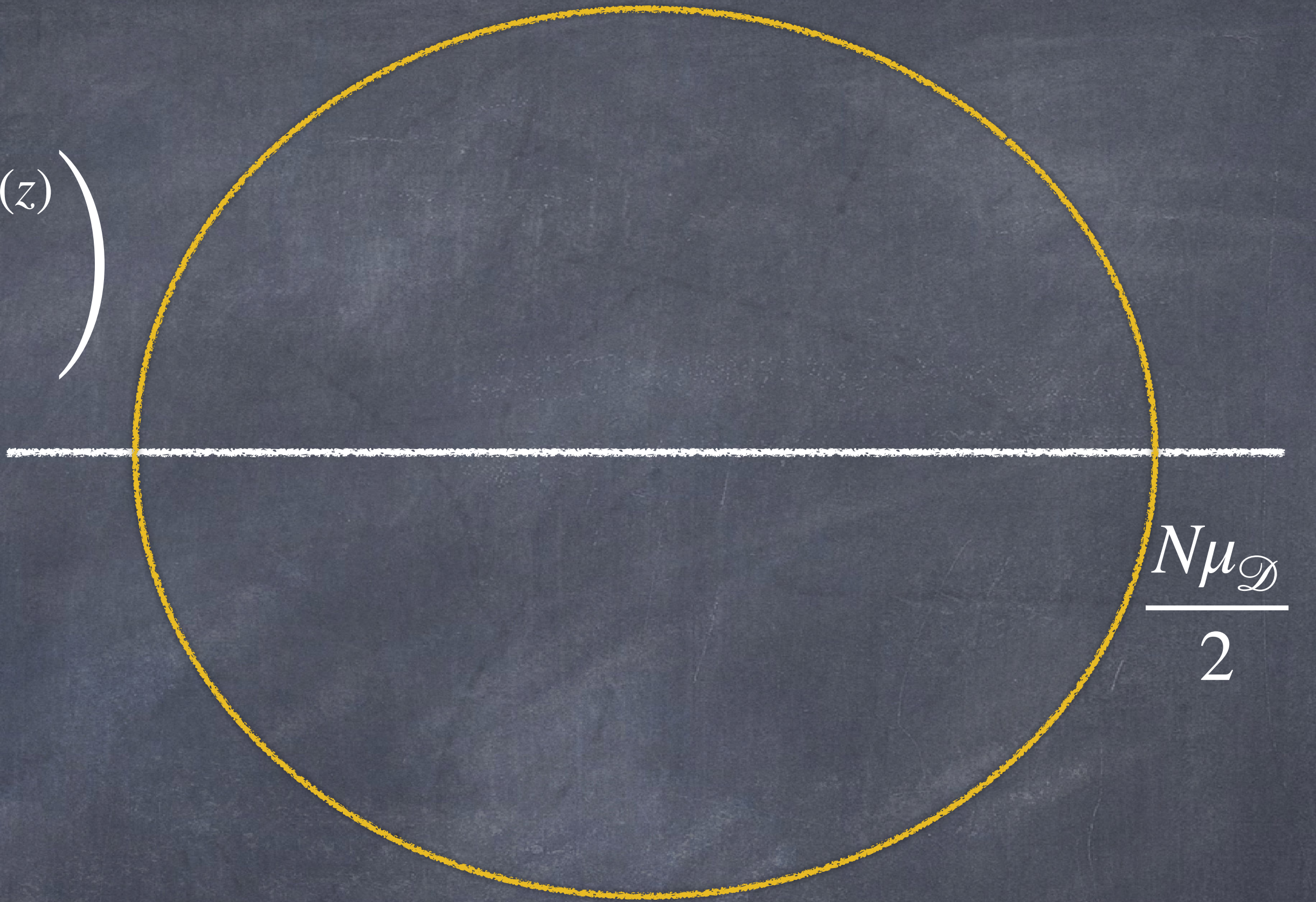
$$B(z)_+ = B(z)_- \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & a(z)e^{-2i\theta(z)} \\ -a^{-1}(z)e^{2i\theta(z)} & 1 \end{pmatrix}$$

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$$\frac{N\mu_{\mathcal{D}}}{2}$$

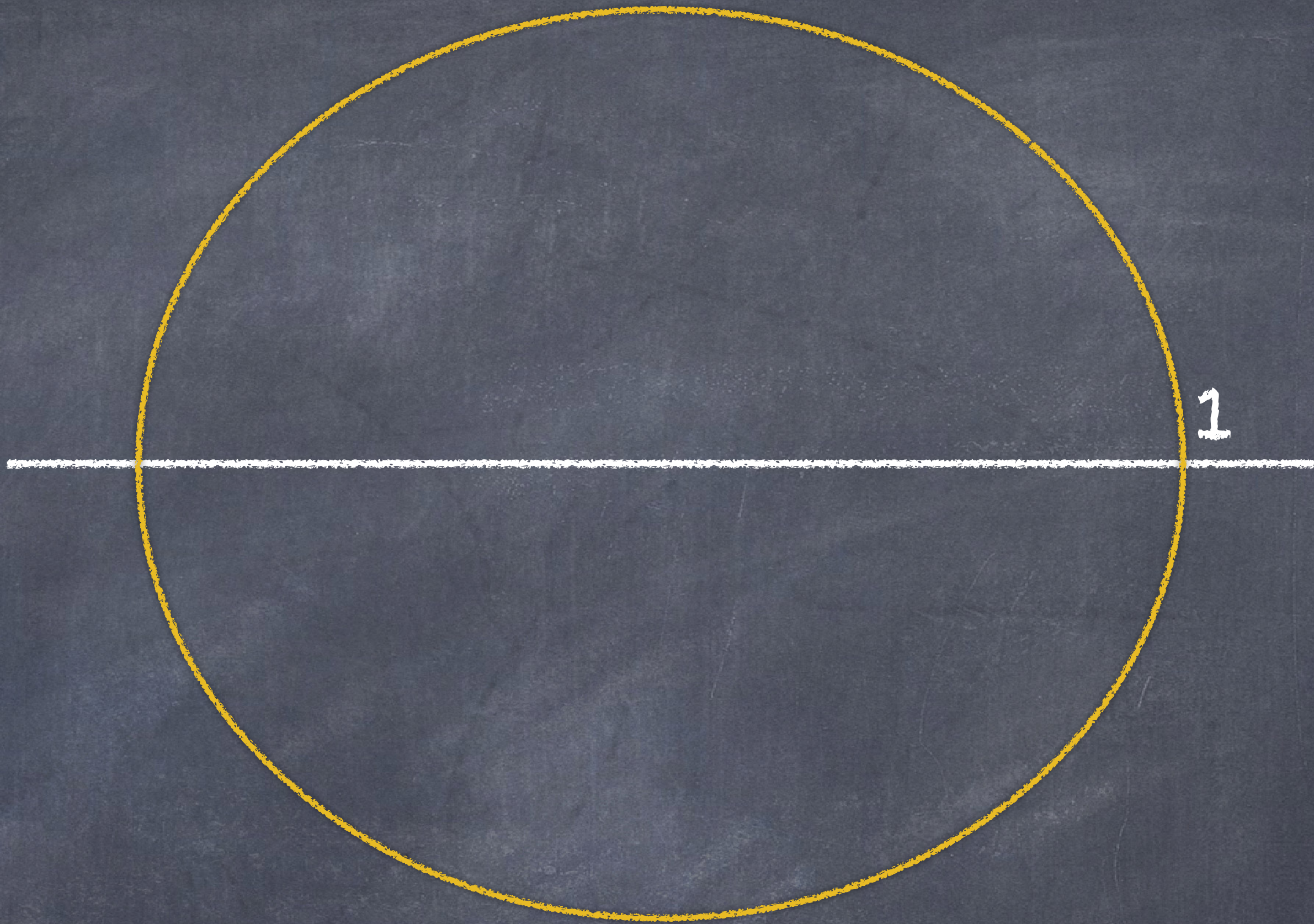
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$$z = \frac{N\mu_{\mathcal{D}}Z}{2}, \quad x = \frac{2X}{N\mu_{\mathcal{D}}}, \quad t = \frac{4T}{\mu_{\mathcal{D}}^2 N^2}.$$

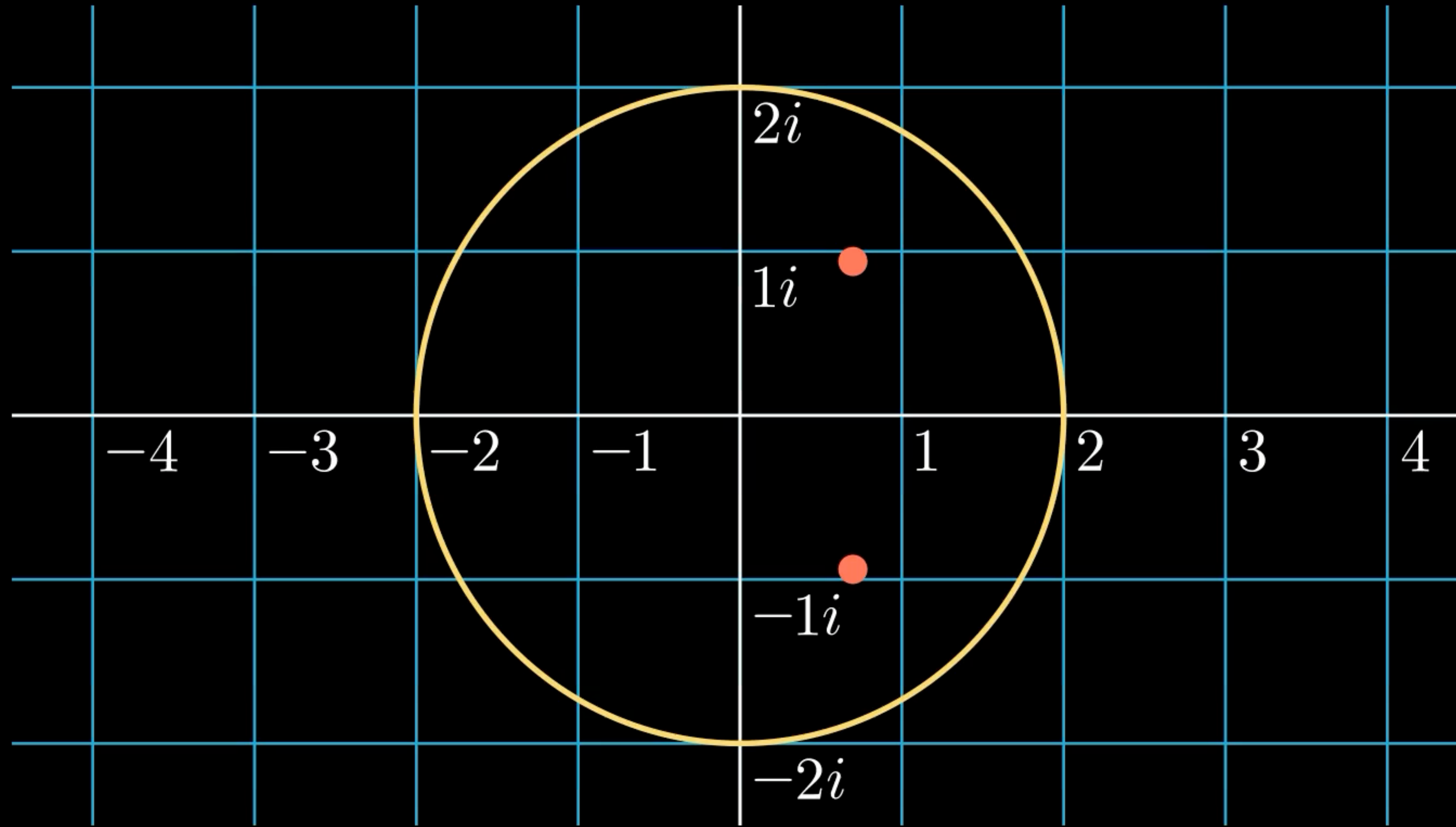
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$$R_+(Z) = R_-(Z) \frac{1}{\sqrt{2}} e^{-i(XZ + TZ^2 + 2Z^{-1})\sigma_3} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} e^{i(XZ + TZ^2 + 2Z^{-1})\sigma_3}$$

