

A Global Carleman inequality for the Oseen system

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Introduction

Notation

Ω smooth bounded domain of $\mathbb{R}^d$, $d = 2, 3$,

$\omega \subset \Omega$ nonempty open set, $T > 0$.

Motivation: local exact controllability of the Navier-Stokes system

$$\left\{ \begin{array}{ll} \partial_t \tilde{v} - \Delta \tilde{v} + (\tilde{v} \cdot \nabla) \tilde{v} + \nabla \tilde{p} = u \mathbf{1}_\omega & \text{in } (0, T) \times \Omega, \\ \operatorname{div} \tilde{v} = 0 & \text{in } (0, T) \times \Omega, \\ \tilde{v} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{v}(0, \cdot) = \tilde{v}^0 & \text{in } \Omega. \end{array} \right.$$

Controllability to the trajectories

Motivation

Let us consider a trajectory

$$\left\{ \begin{array}{ll} \partial_t \bar{v} - \Delta \bar{v} + (\bar{v} \cdot \nabla) \bar{v} + \nabla \bar{p} = 0 & \text{in } (0, T) \times \Omega, \\ \operatorname{div} \bar{v} = 0 & \text{in } (0, T) \times \Omega, \\ \bar{v} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \bar{v}(0, \cdot) = \bar{v}^0 & \text{in } \Omega. \end{array} \right.$$

We want to find $u \in L^2((0, T) \times \omega)$ such that

$$\tilde{v}(T, \cdot) = \bar{v}(T, \cdot).$$

Null-controllability

We set

$$v := \tilde{v} - \bar{v}, \quad p := \tilde{p} - \bar{p}.$$

The new system is

$$\begin{cases} \partial_t v - \Delta v + (\bar{v} \cdot \nabla)v + (v \cdot \nabla)\bar{v} + (v \cdot \nabla)v + \nabla p = u1_\omega & \text{in } (0, T) \times \Omega, \\ \operatorname{div} v = 0 & \text{in } (0, T) \times \Omega, \\ v = 0 & \text{on } (0, T) \times \partial\Omega, \\ v(0, \cdot) = \tilde{v}^0 - \bar{v}^0 & \text{in } \Omega. \end{cases}$$

Null-controllability

For $\|\tilde{v}^0 - \bar{v}^0\|$ “small”, find $u \in L^2((0, T) \times \omega)$ such that

$$v(T, \cdot) = 0.$$

Linearization

Linear system

Null-controllability of the **Oseen system**

$$\begin{cases} \partial_t v - \Delta v + (\bar{v} \cdot \nabla)v + (v \cdot \nabla)\bar{v} + \nabla p = \textcolor{red}{F} + u1_\omega & \text{in } (0, T) \times \Omega, \\ \operatorname{div} v = 0 & \text{in } (0, T) \times \Omega, \\ v = 0 & \text{on } (0, T) \times \partial\Omega, \\ v(0, \cdot) = v^0 & \text{in } \Omega. \end{cases}$$

Then a fixed-point argument with the map

$$\textcolor{blue}{F} \mapsto -(v \cdot \nabla)v$$

shows the local controllability of the nonlinear system.

The adjoint system

Definition

The **adjoint system** of the previous system is

$$\begin{cases} -\partial_t W - \Delta W - (\mathbb{D}W) \bar{v} + \nabla \pi_W = G & \text{in } (0, T) \times \Omega, \\ \operatorname{div} W = 0 & \text{in } (0, T) \times \Omega, \\ W = 0 & \text{on } (0, T) \times \partial\Omega, \\ W(T, \cdot) = W^0 & \text{in } \Omega. \end{cases}$$

$\mathbb{D}W$: symmetric gradient.

Change of variables $t \mapsto T - t$:

Our system

$$\begin{cases} \partial_t \tilde{w} - \Delta \tilde{w} + \nabla \tilde{\pi} = \tilde{f} + (\mathbb{D}\tilde{w}) a & \text{in } (0, T) \times \Omega, \\ \operatorname{div} \tilde{w} = 0 & \text{in } (0, T) \times \Omega, \\ \tilde{w} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{w}(0, \cdot) = \tilde{w}^0 & \text{in } \Omega. \end{cases}$$

Final state observability

To prove the null-controllability of the Oseen system, we can consider the adjoint system with $\tilde{f} = 0$:

$$\left\{ \begin{array}{ll} \partial_t \tilde{w} - \Delta \tilde{w} + \nabla \tilde{\pi} = (\mathbb{D} \tilde{w}) a & \text{in } (0, T) \times \Omega, \\ \operatorname{div} \tilde{w} = 0 & \text{in } (0, T) \times \Omega, \\ \tilde{w} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{w}(0, \cdot) = \tilde{w}^0 & \text{in } \Omega. \end{array} \right.$$

Hypotheses on a

$$a \in L^\infty(0, T; L^\infty(\Omega))^d.$$

$$\partial_t a \in L^2(0, T; L^\sigma(\Omega))^d \quad \text{with} \quad \left\{ \begin{array}{ll} \sigma > 1 & \text{if } d = 2, \\ \sigma > 6/5 & \text{if } d = 3. \end{array} \right.$$

