

# Using dileptons to estimate the initial temperature of QCD matter<sup>1,2</sup>

*Greg Jackson*

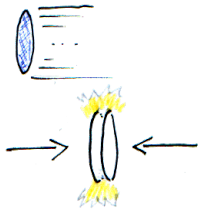
Subatech, CNRS/Nantes U./IMT-Atlantique

– AG du GDR QCD • Tours • May 2024 –

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<sup>1</sup> based on collaboration w/ J. Churchill, L. Du, C. Gale and S. Jeon

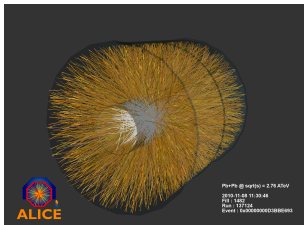
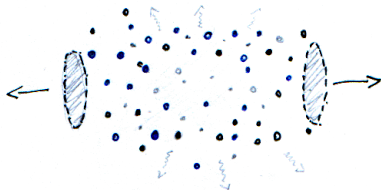
<sup>2</sup> supported by the ANR under grant No. 22-CE31-0018



$$\mathcal{L} = -\frac{1}{4}F^2 + \sum_f \bar{\psi} (i\not{D} - m_f) \psi$$

w/ non-Abelian fields

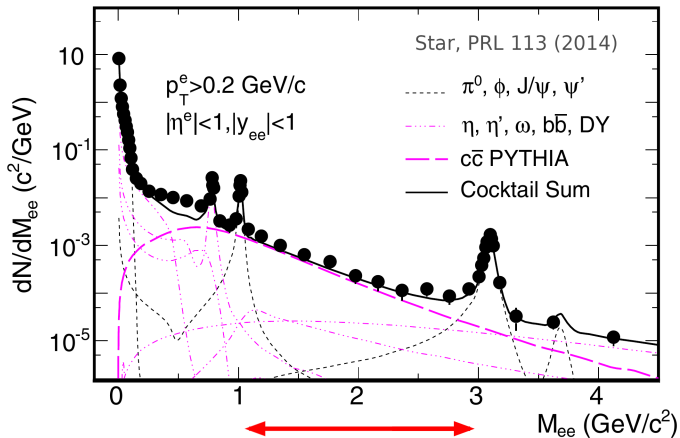
$$F_{\mu\nu} = \frac{1}{ig} [D_\mu, D_\nu]$$



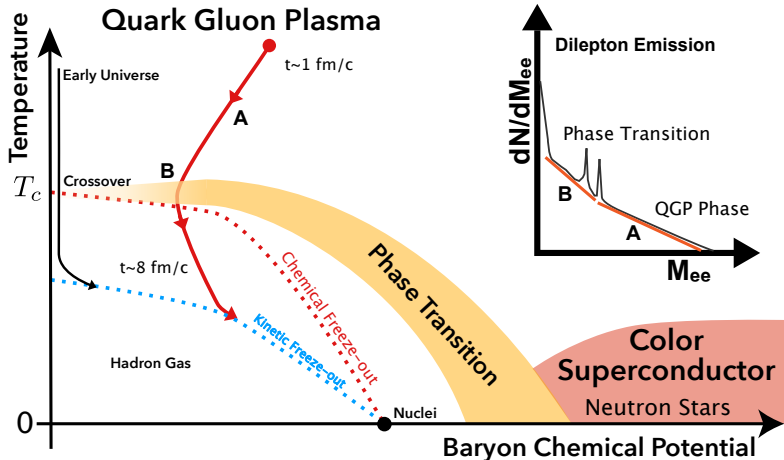
# Electromagnetic probes

Invariant mass spectrum of *dileptons pairs*, e.g. from  $q\bar{q} \rightarrow \gamma^* \rightarrow e^+e^-$

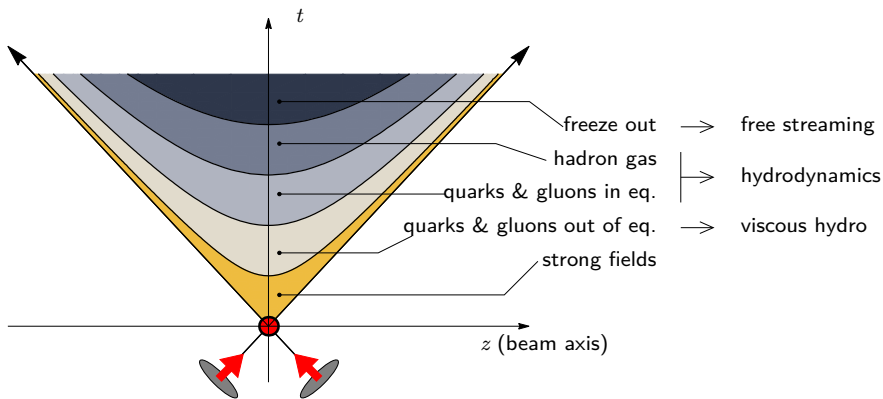
**Au + Au  $\sqrt{s_{NN}} = 200$  GeV (MinBias)**