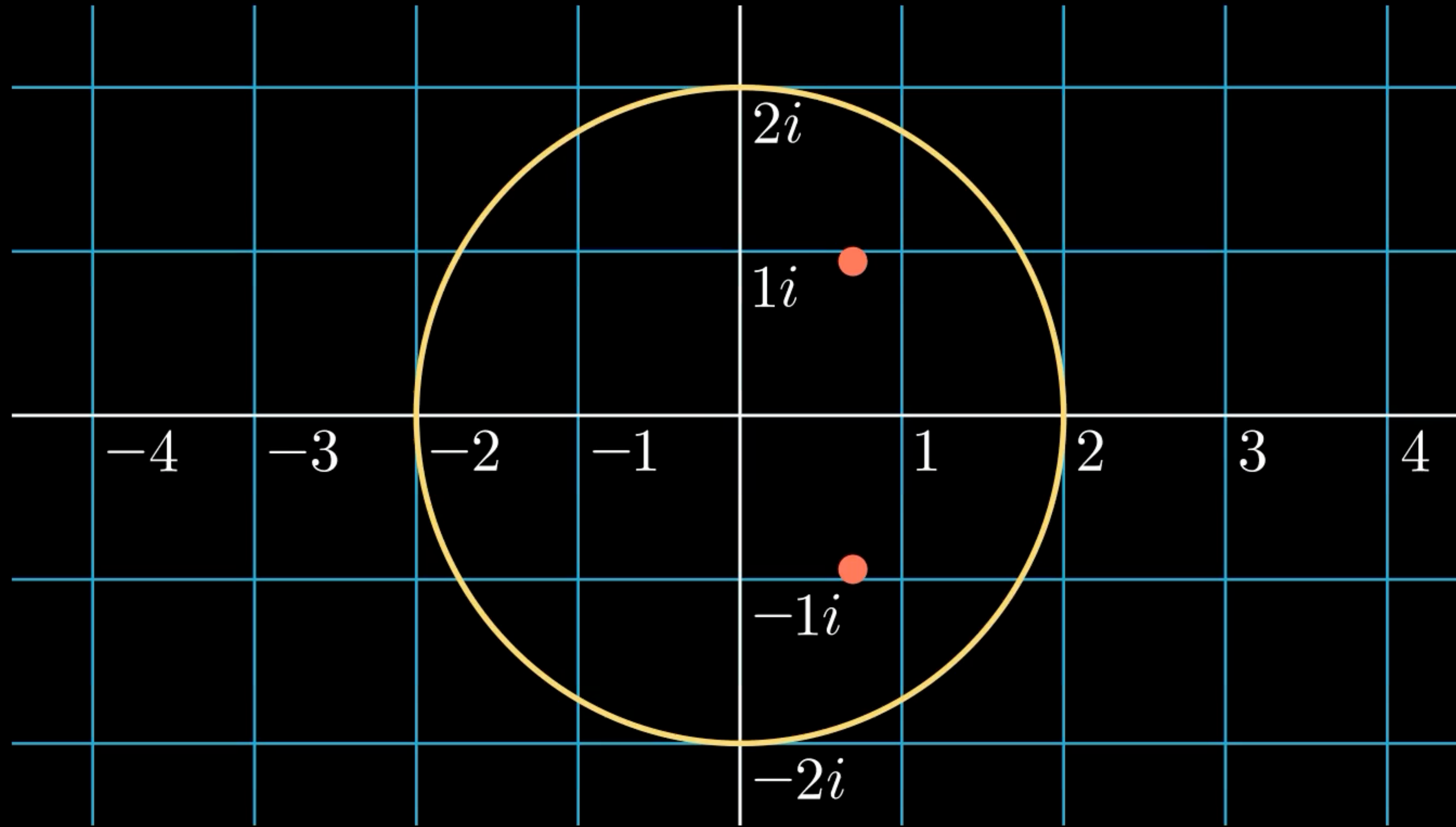
Number of Poles: 1

Order of the pole: 1



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We got lucky

$$R_+(Z) = R_-(Z) \frac{1}{\sqrt{2}} e^{-i(XZ + TZ^2 + 2Z^{-1})\sigma_3} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} e^{i(XZ + TZ^2 + 2Z^{-1})\sigma_3}$$

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15 March 2020

Extreme superposition: Rogue waves of infinite order and the Painlevé-III hierarchy

[Deniz Bilman](#), [Liming Ling](#), [Peter D. Miller](#)

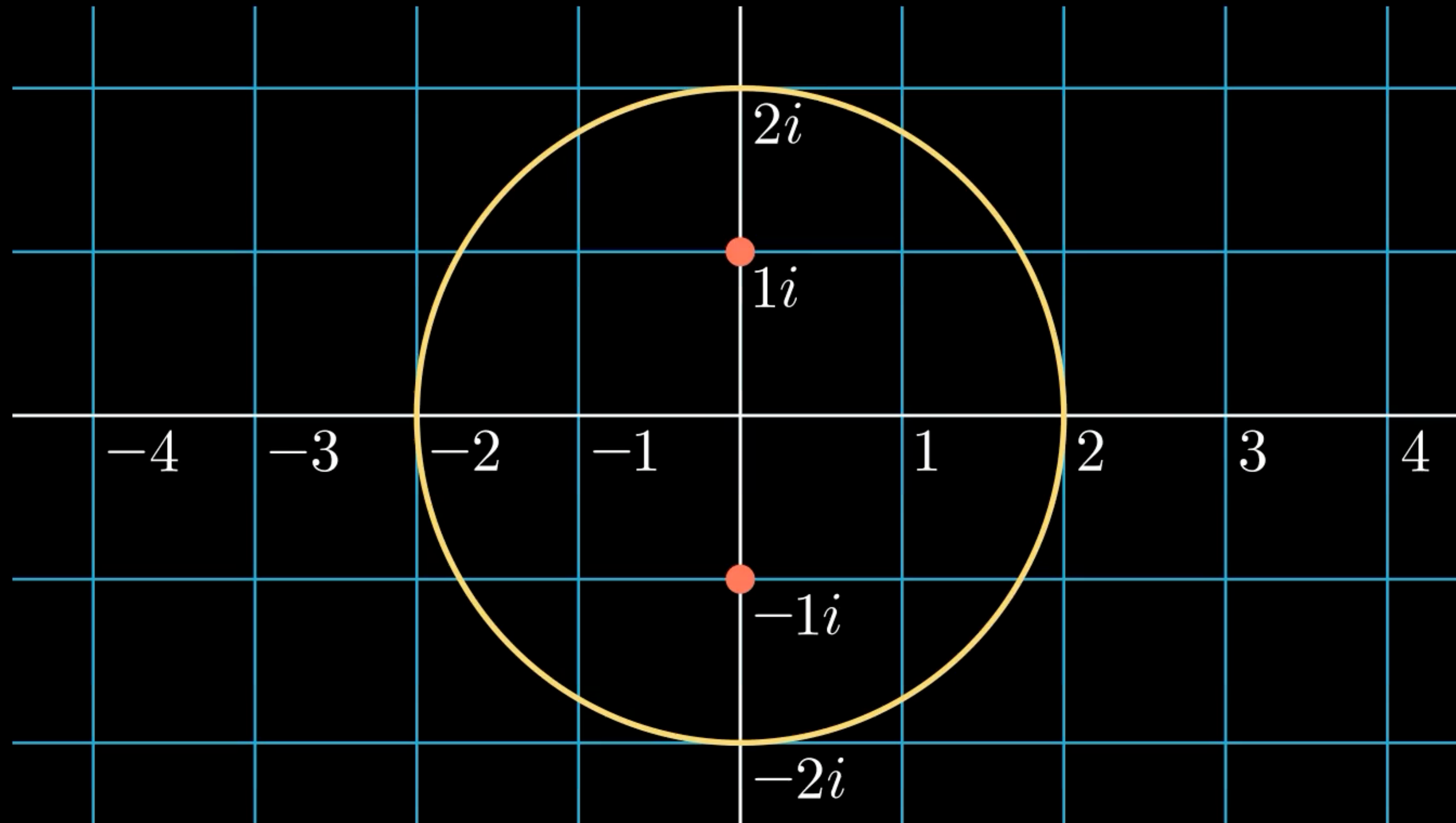
Duke Math. J. 169(4): 671-760 (15 March 2020). DOI: 10.1215/00127094-2019-0066

$$\Psi_{III} \left(-\frac{x^2}{8}, 0 \right) = \frac{\kappa}{x^2} \exp \left(\int_1^x \frac{2}{u(s)} ds \right)$$