Final state observability

$$\left\{ \begin{array}{ll} \partial_t \tilde{w} - \Delta \tilde{w} + \nabla \tilde{\pi} = (\mathbb{D}\tilde{w}) a & \text{in } (0, T) \times \Omega, \\ \operatorname{div} \tilde{w} = 0 & \text{in } (0, T) \times \Omega, \\ \tilde{w} = 0 & \text{on } (0, T) \times \partial\Omega, \\ \tilde{w}(0, \cdot) = \tilde{w}^0 & \text{in } \Omega. \end{array} \right.$$

Theorem

There exists $C > 0$ such that for any $T > 0$,

$$\int_{\Omega} |\tilde{w}(T, y)|^2 dy \leq C e^{C/T} \iint_{(0,T) \times \omega} |\tilde{w}|^2 dy dt.$$

Some references

- E. Fernández-Cara, S. Guerrero, O. Y. Imanuvilov, and J.-P. Puel. “Local exact controllability of the Navier-Stokes system”. *J. Math. Pures Appl.* (9) (2004).

Remark

Cost of the control of the form Ce^{C/T^4} .

Method : Carleman estimates but with different weights.

- F. W. Chaves-Silva and G. Lebeau. “Spectral inequality and optimal cost of controllability for the Stokes system”. *ESAIM Control Optim. Calc. Var.* (2016).

Remark

Cost of the control of the form $Ce^{C/T}$ on the Stokes system ($a = 0$).

Method : spectral inequality.

Our main result: the weight functions

$$\overline{\omega_0} \subset \omega.$$

$$\tilde{\psi} \in C^\infty(\mathbb{R}^d), \quad \tilde{\psi} = 0 \text{ on } \partial\Omega, \quad \tilde{\psi} \in (0, 1) \text{ in } \Omega,$$

$$\nabla \tilde{\psi}(y) = 0 \implies y \in \omega_0.$$

$$\ell(t) := t(T - t) \quad (t \in [0, T]).$$

$$\text{First parameter: } \lambda \geq \lambda_0.$$

Definition

We set

$$\tilde{\zeta}(t, y) := \frac{e^{\lambda \tilde{\psi}(y)}}{\ell(t)}, \quad \tilde{\varphi}(t, y) := \frac{e^{\lambda \tilde{\psi}(y)} - e^{2\lambda}}{\ell(t)} \quad (t \in (0, T), y \in \Omega).$$

Our main result: the weight functions

Remark

- These are the “classical” weight functions for the Carleman estimates of the heat equation.
- In FGIP2004, they use

$$\ell(t) := t^4(T - t)^4 \quad (t \in [0, T]).$$

Second parameter: $s \geq s_0(T + T^2)$.

We estimate

$$I_1(\tilde{w}) := \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} \left[\left(s\tilde{\zeta} \right)^{-1} \left(|\partial_t \tilde{w}|^2 + |\Delta \tilde{w}|^2 \right) \right. \\ \left. + s\tilde{\zeta} |\nabla \tilde{w}|^2 + \left(s\tilde{\zeta} \right)^3 |\tilde{w}|^2 \right] dy \, dt.$$

Our main result: the weight functions

We estimate

$$I_1(\tilde{w}) := \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} \left[\left(s\tilde{\zeta} \right)^{-1} \left(|\partial_t \tilde{w}|^2 + |\Delta \tilde{w}|^2 \right) \right. \\ \left. + s\tilde{\zeta} |\nabla \tilde{w}|^2 + \left(s\tilde{\zeta} \right)^3 |\tilde{w}|^2 \right] dy \, dt.$$

Remark

We have

$$\left(s\tilde{\zeta} \right)^3 e^{2s\tilde{\varphi}} = \left(\frac{se^{\lambda\tilde{\psi}(y)}}{t(T-t)} \right)^3 \exp \left(2s \frac{e^{\lambda\tilde{\psi}(y)} - e^{2\lambda}}{t(T-t)} \right),$$

with

$$\lambda \geq \lambda_0, \quad s \geq s_0(T + T^2).$$

Our main result: the weight functions

We estimate

$$I_1(\tilde{w}) := \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} \left[\left(s\tilde{\zeta} \right)^{-1} \left(|\partial_t \tilde{w}|^2 + |\Delta \tilde{w}|^2 \right) \right. \\ \left. + s\tilde{\zeta} |\nabla \tilde{w}|^2 + \left(s\tilde{\zeta} \right)^3 |\tilde{w}|^2 \right] dy \, dt$$

and

$$I_2(\tilde{\pi}) := \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} \left[|\nabla \tilde{\pi}|^2 + \left(s\tilde{\zeta} \right)^2 |\tilde{\pi}|^2 \right] dy \, dt.$$

Our main result: statement

Theorem

Let $a \in L^\infty((0, T) \times \Omega)^d$.

