Asymptotic expansion of the Laplace Functional for the Sine Process

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Let Conf(\mathbb{R}) be the space of all the configurations of points $\{\xi_j\}_1^n$ in the real line. We have a "*n-point process*" when we introduce a probability measure $d\mathbb{P}_n$ on $\text{Conf}(\mathbb{R})$. A n-point process is a Determinantal Point Process (DPP) if the probability measure $d\mathbb{P}_n$ has the form

$$
\mathrm{d}\mathbb{P}_n(\xi_1,\ldots,\xi_n)=\frac{1}{n!}\mathrm{det}[K(\xi_j,\xi_k)]_{j,k=1}^n\mathrm{d}\xi_1\otimes\ldots\mathrm{d}\xi_n
$$

where $K(x, y)$ is called "correlation kernel".

A important example of DPP is an ∞ -point process, the Sine Process, which has a correlation Kernel of the form:

$$
K^{\sin}(x, y) = \frac{\sin \left[\pi(x - y)\right]}{\pi(x - y)} \quad x, y \in \mathbb{R}.
$$
 (1)

This process is well studied and has a lot of connection with other fields

- 1 Random Hermitian matrices [\[Erdös\]](#page-0-0);
- 2 Orthogonal Polynomials;
- 3 Integrable systems [\[Tracy-Widown,](#page-29-0) [Claeys-Tarricone\]](#page-29-1);

4 ...

Let $\mathcal{B}_+(\mathbb{R})$ be the space of bounded non-negative measurable functions $f : \mathbb{R} \to [0, +\infty)$. In general theory of point process, the "Laplace functional" is defined as the map

$$
\mathscr{L}_{\lambda}: \mathcal{B}_{+}(\mathbb{R}) \to \mathbb{R}^{+}, \quad f \mapsto \mathscr{L}_{\lambda}[f], \quad \mathscr{L}_{\lambda}[f] := \mathbb{E}\left[e^{-\lambda \int_{\mathbb{R}} f \,d\xi}\right], \tag{2}
$$

where $\int_{\mathbb{R}} f \mathrm{d}\xi = \sum_{x \in \{\xi_j\}_1^n} f(x)$ is a random variable and the average is over the probability measure $d\mathbb{P}_n$.

For DPP, the Laplace Functional is a Fredholm determinant

$$
\mathscr{L}_{\lambda}[f] = \det\left[\mathrm{id} - \sqrt{\sigma_{\lambda}} \mathscr{K}\sqrt{\sigma_{\lambda}}\right],\tag{3}
$$

where $\sigma_\lambda(x):=1-{\rm e}^{-\lambda f(x)}$ and the operator $\mathscr{K}:L^2(\mathbb{R})\to L^2(\mathbb{R})$ has kernel $K(x,y)$ the correlation kernel of the DPP itself.

- 1 The Laplace functional characterize the point process, i.e. each point process has it's own unique \mathscr{L}_λ ;
- 2 We can use to prove important results about the distribution of points and the linear statistics of the process, e.g the Chernoff bound

$$
\mathbb{P}\left[\int_{\mathbb{R}} f(x) \leq A\right] \leq \mathscr{L}_{\lambda}[f] e^{\lambda A}.
$$

In our work, we study the asymptotic expansion of the Laplace functional $\mathscr{L}_{\lambda}[f]$ for the Sine Point Process.

We will consider $f \in \mathcal{B}_+(\mathbb{R})$ which satisfies the following conditions:

H.1 f can be extended analytically in a sector $S \subset \mathbb{C}$ containing \mathbb{R} ;

f

H.2 $f(-x) = f(x)$, with a single maximum at $x = 0$;

H.3 $f(x)$ and $f'(x)$ decays exponentially at $x \sim \pm \infty$:

$$
f(x) \sim c_1 e^{-c_2|x|}, \text{ as } x \sim \pm \infty,
$$

(4)

$$
f(x) \sim \mp c_1 c_2 e^{-c_2|x|}, \text{ as } x \sim \pm \infty,
$$

(5)

with $c_1, c_2 > 0$.