The RWIO

Number of Poles: 1

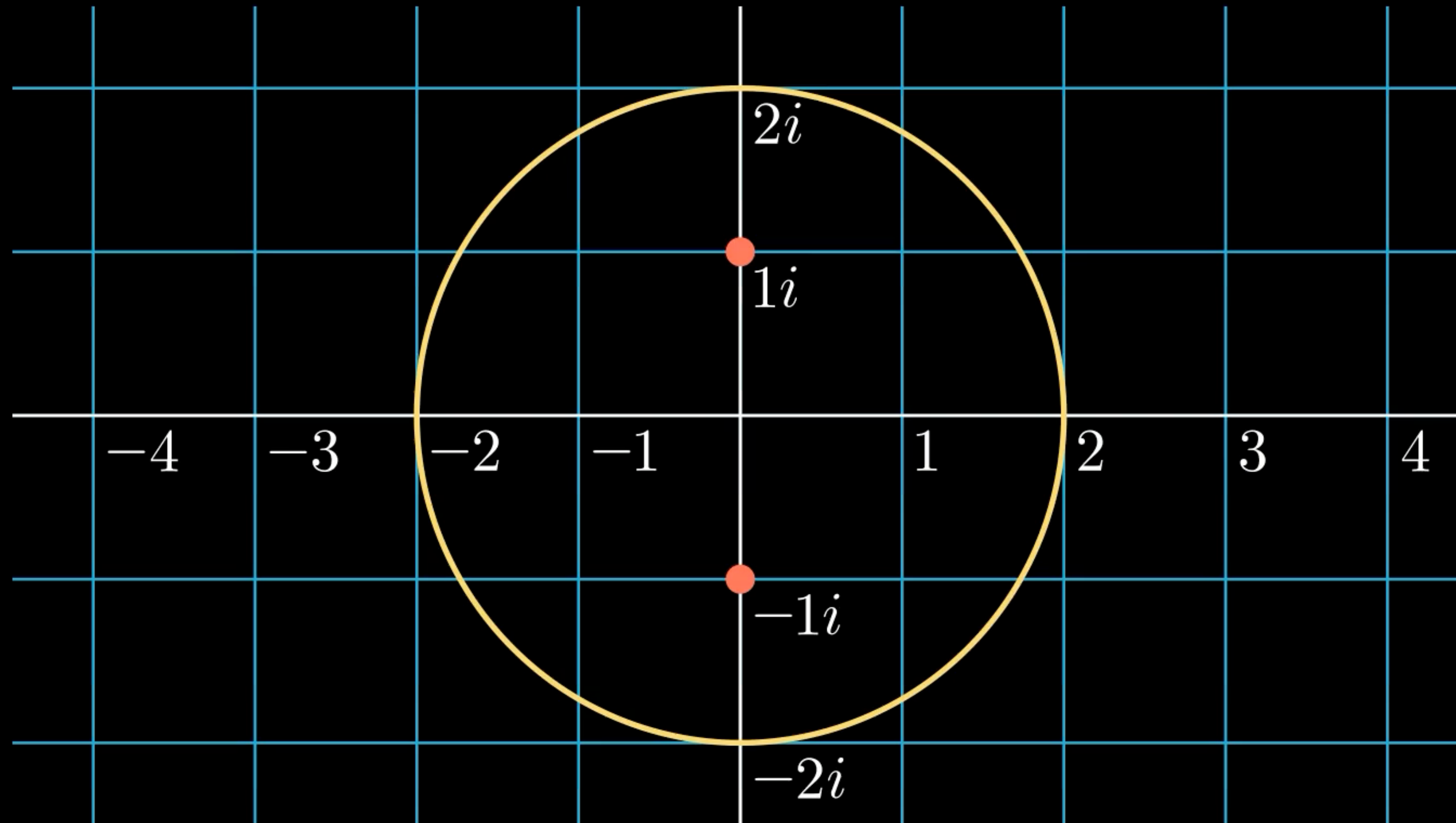
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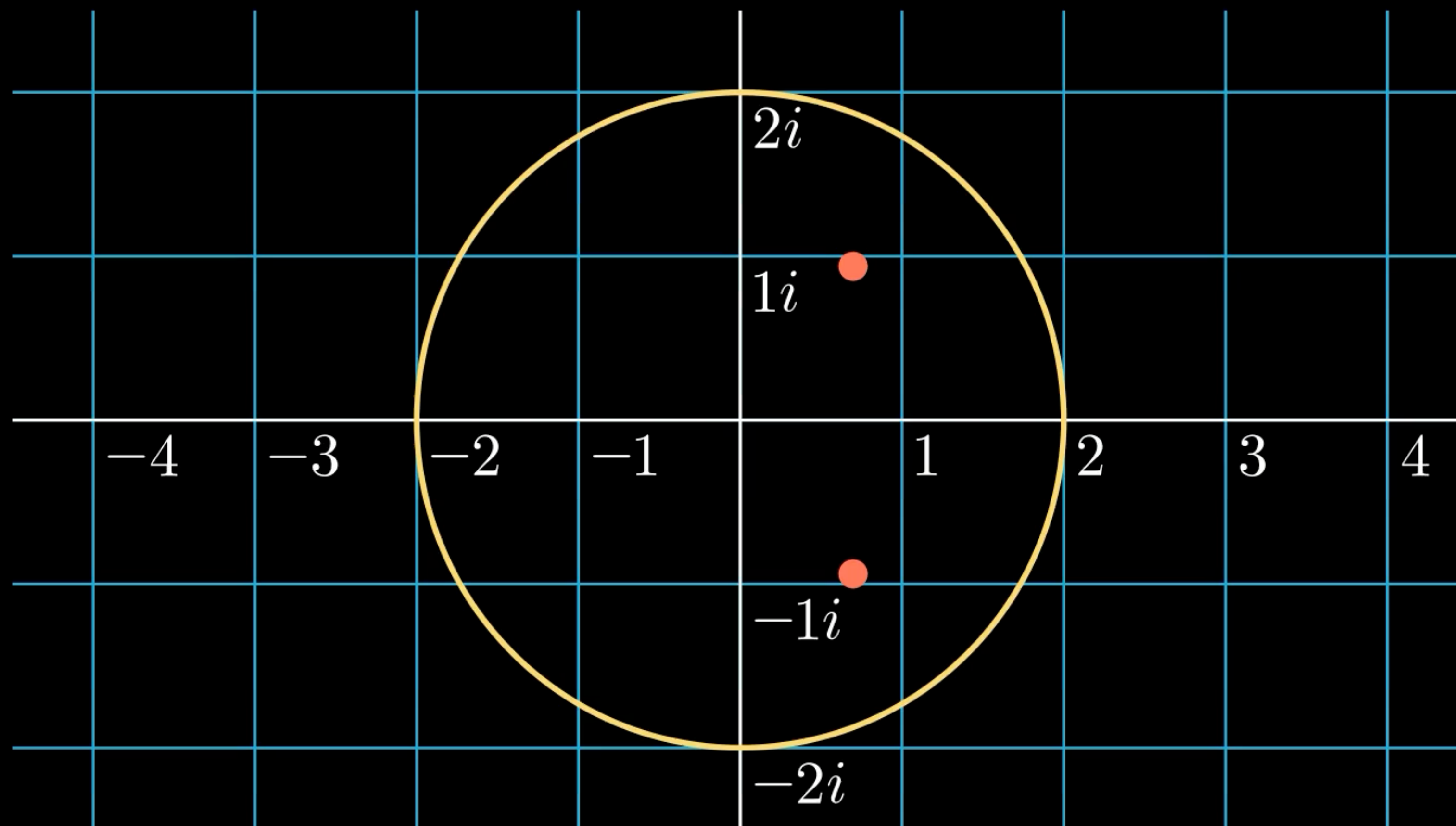
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The PIII-Rogue wave

Number of Poles: 1

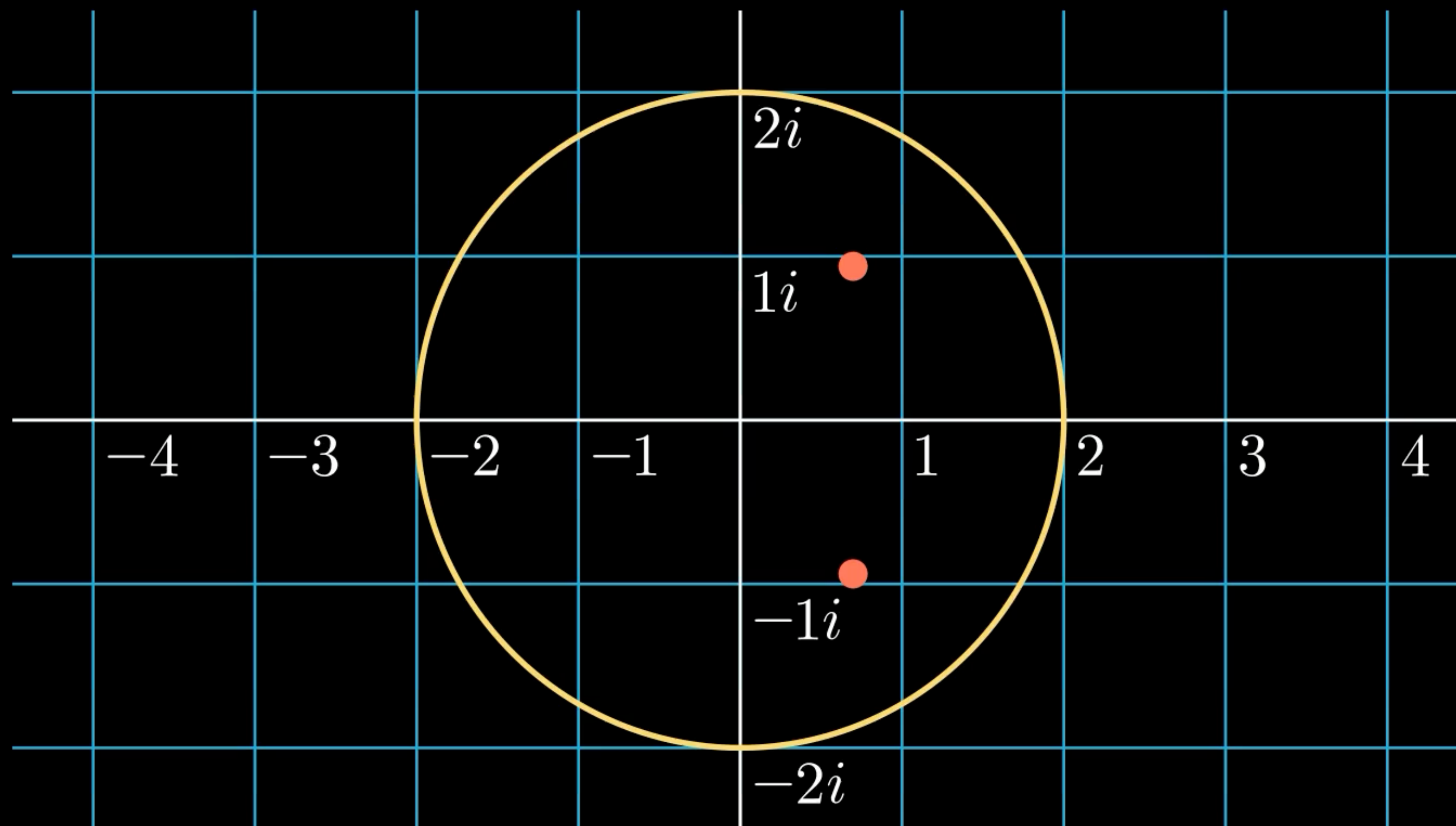
Order of the pole: 1



The PIII-Rogue wave

Number of Poles: 1

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$$A = \left\{ |\mu_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \left\{ |v_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \{\text{too technical}\}$$