$$\exists \lambda_0 \gtrsim 1 + \|a\|_{L^\infty((0, T) \times \Omega)^d}, \forall \lambda \geq \lambda_0,$$

$$\exists C, s_0 > 0, \forall T > 0, \forall s \geq s_0(T + T^2)$$

$$\begin{aligned} I_1(\tilde{w}) + I_2(\tilde{\pi}) \leq C & \left(\iint_{(0, T) \times \Omega} s \tilde{\zeta} e^{2s\tilde{\varphi}} |\tilde{f}|^2 dy dt \right. \\ & \left. + \iint_{(0, T) \times \omega_1} e^{2s\tilde{\varphi}} \left[(s\tilde{\zeta})^3 |\tilde{w}|^2 + (s\tilde{\zeta})^2 |\tilde{\pi}|^2 \right] dy dt \right). \end{aligned}$$

From our result to the observability

- To get rid of the pressure term in the observation, we need to impose that

$$\partial_t a \in L^2(0, T; L^\sigma(\Omega))^d \quad \text{with} \quad \begin{cases} \sigma > 1 & \text{if } d = 2, \\ \sigma > 6/5 & \text{if } d = 3. \end{cases}$$

—→ This follows completely from FGIP2004.

- To obtain the observability of the Oseen system, we cut the integrals in $I_1(\tilde{w})$ and use energy estimates for the Oseen system.

First step: global Carleman estimates

The Oseen system as a heat equation

We write

$$\begin{cases} \partial_t \tilde{w} - \Delta \tilde{w} = -\nabla \tilde{\pi} + \tilde{f} + (\mathbb{D}\tilde{w}) a & \text{in } (0, T) \times \Omega, \\ \tilde{w} = 0 & \text{on } (0, T) \times \partial\Omega. \end{cases}$$

Carleman estimate for the heat equation

$$\begin{aligned} I_1(\tilde{w}) \lesssim & \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} |\nabla \tilde{\pi}|^2 \, dy \, dt + \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} |\tilde{f}|^2 \, dy \, dt \\ & + \iint_{(0,T) \times \Omega} e^{2s\tilde{\varphi}} |(\mathbb{D}\tilde{w}) a|^2 \, dy \, dt \\ & + \iint_{(0,T) \times \omega_0} s^3 \lambda^4 \tilde{\zeta}^3 e^{2s\tilde{\varphi}} |\tilde{w}|^2 \, dy \, dt. \end{aligned}$$

First step: global Carleman estimates

Equation for the pressure

Using that $\operatorname{div} \tilde{w} = 0$, we have that $\tilde{\pi}$ satisfies

$$\Delta \tilde{\pi} = \operatorname{div} \left(\tilde{f} + (\mathbb{D} \tilde{w}) a \right) \quad \text{in } (0, T) \times \Omega.$$

Difficulty: there is no boundary condition!

We apply a Carleman estimate obtained in

O. Y. Imanuvilov and J.-P. Puel. “Global Carleman estimates for weak solutions of elliptic nonhomogeneous Dirichlet problems”. *Int. Math. Res. Not.* (2003).

First step: global Carleman estimates

$$\Delta \tilde{\pi} = \operatorname{div} \left(\tilde{f} + (\mathbb{D}\tilde{w}) a \right) \quad \text{in } (0, T) \times \Omega.$$

We set

$$\varphi_0 := \tilde{\varphi}|_{\partial\Omega}, \quad \zeta_0 := \tilde{\zeta}|_{\partial\Omega}.$$

Carleman estimate for the Laplace equation:

$$\begin{aligned} I_2(\tilde{\pi}) &\lesssim \iint_{(0,T) \times \Omega} s \tilde{\zeta} e^{2s\tilde{\varphi}} |\tilde{f}|^2 \, dy \, dt + \iint_{(0,T) \times \Omega} s \tilde{\zeta} e^{2s\tilde{\varphi}} |(\mathbb{D}\tilde{w}) a|^2 \, dy \, dt \\ &\quad + \iint_{(0,T) \times \omega_1} s^2 \lambda^2 \tilde{\zeta}^2 e^{2s\tilde{\varphi}} |\tilde{\pi}|^2 \, dy \, dt \\ &\quad + \int_0^T (s \zeta_0)^{1/2} e^{2s\varphi_0} |\tilde{\pi}|_{H^{1/2}(\partial\Omega)}^2 \, dt. \end{aligned}$$

First step: global Carleman estimates

Gathering the two Carleman inequalities (heat equation and Laplace equation), we find

Proposition

$$\begin{aligned} I_1(\tilde{w}) + I_2(\tilde{\pi}) &\lesssim \iint_{(0,T) \times \Omega} s \tilde{\zeta} e^{2s\tilde{\varphi}} |\tilde{f}|^2 \, dy \, dt \\ &\quad + \iint_{(0,T) \times \omega_1} e^{2s\tilde{\varphi}} \left[s^3 \lambda^4 \tilde{\zeta}^3 |\tilde{w}|^2 + s^2 \lambda^2 \tilde{\zeta}^2 |\tilde{\pi}|^2 \right] \, dy \, dt \\ &\quad + \int_0^T (s \zeta_0)^{1/2} e^{2s\varphi_0} |\tilde{\pi}|_{H^{1/2}(\partial\Omega)}^2 \, dt. \end{aligned}$$

→ we need to get rid of the boundary term with the pressure.