From the definition of the Lapalce functional, the configurations that contribute more on \mathscr{L}_λ are the one with few or no points near the maximum of f .

Theorem

Given $f(z) \in \mathcal{B}_{+}(\mathbb{R})$ which satisfies the hypothesis [H.1](#page-5-0), [H.2](#page-5-1) and [H.3](#page-5-2), then the Laplace functional $\mathscr{L}_{\lambda}[f]$ for the sine process has the following expansion as $\lambda \to \infty$:

$$
\log \mathcal{L}_{\lambda}[f] = -\frac{2\pi^2}{c_2^2} (\log \lambda)^2 + \mathcal{O}((\log \lambda)^{\frac{3}{2}}). \tag{6}
$$

Let $\mathscr{K}^{\sf sin}$ be the integral operator with kernel $\mathsf{K}^{\sf sin}(\mathsf{x},\mathsf{y})$, this operator is an *"integrable* operator", i.e the Kernel has the form

$$
K^{\sin}(x, y) = \frac{\vec{h}(x)^T \vec{t}(y)}{x - y}, \quad \vec{h}(x)^T \vec{t}(x) \equiv 0 \tag{7}
$$

where $\vec{h}(x)$, $\vec{t}(y)$ are vectors defined as

$$
\vec{h}(x) := \frac{1}{2\pi i} \begin{pmatrix} e^{i\pi x} \\ e^{-i\pi x} \end{pmatrix} \quad \vec{t}(y) := \begin{pmatrix} e^{-i\pi y} \\ -e^{i\pi y} \end{pmatrix}.
$$
 (8)

We can use the results of Its, Izergin, Korepin and Slavnov [\[IIKS\]](#page-29-2) to study the Fredholm determinant of $\mathscr{K}^{\mathsf{sin}}$.

If $K: L^2(\Sigma, |dz|) \to L^2(\Sigma, |dz|)$ is an integrable operators, the resolvent operator $\mathcal{R}:=-{\rm Id\,}+(\mathop{\rm Id\,}-\mathcal{K})^{-1}$ is still an integrable operator with kernel R given by:

$$
R(z, w) = \frac{p(z)^T \Gamma(z)^T (\Gamma^{-1}(w))^T q(w)}{z - w}, \qquad (9)
$$

where $\Gamma(z) \in Mat(r \times r, \mathbb{C})$ is a sectional analytic matrix function which satisfies the Riemann-Hilbert problem:

$$
\Gamma_+(z) = \Gamma_-(z)M(z) \text{ for } z \in \Sigma, \qquad M(z) = 1 + 2\pi i p(z) q(z)^T,
$$

$$
\Gamma(z) \sim 1 + \mathcal{O}(z^{-1}), \text{ as } z \to \infty.
$$
 (10)

If the solution of [\(10\)](#page-9-0) exists then

$$
\det(\mathrm{Id} - \mathcal{K}) \neq 0.
$$

From the Jacobi Identity we get that

$$
\partial_{\lambda} \log \mathscr{L}_{\lambda}[f] = \partial_{\lambda} \log \det \left[id - \sqrt{\sigma_{\lambda}} \mathscr{K}^{\sin} \sqrt{\sigma_{\lambda}} \right]
$$

= $-\text{Tr} \left[(\text{id} - \sqrt{\sigma_{\lambda}} \mathscr{K}^{\sin} \sqrt{\sigma_{\lambda}})^{-1} \circ \partial_{\lambda} (\sqrt{\sigma_{\lambda}} \mathscr{K}^{\sin} \sqrt{\sigma_{\lambda}}) \right]$
= $-\text{Tr} [\partial_{\lambda} \sigma_{\lambda} \mathscr{K}^{\sin} (\text{Id} - \sigma_{\lambda} \mathscr{K}^{\sin})^{-1}]$
= $-\int_{\mathbb{R}} f(x) L_{\lambda}(x, x) dx$