$$\frac{2}{N\mu_{\mathcal{D}}} \psi \left(\frac{2X}{N\mu_{\mathcal{D}}}, \frac{4T}{\mu_{\mathcal{D}}^2 N^2} \right) - \Psi_{III}(X, T) \xrightarrow{N \rightarrow \infty} 0$$

$$A = \left\{ |\mu_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \left\{ |v_j| < \frac{\mu_{\mathcal{D}}}{4} N \right\} \cap \{\text{too technical}\}$$

$$\frac{2}{N\mu_{\mathcal{D}}} \psi \left(\frac{2X}{N\mu_{\mathcal{D}}}, \frac{4T}{\mu_{\mathcal{D}}^2 N^2} \right) - \Psi_{III}(X, T) \xrightarrow{N \rightarrow \infty} 0$$

$$\mathbb{E} \left[\left| \frac{2}{N\mu_{\mathcal{D}}} \psi \left(\frac{2X}{N\mu_{\mathcal{D}}}, \frac{4T}{\mu_{\mathcal{D}}^2 N^2} \right) - \Psi_{III}(X, T) \right| \right] \xrightarrow{N \rightarrow \infty} 0$$

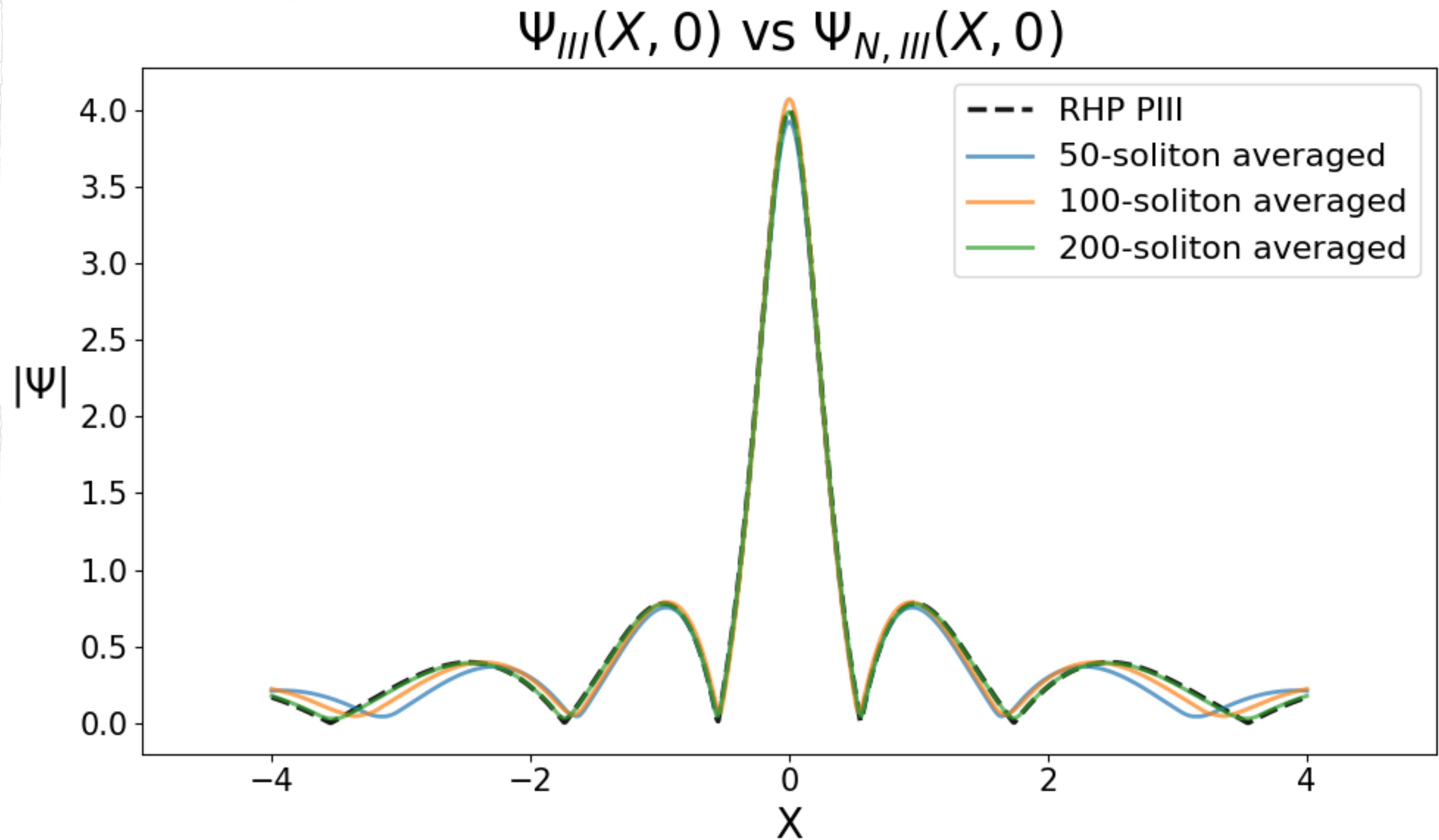
$$z_j = v_j + i\mu_j$$

Iid sub-exp

$$c_j = \frac{1}{a'(z_j)}$$

$$a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

$$\mathbb{E} \left[\left| \frac{2}{N\mu_{\mathcal{D}}} \psi \left(\frac{2X}{N\mu_{\mathcal{D}}}, \frac{4T}{\mu_{\mathcal{D}}^2 N^2} \right) - \Psi_{III}(X, T) \right| \right] \xrightarrow{N \rightarrow \infty} 0$$



$$\frac{2}{N\mu_{\mathcal{D}}}\psi\left(\frac{2X}{N\mu_{\mathcal{D}}}, \frac{4T}{\mu_{\mathcal{D}}^2 N^2}\right), \Psi_{III}(X, T)$$

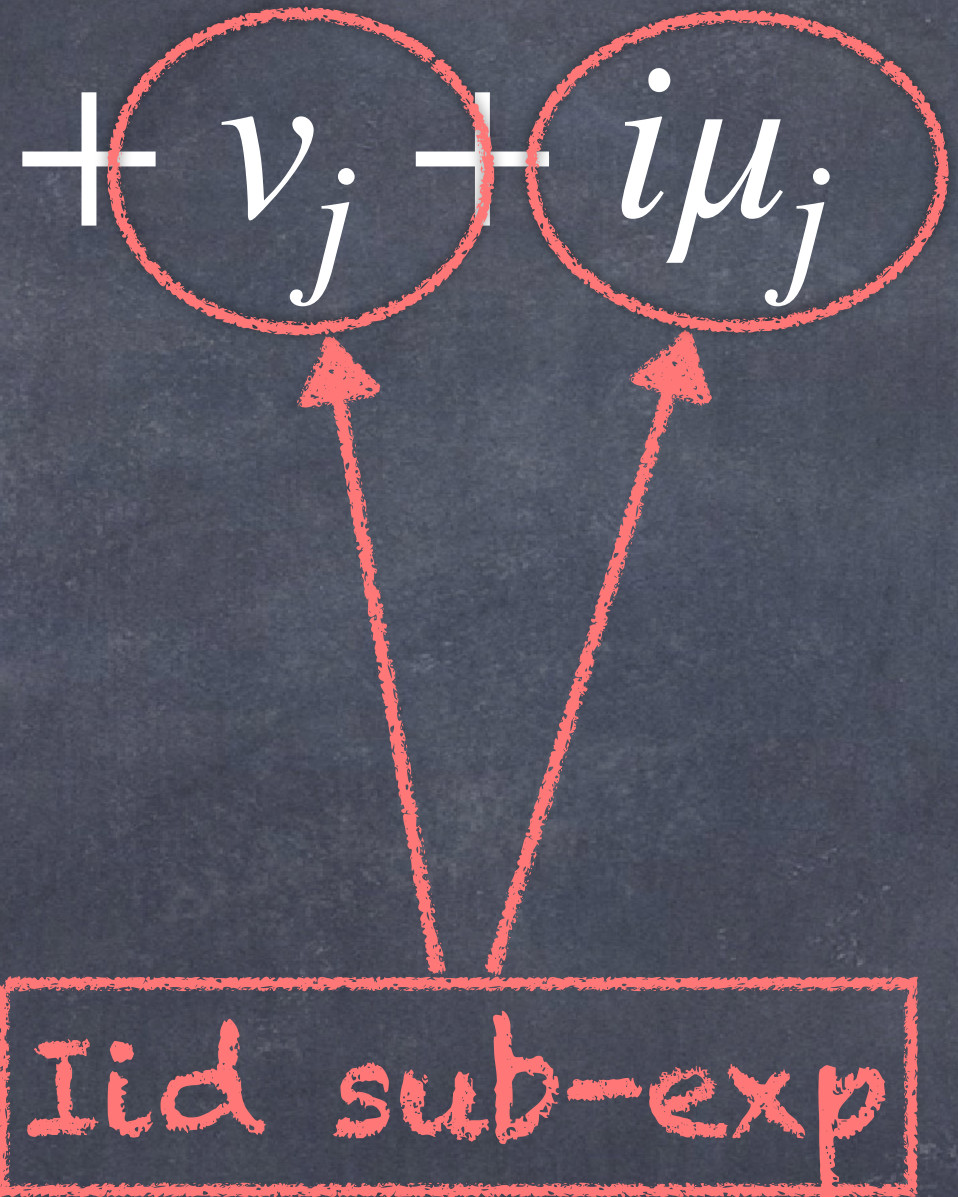
More exotic RW

$$z_j = -\zeta_j + v_j + i\mu_j \quad c_j = \frac{1}{a'(z_j)} \quad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