Truncation and change of variables

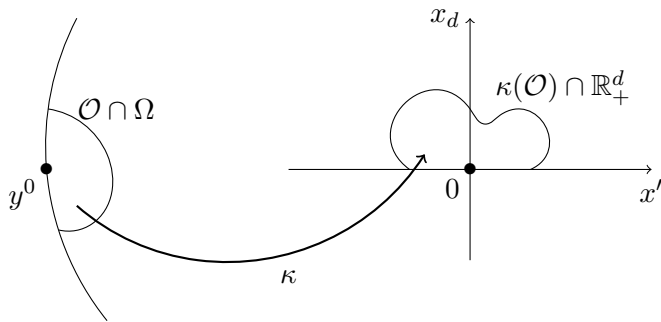


Figure: Local chart $(\mathcal{O}, \kappa)$ in a neighborhood of $y^0 \in \partial\Omega$.

We set

$$\mathbb{R}_+^d := \left\{ (x', x_d) \in \mathbb{R}^d : x_d > 0 \right\}.$$

Truncation and change of variables

Truncation in space

$$\theta \in C_0^\infty(\mathbb{R}^d), \quad \text{supp } \theta \subset \mathcal{O}.$$

We set

$$\check{w} := \theta \tilde{w}, \quad \check{\pi} := \theta \tilde{\pi}.$$

$$\Delta \tilde{\pi} = \text{div} \left(\tilde{f} + (\mathbb{D} \tilde{w}) a \right) \quad \text{in } (0, T) \times \Omega.$$

becomes

$$-\Delta \check{\pi} = \check{H}_0 + \text{div } \check{H},$$

where

$$\check{H}_0 := [\theta, \Delta] \tilde{\pi} + \nabla \theta \cdot \left(\tilde{f} + (\mathbb{D} \tilde{w}) a \right), \quad \check{H} := -\theta \left(\tilde{f} + (\mathbb{D} \tilde{w}) a \right).$$

Truncation and change of variables

Change of variables

$$\pi := \breve{\pi} \circ \kappa^{-1}$$

$$-\Delta \breve{\pi} = \breve{H}_0 + \operatorname{div} \breve{H} \quad \longrightarrow \quad \mathcal{M}\pi = H_0 + \operatorname{div} H \quad \text{in } (0, T) \times \mathbb{R}_+^d.$$

Lemma

One can choose $(\mathcal{O}, \kappa)$ such that

$$\mathcal{M}\pi := -\frac{\partial^2 \pi}{\partial x_d^2} - \sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \frac{\partial^2 \pi}{\partial x_j \partial x_k} - \mathfrak{h} \cdot \nabla \pi,$$

with

$$\sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \xi_j \xi_k \geq C |\xi'|^2 \quad \left(\xi' = (\xi_1, \dots, \xi_{d-1}) \in \mathbb{R}^{d-1} \right).$$

Objective

Estimate

$$\int_0^T (s\zeta_0)^{1/2} e^{2s\varphi_0} |\tilde{\pi}|_{H^{1/2}(\partial\Omega)}^2 dt.$$

We thus need to consider

$$e^{s\tilde{\varphi}}\tilde{\pi} = e^{\frac{s}{\ell}[e^{\lambda\tilde{\psi}} - e^{2\lambda}]} \tilde{\pi} = e^{\tau\tilde{\phi}}\tilde{\pi} e^{-\tau e^{2\lambda}},$$

with

$$\tau := \frac{s}{\ell}, \quad \tilde{\phi} := e^{\lambda\tilde{\psi}}.$$

Change of variables

Weights on Ω

$$\tilde{\psi} \in C^\infty(\mathbb{R}^d), \quad \tilde{\psi} = 0 \text{ on } \partial\Omega,$$

$$\nabla \tilde{\psi}(y) = 0 \implies y \in \omega_0.$$

$$s, \quad \tilde{\varphi} := \frac{1}{\ell} \left[e^{\lambda \tilde{\psi}} - e^{2\lambda} \right].$$
$$e^{s\tilde{\varphi}} \tilde{\pi}$$

Weights on $\mathbb{R}_+^d$

$$\psi := \tilde{\psi} \circ \kappa^{-1},$$

$$\inf_{\kappa(\mathcal{O})} \partial_{x_d} \psi > 0.$$

$$\tau := \frac{s}{\ell}, \quad \phi := e^{\lambda \psi}.$$

$$q := e^{\tau \phi} \pi$$

Conjugaison

Before conjugaison

$$\mathcal{M}\pi = H_0 + \operatorname{div} H \quad \text{in } (0, T) \times \mathbb{R}_+^d,$$

with

$$\mathcal{M}\pi := -\frac{\partial^2 \pi}{\partial x_d^2} - \sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \frac{\partial^2 \pi}{\partial x_j \partial x_k} - \mathfrak{h} \cdot \nabla \pi,$$