where $L_{\lambda}(x, y)$ is a smooth kernel of the operator

$$
\mathcal{L}_{\lambda} = \sqrt{1 - \sigma_{\lambda}} \mathscr{K}^{\sin}(\mathrm{Id} - \sigma_{\lambda} \mathscr{K}^{\sin})^{-1} \sqrt{1 - \sigma_{\lambda}} = \sqrt{\sigma_{\lambda}^{-1} - 1} \mathscr{K}^{\sin} \sqrt{\sigma_{\lambda}^{-1} - 1}
$$
(11)

where $\mathscr{R}^{\mathsf{sin}}$ is the resolvent of the operator $\mathscr{K}^{\mathsf{sin}}$.

We apply the result [\(9\)](#page-9-1) to the operator \mathcal{R}^{sin} and we get that \mathcal{L}_{λ} has Kernel

$$
L_{\lambda}(x,y) = \sqrt{1 - \sigma_{\lambda}(x)} \frac{\vec{h}(x)^{\top} Y^{\top}(x) (Y^{-1}(y))^{\top} \vec{t}(y)}{x - y} \sqrt{1 - \sigma_{\lambda}(y)};
$$
 (12)

Where $Y(z)$ is a 2 × 2 matrix, analytic in $z \in \mathbb{C} \setminus \mathbb{R}$, which satisfy the Riemann-Hilbert problem $Y_{+}(x) = Y_{-}(x)J_{Y}(x)$ as $x \in \mathbb{R}$, where the jump matrix has the form

$$
J_Y(x) := \begin{bmatrix} 1 - \sigma_\lambda(x) & \sigma_\lambda(x) e^{2i\pi x} \\ -\sigma_\lambda(x) e^{-2i\pi x} & 1 + \sigma_\lambda(x) \end{bmatrix}
$$
(13)

 $Y(z) \sim 1 + \mathcal{O}(z^{-1}), \text{ as } z \to \infty.$

We apply the following transformation

$$
\Psi(z) = Y(z) \begin{cases} \begin{bmatrix} e^{2i\pi z} & e^{2i\pi z} \\ e^{-2i\pi z} & 0 \end{bmatrix}; & \text{Im}[z] > 0; \\ \begin{bmatrix} e^{2i\pi z} & 0 \\ e^{-2i\pi z} & -e^{-2i\pi z} \end{bmatrix}; & \text{Im}[z] < 0; \end{cases}
$$
(14)

Then the Riemann-Hilbert problem becomes

$$
\Psi_{+}(x) = \Psi_{-}(x) \begin{bmatrix} 1 & e^{-\lambda f(x)} \\ 0 & 1 \end{bmatrix} \text{ as } x \in \mathbb{R}
$$
\n
$$
\Psi(z) \sim e^{iz\pi\sigma_3} \begin{cases} \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}; & \text{Im}[z] > 0; \\ \begin{bmatrix} 1 & 0 \\ 1 & -1 \end{bmatrix}; & \text{Im}[z] < 0; \end{cases}
$$
\n
$$
(16)
$$

Then we have that the diagonal term of the Kernel $L_{\lambda}(x, y)$ depends on the matrix $\Psi(x)$ as follows:

$$
L_{\lambda}(x,x) = \frac{1 - \sigma_{\lambda}(x)}{2\pi i} \left[\Psi(x)^{-1} \partial_x \Psi(x) \right]_{21}
$$
 (17)

and the Fredholm determinant becomes

$$
\partial_{\lambda} \log \mathscr{L}_{\lambda}[f] = \int_{\mathbb{R}} \partial_{\lambda} (e^{-\lambda f(x)}) \left[\Psi(x)^{-1} \partial_{x} \Psi(x) \right]_{21} \frac{dx}{2\pi i}
$$
(18)

This implies that, by knowing the large asymptotics of the matrix $\Psi(x)$ and $\partial_x \Psi(x)$ we can find the asymptotic of the Laplace functional $\mathscr{L}_{\lambda}[f]$.