More exotic RW

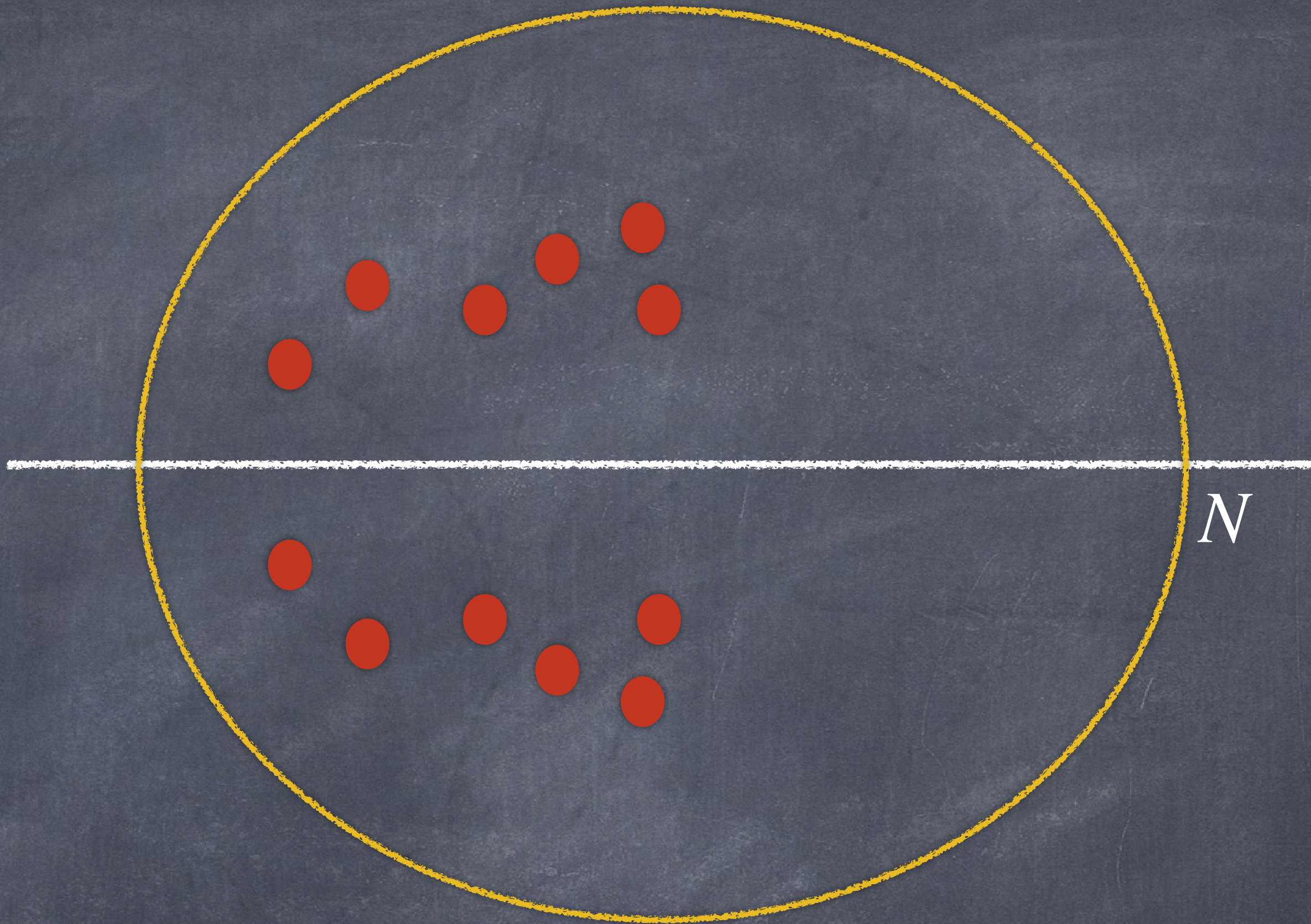
$$z_j = -\zeta_j + v_j + i\mu_j \quad c_j = \frac{1}{a'(z_j)} \quad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

Iid sub-exp



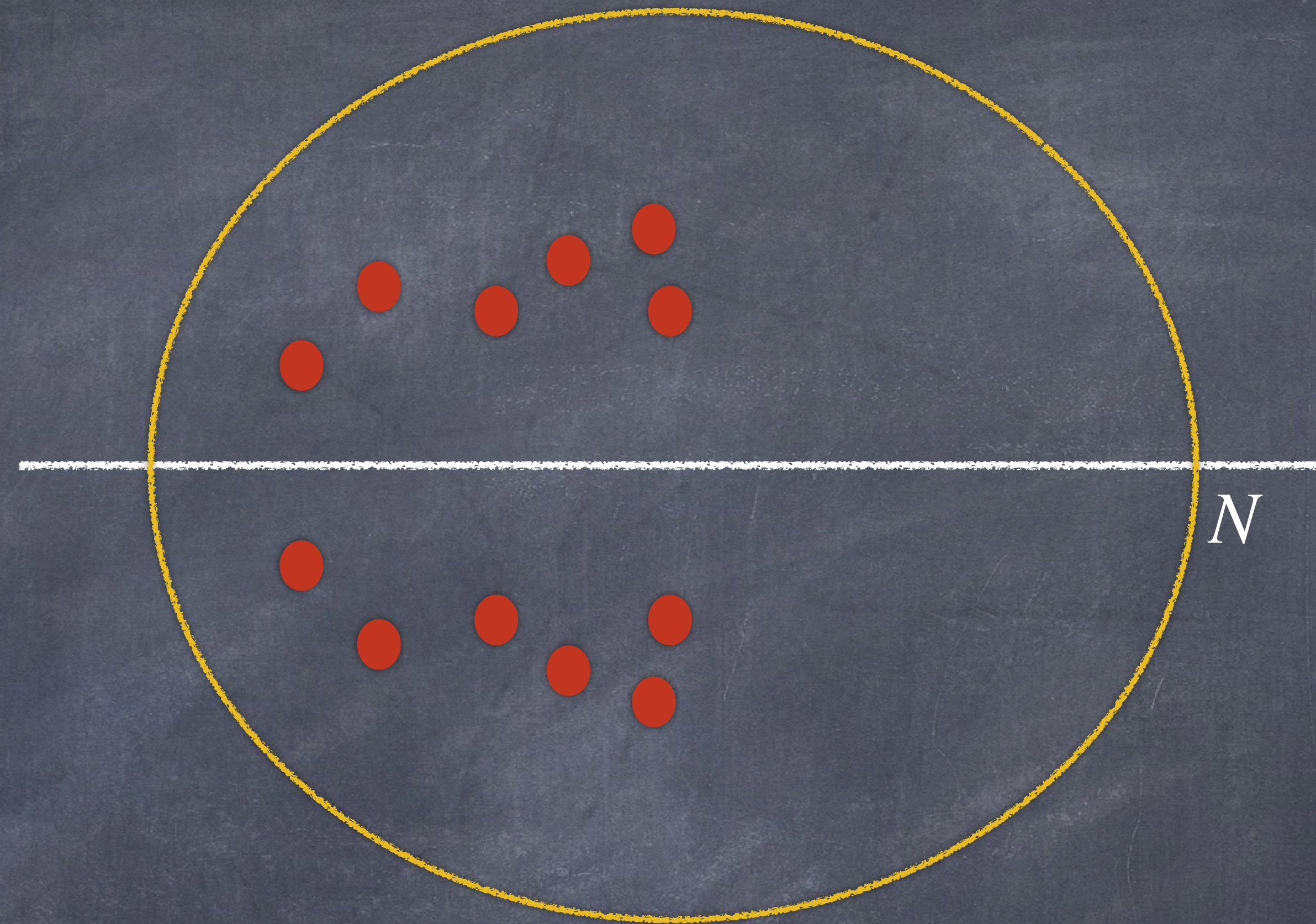
$$\operatorname{Res}_{z_j} M(z) = \lim_{z \rightarrow z_j} c_j e^{2i\theta(z,x,t)} M(z) \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$\operatorname{Res}_{\bar{z}_j} M(z) = \lim_{z \rightarrow \bar{z}_j} -\bar{c}_j e^{-2i\theta(z,x,t)} M(z) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$