After conjugaison

$$q := e^{\tau \phi} \pi$$

$$\mathcal{M}_\phi q = G_0 + \operatorname{div}(G) \quad \text{in } (0, T) \times \mathbb{R}_+^d.$$

$$\mathcal{M}_\phi q = G_0 + \operatorname{div}(G) \quad \text{in } (0, T) \times \mathbb{R}_+^d.$$

where

$$\begin{aligned} \mathcal{M}_\phi q = & -\frac{\partial^2 q}{\partial x_d^2} - \sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \frac{\partial^2 q}{\partial x_j \partial x_k} \\ & -\tau^2 \left[\left(\frac{\partial \phi}{\partial x_d} \right)^2 + \sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \frac{\partial \phi}{\partial x_j} \frac{\partial \phi}{\partial x_k} \right] q \\ & + 2\tau \left[\frac{\partial \phi}{\partial x_d} \frac{\partial q}{\partial x_d} + \sum_{j,k=1}^{d-1} \mathfrak{g}^{j,k} \frac{\partial \phi}{\partial x_j} \frac{\partial q}{\partial x_k} \right] + \dots \end{aligned}$$

Ideas of the proof

Simplifications

Assume

$$g^{j,k} = \delta_{j,k}, \quad \phi = \phi(x_d), \quad \frac{\partial \phi}{\partial x_d} = 1 > 0,$$

so that

$$\mathcal{M}_{\phi} q = -\frac{\partial^2 q}{\partial x_d^2} - \sum_{j=1}^{d-1} \frac{\partial^2 q}{\partial x_j^2} - \tau^2 q + 2\tau \frac{\partial q}{\partial x_d} + \dots$$

Important property of ψ

Recall that

$$\phi := e^{\lambda \psi}, \quad \inf_{\kappa(\mathcal{O})} \partial_{x_d} \psi > 0.$$

Thus

$$\inf_{\kappa(\mathcal{O})} \frac{\partial \phi}{\partial x_d} > 0.$$

Ideas of the proof

Fourier transform

$$q(x', x_d) \longrightarrow \widehat{q}(\xi', x_d).$$

$$\begin{aligned}\widehat{\mathcal{M}_\phi q} &= -\frac{\partial^2 \widehat{q}}{\partial x_d^2} + |\xi'|^2 \widehat{q} - \tau^2 \widehat{q} + 2\tau \frac{\partial \widehat{q}}{\partial x_d} \\ &= -\left(\frac{\partial}{\partial x_d} - (\tau + |\xi'|)\right)\left(\frac{\partial}{\partial x_d} - (\tau - |\xi'|)\right)\widehat{q}.\end{aligned}$$

We need to obtain estimates for the equations

$$\frac{\partial}{\partial x_d} \hat{u} - (\tau + |\xi'|) \hat{u} = \hat{f}$$

and

$$\frac{\partial}{\partial x_d} \hat{v} - (\tau - |\xi'|) \hat{v} = \hat{g}$$

with $\hat{u}(\xi', x_d) = 0$ and $\hat{v}(\xi', x_d) = 0$ for $x_d \geq R(\mathcal{O})$.

Ideas of the proof

First equation

$$\frac{\partial}{\partial x_d} \hat{u} - (\tau + |\xi'|) \hat{u} = \hat{f}$$

We multiply by $(\tau + |\xi'|) \hat{u}$:

$$\begin{aligned} & - \int_{\mathbb{R}^{d-1}} (\tau + |\xi'|) \int_0^\infty \frac{1}{2} \frac{\partial}{\partial x_d} |\hat{u}|^2 dx_d d\xi' \\ & + \frac{1}{2} \int_{\mathbb{R}_+^d} (\tau + |\xi'|)^2 |\hat{u}|^2 d\xi' dx_d \leq \frac{1}{2} \int_{\mathbb{R}_+^d} |\hat{f}|^2 d\xi' dx_d \end{aligned}$$

Ideas of the proof

First equation

$$\frac{\partial}{\partial x_d} \widehat{u} - (\tau + |\xi'|) \widehat{u} = \widehat{f}$$

$$\begin{aligned} & \int_{\mathbb{R}^{d-1}} (\tau + |\xi'|) |\widehat{u}(\xi', 0)|^2 d\xi' \\ & + \int_{\mathbb{R}^{d-1} \times \mathbb{R}} \tau^2 |\widehat{u}|^2 + |i\xi' \widehat{u}|^2 + \left| \frac{\partial}{\partial x_d} \widehat{u} \right|^2 d\xi' dx_d \\ & \lesssim \int_{\mathbb{R}^{d-1} \times \mathbb{R}} |\widehat{f}|^2 d\xi' dx_d \end{aligned}$$

Remark

We don't need any boundary condition !