We define the transformation:

$$
\hat{\Psi}(z) = \Psi(z) e^{-\lambda (g(z) + \frac{1}{2})\sigma_3} \tag{19}
$$

where *l* is a constant and $g(z)$ is a complex function, analytic in $z \in \mathbb{C} \setminus \mathbb{R}$ Then the RHP [\(15\)](#page-12-0) becomes

$$
\hat{\Psi}_{+}(x) = \hat{\Psi}_{-}(x) \begin{bmatrix} e^{\lambda(g_{-}(x) - g_{+}(x))} & e^{\lambda(g_{+}(x) + g_{-}(x) - f(x) + l)} \\ 0 & e^{-\lambda(g_{-}(x) - g_{+}(x))} \end{bmatrix} \text{ for } z \in \mathbb{R}
$$
\n
$$
\hat{\Psi}(z) = (1 + \mathcal{O}(z^{-1})) \begin{cases} \begin{bmatrix} e^{2i\pi z} & 1 \\ 1 & 0 \end{bmatrix}; & \text{Im}[z] > 0; \\ \begin{bmatrix} 1 & 0 \\ -e^{-2i\pi z} & -1 \end{bmatrix}; & \text{Im}[z] < 0; \end{cases}
$$
\n
$$
(21)
$$

the g-function

We define $B_\lambda \subset \mathbb{R}$ the set on which the approximation of $L_\lambda(x, x)$ is supported. According to the properties that the function $f(x)$ has, we suppose the

$$
B_{\lambda}=(-\infty;-a(\lambda))\cup(a(\lambda);+\infty).
$$

We are looking for a g-function $g(z)$ which solves the scalar Riemann-Hilbert problem:

$$
g_{+}(x) + g_{-}(x) - f(x) + l = 0 \text{ for } x \in B_{\lambda}
$$

\n
$$
g_{+}(x) - g_{-}(x) = 0 \text{ for } x \in \mathbb{R} \setminus B_{\lambda}
$$
\n(23)

with boundary condition

$$
g(z) = \pm \frac{i\pi z}{\lambda} - \frac{1}{2} + \frac{g_1}{z} + \mathcal{O}(z^{-2}) \text{ as } z \sim \infty \quad \pm \text{ Im}[z] > 0 \tag{24}
$$

$$
g(z) \sim \pm i(z \pm a(\lambda))^{3/2} \text{ for } z \sim \pm a(\lambda). \tag{25}
$$

By deriving the scalar RHP respect to x and solving the new RHP we get

$$
g'(z) = \frac{R(z)}{2\pi i} \int_{B_\lambda} \frac{f'(w)}{R_+(w)(w-z)} dw \qquad (26)
$$

where $R(z) := \sqrt{(z + a(\lambda))(a(\lambda) - z)}$ is a multi valued function, analytic in $z \in \mathbb{C} \setminus B_\lambda$. From the boundary condition at infinity we get a equation for the point $a(\lambda)$ we obtain an equation for the point $a(\lambda)$

$$
\frac{\pi^2}{\lambda} = \int_{a}^{+\infty} \frac{|f'(w)|}{|R(w)|} dw.
$$
 (27)

By analysing the equation [\(27\)](#page-16-0) as $\lambda \sim +\infty$, we get that $a(\lambda)$ has the following expansion:

$$
a(\lambda) = \frac{\log \lambda}{c_2} + C_f + \mathcal{O}(\log \log \lambda). \tag{28}
$$

The φ -function

We define $\varphi(z) := 2g(z) - f(z) + l$

Proposition

The function $\varphi(z)$ has the following expansions as $\lambda \sim +\infty$:

$$
\varphi(z) \sim \mp \frac{2\pi c_1}{\lambda} R(z) \text{ for } |z| >> a(\lambda)
$$
 (29)

$$
\varphi(z) \sim \mp 2\pi i c_1^2 \frac{[R(z)]^3}{\lambda} \text{ for } |z \pm a(\lambda)| << 1 \tag{30}
$$

where we have the $-$ sign when $\text{Re}(z) \in (a(\lambda); +\infty)$ and $+$ when $\text{Re}(z) \in (-\infty, a(\lambda))$