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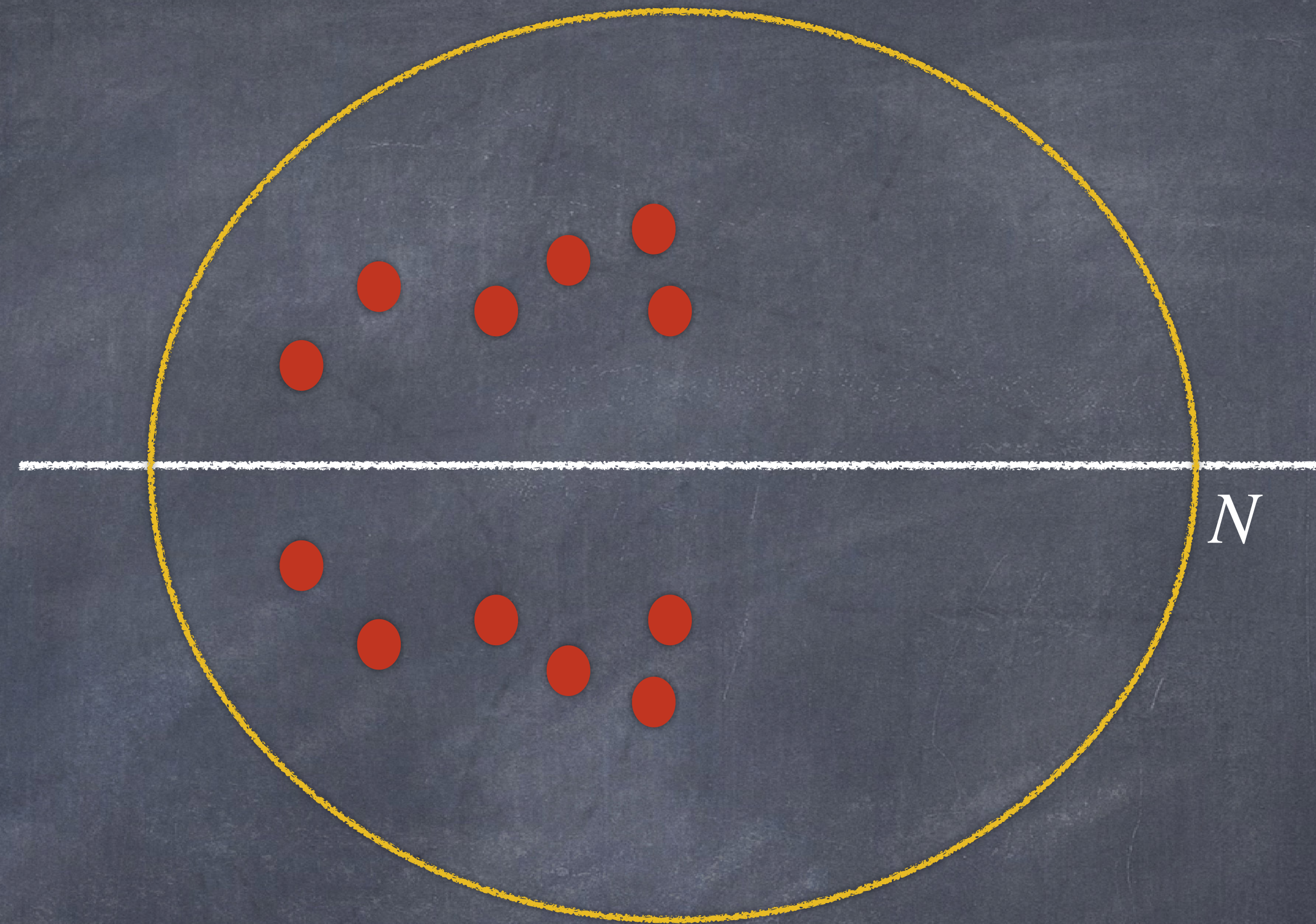
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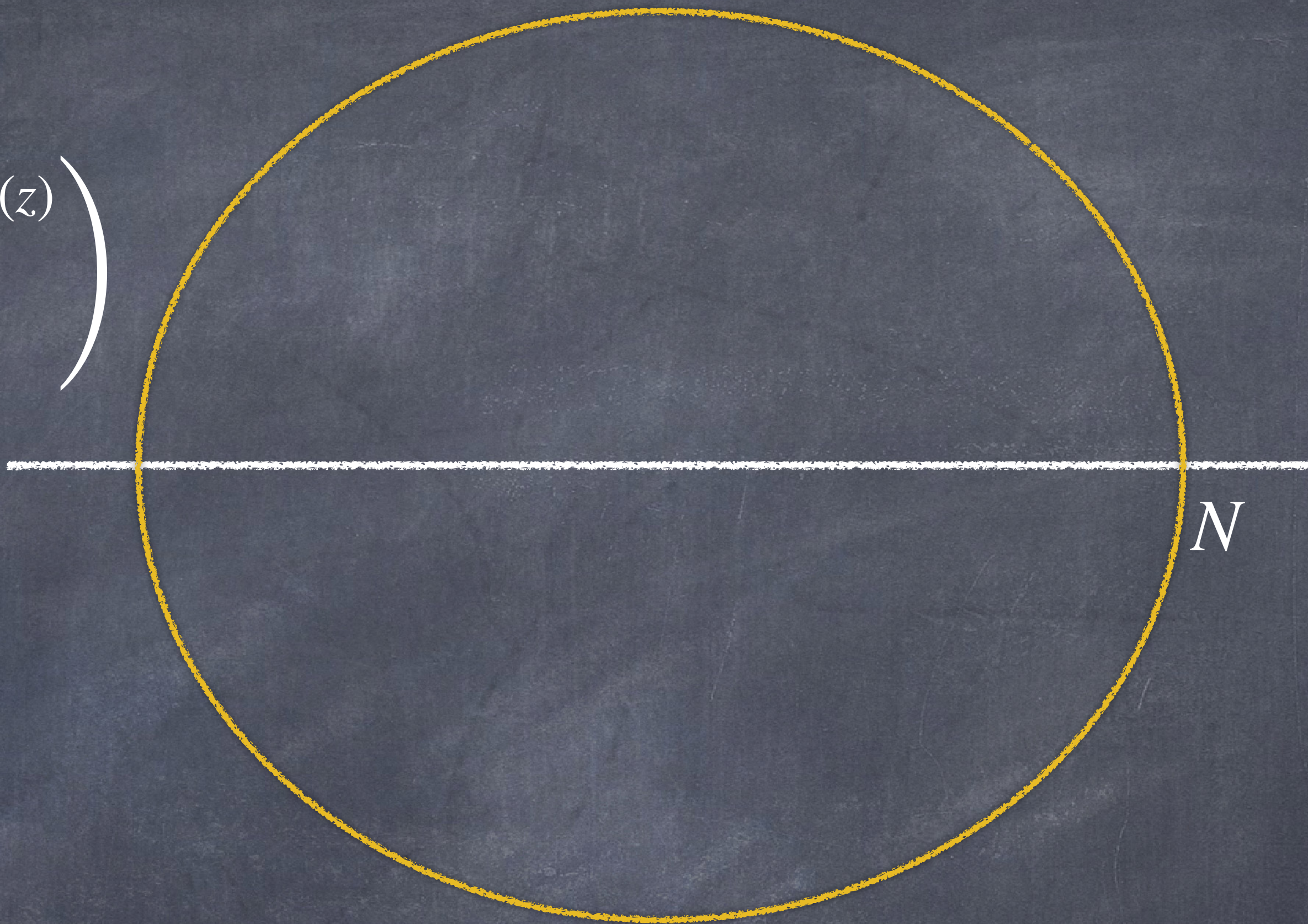


$$A = \{ |\mu_j| < N \} \cap \{ |v_j| < N \} \cap \{ \text{too technical} \}$$

$$\mathbb{P}(A^c) < e^{-dN}$$

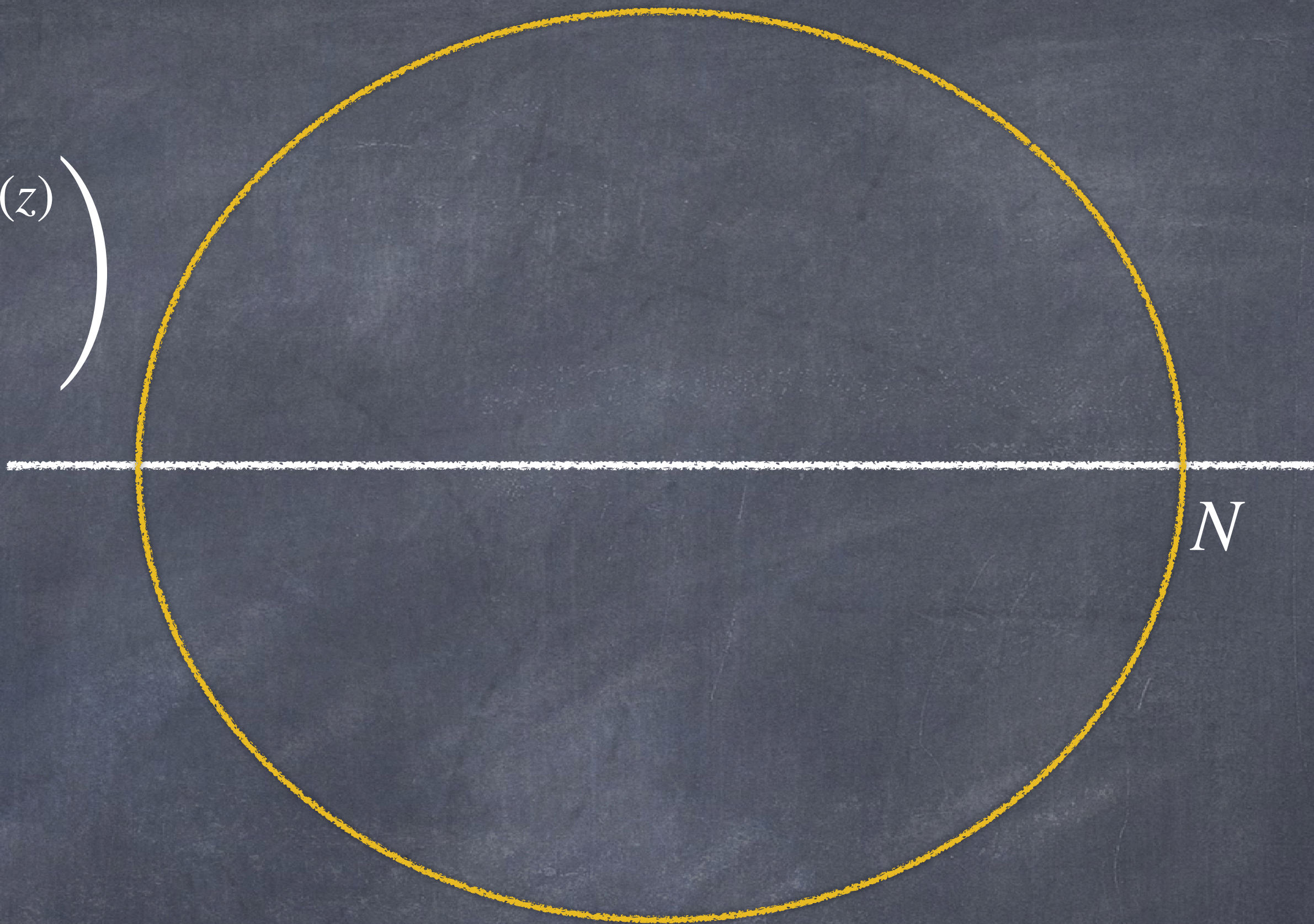
$$B(z)_+ = B(z)_- \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & a(z)e^{-2i\theta(z)} \\ -a^{-1}(z)e^{2i\theta(z)} & 1 \end{pmatrix}$$