Ideas of the proof

In order to the same for

$$\frac{\partial}{\partial x_d} \hat{v} - \underbrace{(\tau - |\xi'|)}_{???} \hat{v} = \hat{g}$$

we need to assume

$$|\xi'| < C\tau \longrightarrow \text{LOW FREQUENCY}$$

Pseudo-differential operators



To justify the above ideas in the case **without simplifications**, we need to use **pseudo-differential operators**

Decomposition

$$q \approx \text{Op}_T(\chi^-) q + \text{Op}_T(\chi^+) q.$$

Splitting Low and High frequencies

Decomposition

$$q \approx \text{Op}_T(\chi^-) q + \text{Op}_T(\chi^+) q.$$

Lemma

There exist $\chi^-, \chi^+ \in \mathbf{S}_{T,\tau}^0$ such that

$$(x, \xi', \tau) \in \text{supp}(\chi^+) \implies |\xi'| \geq C\tau \quad \textit{High Frequency}$$

$$(x, \xi', \tau) \in \text{supp}(\chi^-) \implies |\xi'| \leq C\tau \quad \textit{Low Frequency}$$

Semiclassical norms

Definition

Semiclassical norms: in $\mathbb{R}_+^d$

$$\|u\|_{n,\tau} := \left(\sum_{|\beta|+j \leq n} \tau^{2j} \left\| \partial_x^\beta u \right\|_{L^2(\mathbb{R}_+^d)}^2 \right)^{1/2}.$$

At the boundary,

$$|u|_{e,\tau} := \left| \text{Op}_\tau (|\xi'|^2 + \tau^2)^{e/2} u \right|_{L^2(\mathbb{R}^{d-1})}.$$

Estimates of the low frequency part

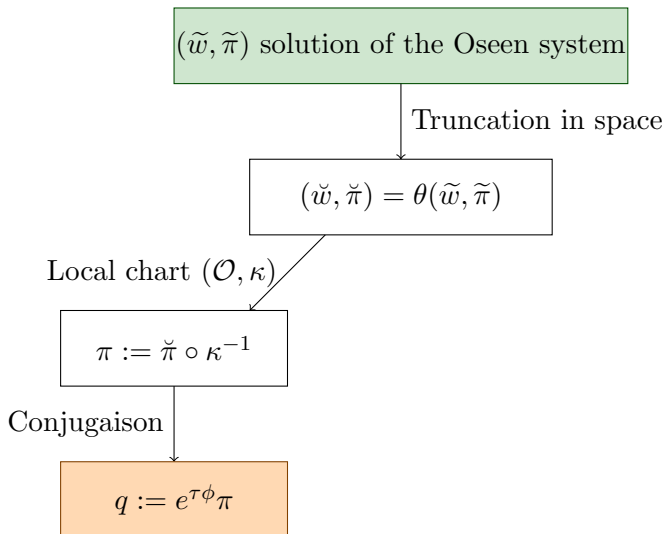
We assume

$$\mathcal{M}_\phi q = G_0 + \operatorname{div}(G) \quad \text{in } (0, T) \times \mathbb{R}_+^d.$$

Theorem

There exist $\tau^ > 0$ and $C > 0$ such that for any $\tau \geq \tau^*$,*

$$\begin{aligned} \tau^2 \|\operatorname{Op}_T(\chi^-) q\|_{1,\tau}^2 + \tau^2 \|\operatorname{Op}_T(\chi^-) q\|_{1/2,\tau}^2 \\ \leq C \left(\|G_0\|_{0,\tau}^2 + \tau^2 \|G\|_{0,\tau}^2 + \|q\|_{1,\tau}^2 \right). \end{aligned}$$



Low frequency arguments

Estimates of the low frequency part

$$\Delta \tilde{\pi} = \operatorname{div} \left(\tilde{f} + (\mathbb{D}\tilde{w}) a \right) \quad \text{in } (0, T) \times \Omega.$$

Corollary

There exists $s_0 > 0$ such that for any $s \geq s_0 T^2$,

$$\begin{aligned} & \left| (s\zeta_0)^{1/4} e^{s\varphi_0} \mathbf{Op}_T(\chi^-) ([\theta \tilde{\pi}] \circ \kappa^{-1}) \right|_{L^2(0, T; H^{1/2}(\mathbb{R}^{d-1}))}^2 \\ & \lesssim \left\| \left(s\tilde{\zeta} \right)^{1/4} e^{s\tilde{\varphi}} \tilde{f} \right\|_{L^2(0, T; L^2(\Omega))^d}^2 + \left\| \left(s\tilde{\zeta} \right)^{1/4} e^{s\tilde{\varphi}} \nabla \tilde{w} \right\|_{L^2(0, T; L^2(\Omega))^{d \times d}}^2 \\ & \quad + \left\| \left(s\tilde{\zeta} \right)^{1/4} e^{s\tilde{\varphi}} \tilde{\pi} \right\|_{L^2(0, T; L^2(\Omega))}^2 + \left\| \left(s\tilde{\zeta} \right)^{-3/4} e^{s\tilde{\varphi}} \nabla \tilde{\pi} \right\|_{L^2(0, T; L^2(\Omega))^d}^2. \end{aligned}$$