Idea of proof

Another way to represent $g'(z)$ is by integrating [\(26\)](#page-16-1) by parts

$$
g'(z) = \frac{R(z)}{2\pi i} \int_{B_{\lambda}} \frac{f'(w)}{R_+(w)(w-z)} dw
$$

=
$$
\frac{z}{2\pi i} \int_{a(\lambda)}^{+\infty} \frac{d}{dw} \left(\frac{f'(w)}{w}\right) \log \left[\frac{1 + \frac{\sqrt{w^2 - a^2}}{iR(z)}}{1 - \frac{\sqrt{w^2 - a^2}}{iR(z)}}\right] dw
$$
(31)

Since $R(z)=\pm i\sqrt{z^2-a^2}$ and $R_+(z)=-R_-(z)$ for $z\in B_\lambda$, then we have that

$$
g'_{+}(z) - g'_{-}(z) = \pm \frac{z}{\pi i} \int_{a(\lambda)}^{+\infty} \frac{d}{dw} \left(\frac{f'(w)}{w} \right) \log \left[\frac{1 + \sqrt{\frac{w^2 - a^2}{z^2 - a^2}}}{1 - \sqrt{\frac{w^2 - a^2}{z^2 - a^2}}} \right] dw \tag{32}
$$

where we have the sign + when $z \in (a(\lambda); +\infty)$ and the sign – when $z \in (-\infty; -a(\lambda)).$

From the definition of $\varphi(z)$

$$
\varphi_{+}'(z) = g_{+}'(z) - g_{-}'(z) = \pm \frac{z}{\pi i} \int_{a(\lambda)}^{+\infty} \frac{d}{dw} \left(\frac{f'(w)}{w} \right) \log \left[\frac{1 + \sqrt{\frac{w^2 - a^2}{z^2 - a^2}}}{1 - \sqrt{\frac{w^2 - a^2}{z^2 - a^2}}} \right] dw
$$

Then we define $\varphi'(z)$ as the analytic continuation of $\varphi_+'(z)$ analyze the asymptotic expansion as $\lambda \sim +\infty$. We have to consider two kinds of limits :

$$
\lambda \sim +\infty \text{ and } z >> a(\lambda) \text{ (or } z << -a(\lambda))
$$

$$
\varphi'(z) = \mp \frac{2\pi i c_1 z}{\lambda \sqrt{z^2 - a^2}} + \mathcal{O}((\lambda a(\lambda))^{-1})
$$
\n(33)

 $\lambda \sim +\infty$ and $|z \pm a(\lambda)| << 1$

$$
\varphi'(z) = \mp 2\pi i c_1^2 \frac{z\sqrt{z^2 - a^2(\lambda)}}{a(\lambda)\lambda} [1 + \mathcal{O}((\lambda a(\lambda))^{-1})] + \mathcal{O}((\lambda \sqrt{a(\lambda)})^{-1}) \tag{34}
$$

Opening of the Lenses

Proposition

As $\lambda \sim +\infty$ function $\varphi(z)$ satisfy the following inequalities:

 $\text{Re}(\lambda \varphi(z)) > 0$ for $z \in \mathcal{U}_r$; (35) $\text{Re}(\lambda \varphi(z)) < 0$ for $z \in (-a(\lambda) + r, a(\lambda) - r)$. (36) where $U_r := \{ z \in \mathbb{C} : |\operatorname{Im}(z)| = r, |\operatorname{Re}(z)| > r, r > 0 \}.$

This means that the Jump matrices along the lenses tends to the identity as $\lambda \sim +\infty$.