$$a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

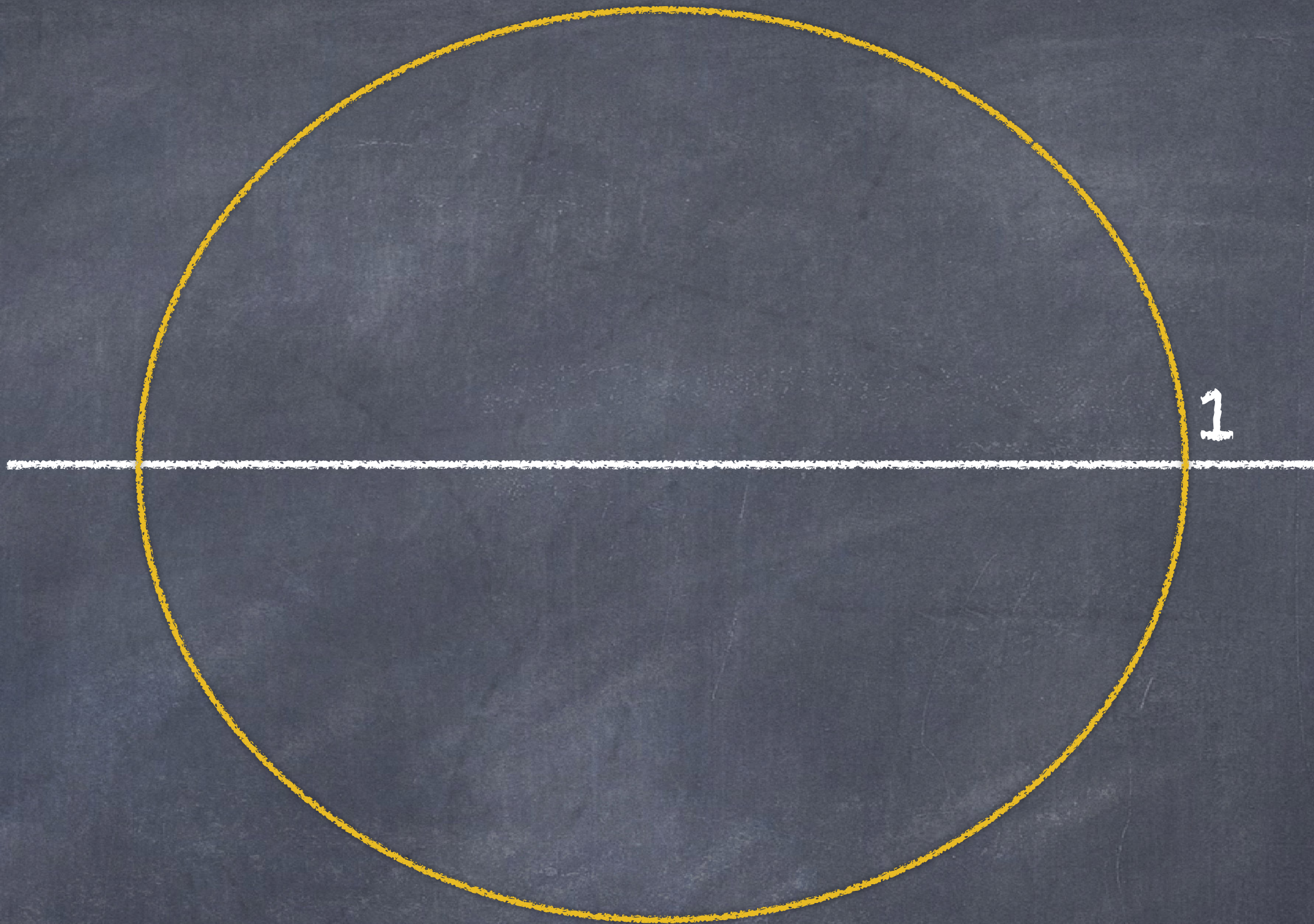


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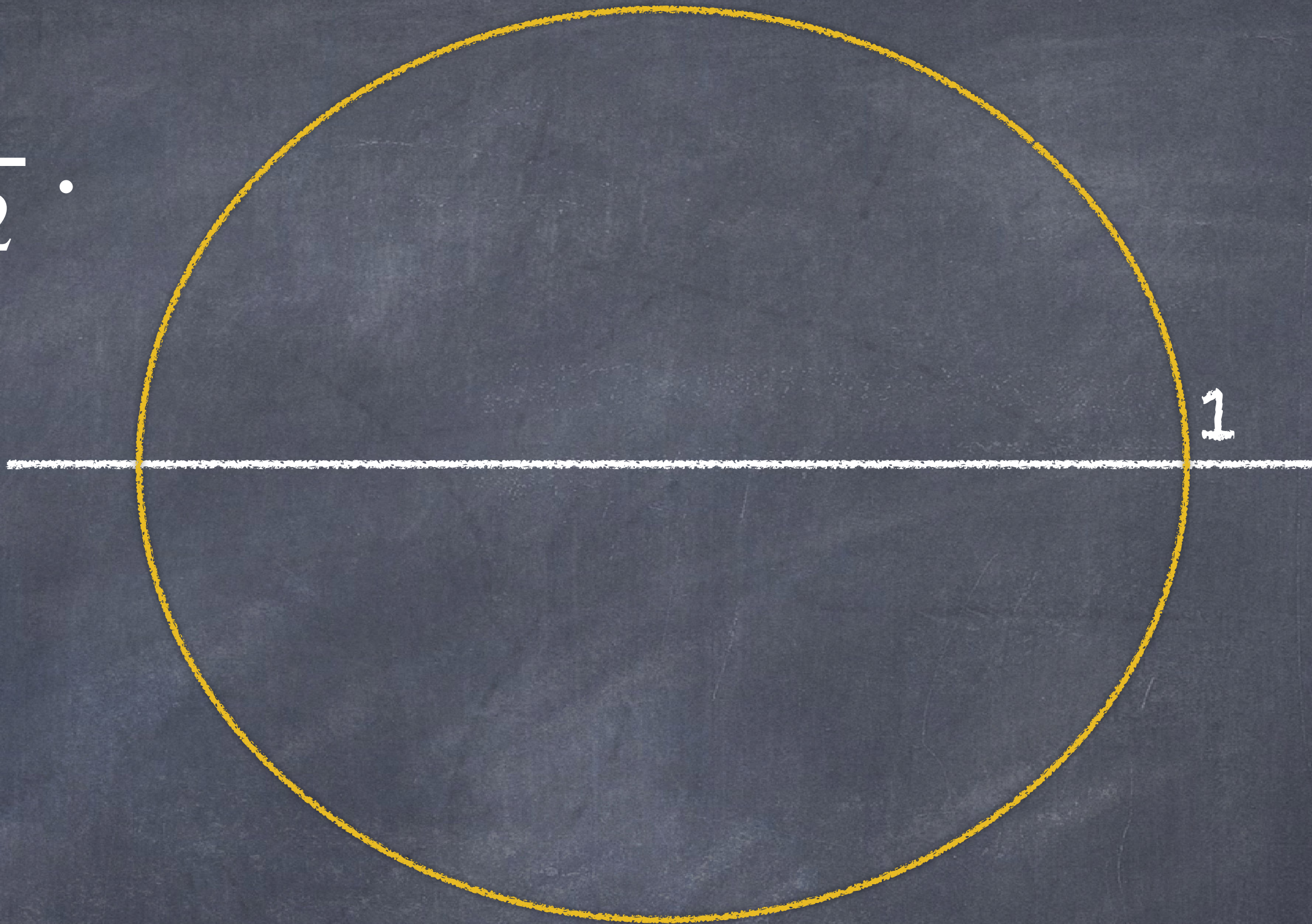


$$z = NZ, \quad x = \frac{X}{N}, \quad t = \frac{T}{N^2}.$$



$$R_+(Z) = R_-(Z) \frac{1}{\sqrt{2}} e^{-i \left(XZ + TZ^2 + \frac{\mu \mathcal{D}}{\xi} \ln \left(\frac{Z+\zeta}{Z} \right) \right) \sigma_3} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} e^{i \left(XZ + TZ^2 + \frac{\mu \mathcal{D}}{\xi} \ln \left(\frac{Z+\zeta}{Z} \right) \right) \sigma_3}$$