High frequency estimates

Objective

Estimate

$$\int_0^T (s\zeta_0)^{1/2} e^{2s\varphi_0} |\widetilde{\pi}|_{H^{1/2}(\partial\Omega)}^2 dt.$$

Truncation in space

$$\theta \in C_0^\infty(\mathbb{R}^d), \quad \text{supp } \theta \subset \mathcal{O}.$$

We set

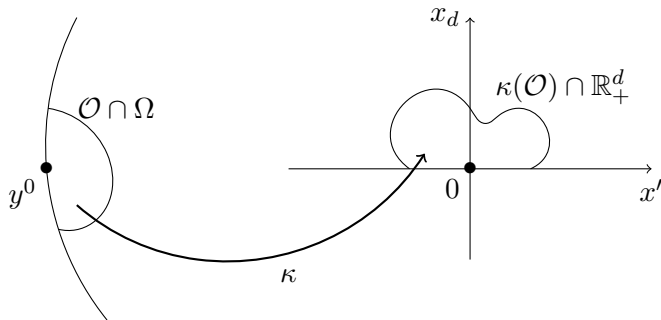
$$\check{w} := \theta \widetilde{w}, \quad \check{\pi} := \theta \widetilde{\pi}.$$

Weak conjugaison

We set

$$w^{\natural} := (s\zeta_0)^{1/4} e^{s\varphi_0} \check{w}, \quad \pi^{\natural} := (s\zeta_0)^{1/4} e^{s\varphi_0} \check{\pi}.$$

High frequency estimates



Definition

For $f : \Omega \cap \mathcal{O} \rightarrow \mathbb{R}$, we set

$$\mathbb{O}^+ f := [\text{Op}_\top (\chi^+) (f \circ \kappa^{-1})] \circ \kappa.$$

High frequency estimates

With this operator, we define

$$w^\sharp := \mathbb{O}^+ w^\natural, \quad \pi^\sharp := \mathbb{O}^+ \pi^\natural$$

so that

$$\begin{cases} \partial_t w^\sharp - \Delta w^\sharp + \nabla \pi^\sharp = f^\sharp & \text{in } (0, T) \times \Omega, \\ \operatorname{div} w^\sharp = g^\sharp & \text{in } (0, T) \times \Omega, \\ w^\sharp = 0 & \text{on } (0, T) \times \partial\Omega, \\ w^\sharp(0, \cdot) = 0 & \text{in } \Omega. \end{cases}$$

High frequency estimates

Lemma

We have the estimate

$$\begin{aligned} & \left\| w^\sharp \right\|_{L^2(0,T;H^2(\Omega))^d}^2 + \left\| \partial_t w^\sharp \right\|_{L^2(0,T;L^2(\Omega))^d}^2 + \left\| \pi^\sharp \right\|_{L^2(0,T;H^1(\Omega))}^2 \\ & \lesssim \left\| f^\sharp \right\|_{L^2(0,T;L^2(\Omega))^d}^2 + \left\| g^\sharp \right\|_{L^2(0,T;H^1(\Omega))}^2 \\ & \quad + \left\| \tau^{-1} \partial_t g^\sharp \right\|_{L^2(0,T;L^2(\Omega))} \left\| \tau \pi^\sharp \right\|_{L^2(0,T;L^2(\Omega))} . \end{aligned}$$

Question: semiclassical norms?

High frequency estimates

Property of χ^+

$$(x, \xi', \tau) \in \text{supp}(\chi^+) \implies |\xi'| \geq C\tau \quad \text{High Frequency}$$

$$\|\text{Op}_T(\chi^+) f\|_{H^1(\mathbb{R}_+^d)}^2 \geq \|\text{Op}_T(\chi^+) f\|_{L^2(\mathbb{R}_+; H^1(\mathbb{R}^{d-1}))}^2.$$

Consequently,

$$\begin{aligned} \|\text{Op}_T(\chi^+) f\|_{H^1(\mathbb{R}_+^d)}^2 \\ \gtrsim \left(\text{Op}_T \left(1 + |\xi'|^2 \right) \text{Op}_T(\chi^+) f, \text{Op}_T(\chi^+) f \right)_{L^2(\mathbb{R}_+^d)}. \end{aligned}$$

High frequency estimates

Property of χ^+

$$(x, \xi', \tau) \in \text{supp}(\chi^+) \implies 1 + |\xi'|^2 \geq C(\tau^2 + |\xi'|^2) \quad \text{High Frequency}$$

We have obtained

$$\begin{aligned} \|\text{Op}_\tau(\chi^+) f\|_{H^1(\mathbb{R}_+^d)}^2 \\ \gtrsim \left(\text{Op}_\tau(1 + |\xi'|^2) \text{Op}_\tau(\chi^+) f, \text{Op}_\tau(\chi^+) f \right)_{L^2(\mathbb{R}_+^d)}. \end{aligned}$$

Gårding inequality:

$$\|\text{Op}_\tau(\chi^+) f\|_{1,\tau}^2 \lesssim \|\text{Op}_\tau(\chi^+) f\|_{H^1(\mathbb{R}_+^d)}^2 + \|f\|_{L^2(\mathbb{R}_+^d)}^2.$$