Model problem $P^{\infty}(z)$

We are looking for a matrix $P^{\infty}(z) \in GL(2,\mathbb{C})$, analytic in $\mathbb{C} \setminus B_{\lambda}$, that solves the RHP

The solution is the following:

$$
P^{\infty}(z) = U \gamma(z)^{\frac{\sigma_3}{4}} e^{-\frac{i\pi}{4}\sigma_3} U^{\dagger} \begin{cases} \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} & \text{for } \text{Im}(z) > 0; \\ 1 & \text{for } \text{Im}(z) < 0. \end{cases}
$$
(37)

where
$$
\gamma(z) := \left(\frac{a+z}{a-z}\right)^{\frac{1}{4}}
$$
 and $U := \frac{1}{\sqrt{2}}\begin{bmatrix} 1 & i \\ i & 1 \end{bmatrix}$.

Local Parametrix near $\pm a(\lambda)$

In the regions near the points $\pm a(\lambda)$, i.e. the disk $\mathbb{D}_r(\pm a) = \{z \in \mathbb{C} : |z \pm a(\lambda)| < r\}$, we define the conformal map

$$
\zeta(z) := \mathrm{e}^{i\frac{2}{3}\pi} \left(\frac{3}{4} \lambda \varphi(z) \right)^{\frac{2}{3}}.
$$

then the RHP near the points $\pm a(\lambda)$ is mapped to the following problem

From our previews results, we can estimate the 21 term of the matrix $\Psi(z)^{-1}\partial_z\Psi(z)$

$$
[\Psi(z)^{-1}\partial_z\Psi(z)]_{21} = -\lambda\varphi'_+(z)e^{+\lambda f(z)}
$$

+ Terms which depends on P^∞ and the local parametric

but this last terms, as $\lambda \sim +\infty$, are smaller respect the the leading order of the first term. Then, by substituting [\(38\)](#page-25-0) inside the formula [\(18\)](#page-13-0) and using the fact that the leading order of the Integral is given by the expansion $x \sim a(\lambda)$, we get that

$$
\partial_\lambda \log \mathscr{L}_\lambda[f] \sim \int_{a(\lambda)}^{+\infty} \lambda f(x) \varphi_+'(x) \frac{\mathrm{d}x}{\pi i} \sim -\int_{a}^{+\infty} 2\pi i c_1^2 \frac{x\sqrt{x^2-a^2(\lambda)}}{a(\lambda)} e^{-c_2 x} \frac{\mathrm{d}x}{\pi i} \sim -\frac{2\pi^2}{c_2^2} \frac{\log \lambda}{\lambda}
$$

By integrating this result we have [\(6\)](#page-7-0).

Large tail probability

From the Chernoff's bound, for $\lambda = 1/A$ and the result [\(6\)](#page-7-0), we have that

$$
\log \mathbb{P}\left[\int_\mathbb{R} f(x) \leq A\right] \leq \log \mathscr{L}_\lambda[f] + \lambda A \sim -\frac{\pi^2}{c_2^2}(1+o(1))\log^2 A
$$

as $A \rightarrow 0$. We would like to find a lower bound such that

$$
\log \mathbb{P}\left[\int_\mathbb{R} f(x) \leq A\right] \sim -\frac{\pi^2}{c_2^2} (1+o(1))\log^2 A
$$

as $A \rightarrow 0$.

Large tail probability; More general functions $f(x) \in \mathcal{B}_+(\mathbb{R})$; Connection with integrable systems.

Thank you

References

[Erdös] L. Erdös (2013)

Universality for random matrices and log-gases Current developments in mathematics 2012.

[Tracy-Widown] C. A. Tracy, H. Widom (1994) Fredholm determinants, differential equations and matrix models Comm. Math. Phys 163, no. 1, 33–72.

[Claeys-Tarricone] T.Claeys, S.Tarricone (2024) On the Integrable Structure of Deformed Sine Kernel Determinants Math Phys Anal Geom 27, 3

[IIKS] A. Its, A. Izergin, V. Korepin, and N. Slavnov (1990) Differential Equations for Quantum correlation functions International Journal of Modern Physics B 04, pp. 1003–1037