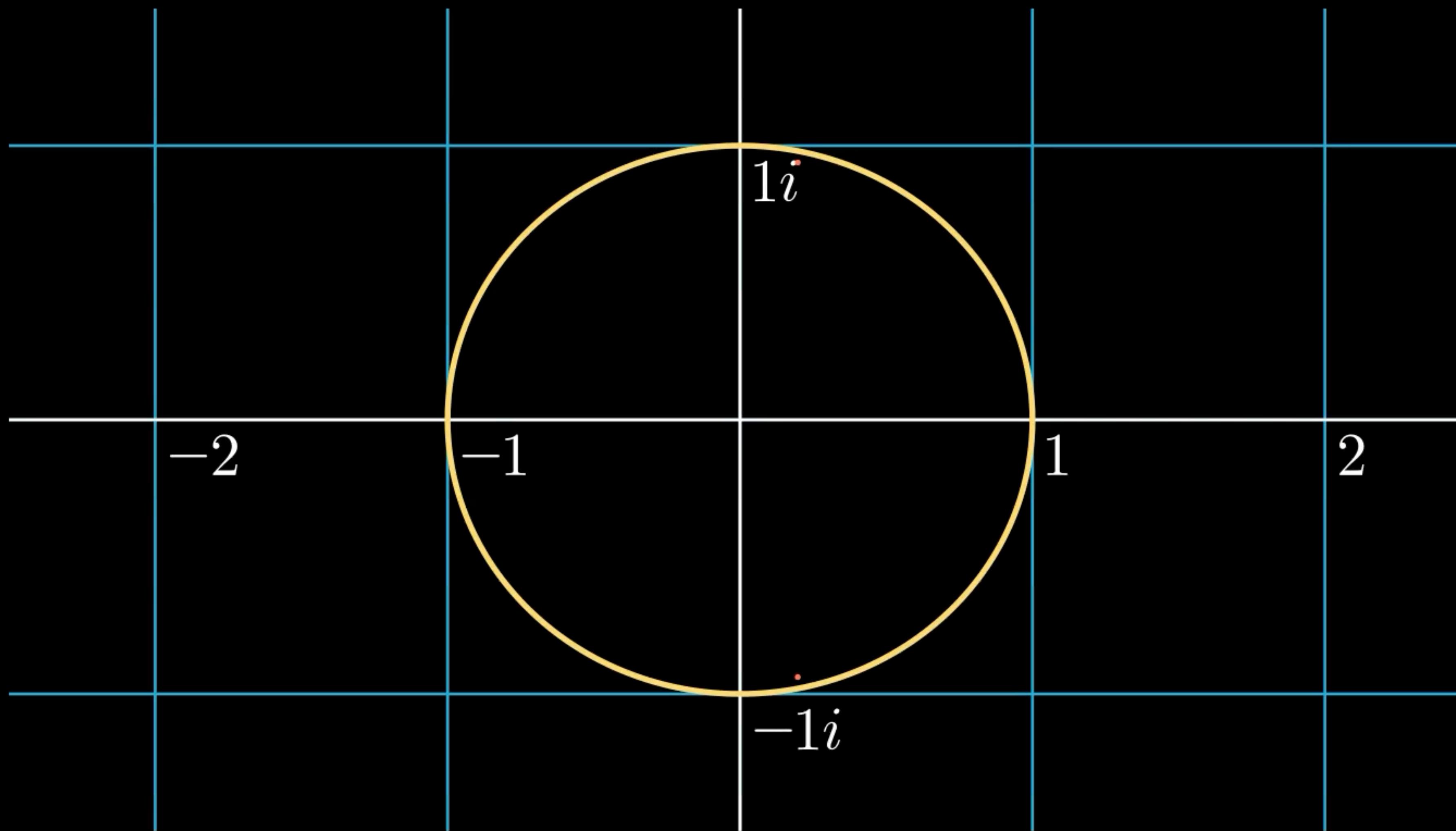
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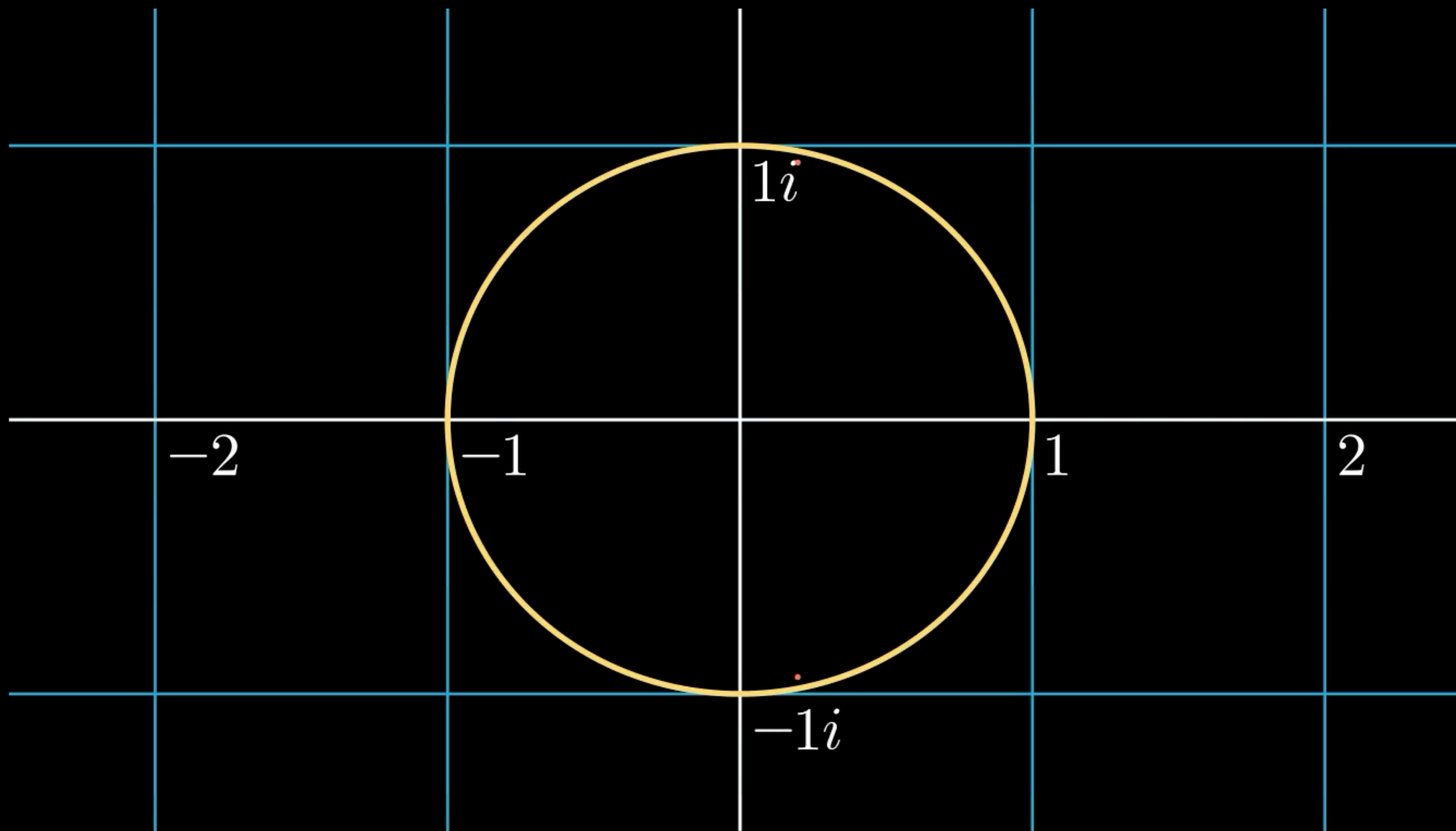
Number of poles: 1

Order of the poles: 1



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$$z_j = -\zeta j + v_j + i\mu_j \quad c_j = \frac{1}{a'(z_j)} \quad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

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↑
"Explicit" solution in terms of PV

$$z_j = -\zeta j + v_j + i\mu_j \quad c_j = \frac{1}{a'(z_j)} \quad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

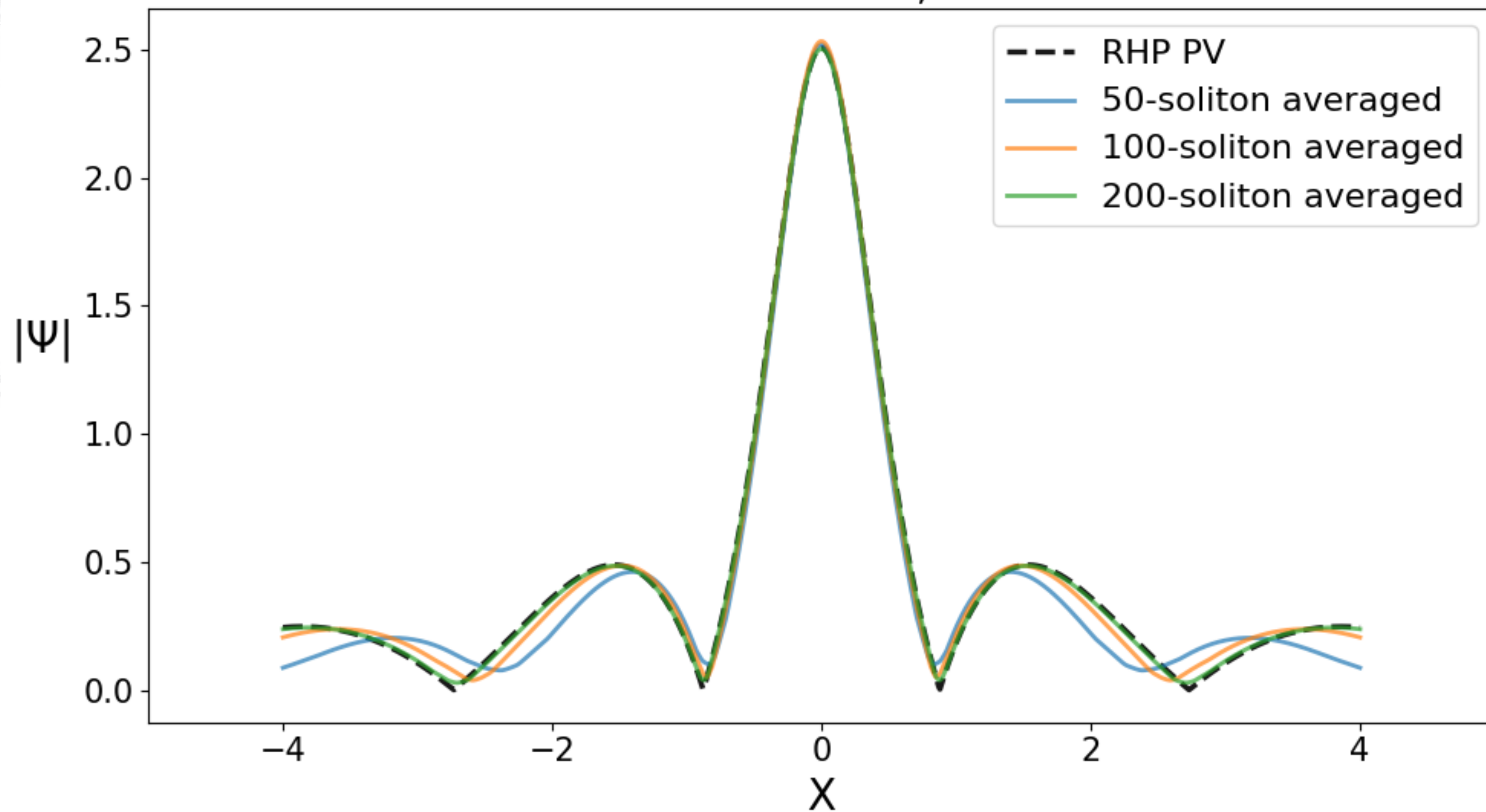
$$\mathbb{E} \left[\left| \frac{1}{N} \psi \left(\frac{X}{N}, \frac{T}{N^2} \right) - \Psi_V(X, T) \right| \right] \xrightarrow{N \rightarrow \infty} 0$$

$$z_j = -\zeta j + v_j + i\mu_j \quad c_j = \frac{1}{a'(z_j)} \quad a(z) = \prod_{j=1}^N \frac{z - z_j}{z - \bar{z}_j}$$

$$\mathbb{E} \left[\left| \frac{1}{N} \psi \left(\frac{X}{N}, \frac{T}{N^2} \right) - \Psi_V(X, T) \right| \right] \xrightarrow{N \rightarrow \infty} 0$$

$$\Psi_V(X,0) = \frac{\kappa}{X} \exp \left(-\frac{1}{2i\zeta} \int_1^{2i\zeta X} \frac{u(s)}{1-u(s)} ds \right)$$

$\Psi_V(X, 0)$ vs $\Psi_{N,V}(X, 0)$



Something to try

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- Describe an order 1 open set

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- Random norming constants

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- Describe an order 1 open set
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- True Probabilistic RHP analysis

Thanks!