High frequency estimates

We thus deduce

Lemma

There exist $C > 0$ and $s_0 > 0$ such that for any $T > 0$, for any $s \geq s_0(T + T^2)$

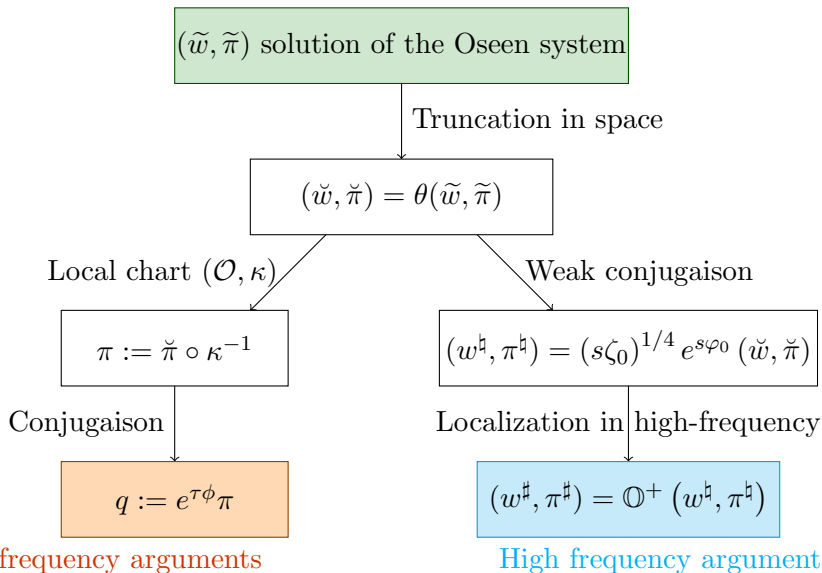
$$\begin{aligned} & \left\| \text{Op}_T(\chi^+) \left(w^\sharp \circ \kappa^{-1} \right) \right\|_{L^2(0,T;H_T^2)}^2 + \left\| \text{Op}_T(\chi^+) \left(\pi^\sharp \circ \kappa^{-1} \right) \right\|_{L^2(0,T;H_T^1)}^2 \\ & \leq C \left(\left\| f^\sharp \right\|_{L^2(0,T;L^2(\Omega))^d}^2 + \left\| g^\sharp \right\|_{L^2(0,T;H^1(\Omega))}^2 + \left\| \tau^{-1} \partial_t g^\sharp \right\|_{L^2(0,T;L^2(\Omega))}^2 \right. \\ & \quad \left. + \left\| w^\sharp \right\|_{L^2(0,T;L^2(\Omega))^d}^2 + \left\| \pi^\sharp \right\|_{L^2(0,T;L^2(\Omega))}^2 \right). \end{aligned}$$

High frequency estimates

Semiclassical trace lemma

$$\begin{aligned} \left| (s\zeta_0)^{1/4} e^{s\varphi_0} \text{Op}_T(\chi^+) (\theta \tilde{\pi} \circ \kappa^{-1}) \right|_{L^2(0,T;H^{1/2}(\mathbb{R}^{d-1}))}^2 \\ \lesssim \left\| \text{Op}_T(\chi^+) (\pi^\sharp \circ \kappa^{-1}) \right\|_{L^2(0,T;H_\tau^1)}^2. \end{aligned}$$

History



End of the proof

- We estimate all the “commutator” terms.
- We add the estimates in low and high frequency modes.
- We consider a finite family $(\mathcal{O}_j, \kappa_j)$, $j = 1, \dots, J$ such that

$$\partial\Omega \subset \bigcup_{j=1}^J \mathcal{O}_j.$$

- We consider $\{\theta_j\}_{j=1}^J$ a partition of unity associated with this covering:

$$\theta_j \in C^\infty(\mathbb{R}^d), \quad \theta_j \geq 0, \quad \sum_{j=1}^J \theta_j \equiv 1 \text{ on } \partial\Omega, \quad \text{supp } \theta_j \subset \mathcal{O}_j.$$

Our main result: statement

Theorem

Let $a \in L^\infty((0, T) \times \Omega)^d$.

$$\exists \lambda_0 \gtrsim 1 + \|a\|_{L^\infty((0, T) \times \Omega)^d}, \forall \lambda \geq \lambda_0,$$

$$\exists C, s_0 > 0, \forall T > 0, \forall s \geq s_0(T + T^2)$$

$$\begin{aligned} I_1(\tilde{w}) + I_2(\tilde{\pi}) \leq & C \left(\iint_{(0, T) \times \Omega} s \tilde{\zeta} e^{2s\tilde{\varphi}} |\tilde{f}|^2 dy dt \right. \\ & \left. + \iint_{(0, T) \times \omega_1} e^{2s\tilde{\varphi}} \left[(s \tilde{\zeta})^3 |\tilde{w}|^2 + (s \tilde{\zeta})^2 |\tilde{\pi}|^2 \right] dy dt \right). \end{aligned}$$