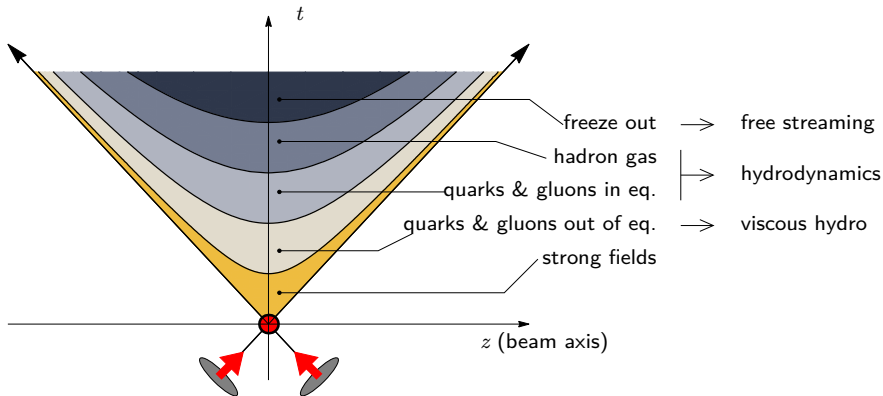
Intermediate Mass Range (IMR) = 1 ... 3 GeV



[STAR collaboration (2024)]



- $\tau \sim 0.2 \text{ fm}/c$  production of light quarks & gluons
- $1 - 2 \text{ fm}/c$  thermalisation rapid (?)
- $2 - 10 \text{ fm}/c$  **quark-gluon plasma**
- $10 - 20 \text{ fm}/c$  hadron gas
- $\tau \rightarrow \infty$  dilute, no further interactions



$\tau \sim 0.2 \text{ fm}/c$	production of light quarks & gluons	} beginning
$1 - 2 \text{ fm}/c$	thermalisation rapid (?)	
$2 - 10 \text{ fm}/c$	<b>quark-gluon plasma</b>	- middle
$10 - 20 \text{ fm}/c$	hadron gas	} end
$\tau \rightarrow \infty$	dilute, no further interactions	

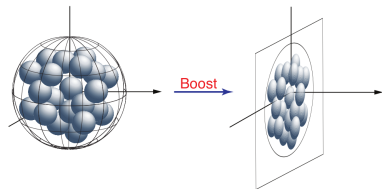
# Beginning: (initial conditions)

## Pedestrian approach:

sample nucleons with,  
Monte Carlo (Glauber)

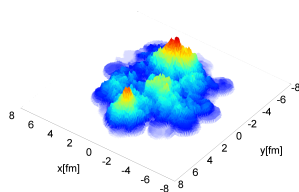
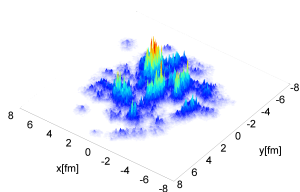
public code: T<sub>R</sub>ENTO

[Moreland, Bernhard, Bass (2014)]



**For the connoisseur:** IP-Glasma / KLN / EKRT / ...

(classical YM action in 2D, sat. scale  $Q_s$  ... valid at high- $E$ )



[Eskola, Kajantie, Ruuskanen, Tuominen (1999)]

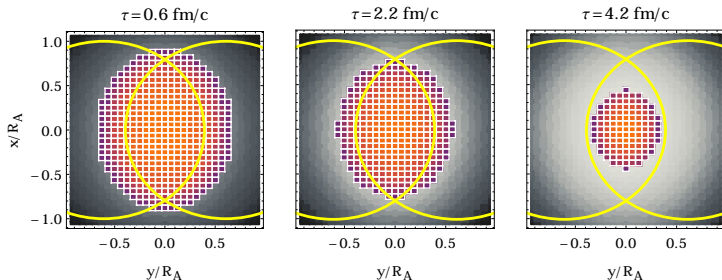
[Schenke, Tribedy, Venugopalan (2012)]

# Middle: (hydrodynamical simulation)

from IC: get energy density  $e(x, y, \dots)$  at  $\tau_0$

[Kurkela, et al. (2019)]

... then discretize & evolve in spacetime:



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**VISH2+1** = **VI**scous **H**ydrodynamics in (2+1) dim. [Song, Heinz (2008)]  
(using SHASTA = SHarp And Smooth Transport Algorithm)

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**MUSIC** = **MUS**(cl) for **I**on **C**ollisions [Schenke, Jeon, Gale (2010)]  
(MUSCL = Monotonic Upstream-centered Schemes for Conservation Laws)



# End: (“particlization”)

convert  $T^{\mu\nu}(X)$  and  $J^\mu(X)$  into **hadrons**  
(in a way that conserves  $E$  and  $\mathbf{p}$ )

[Huovinen, Petersen (2012)]

[Cooper, Frye (1974)]

$$\text{freeze-out: } E \frac{dN}{d^3\mathbf{p}} = \int_{\Sigma} d\sigma_{\mu} P^{\mu} f_{B/F} \left( \frac{P \cdot u(X)}{T(X)} \right)$$

( for MUSIC, this is done with iS3D:

<https://github.com/derekeverett/iS3D> )

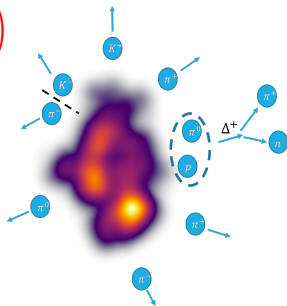
... then hadronic transport, e.g.:

*UrQMD* = **U**ltra-relativistic **Q**uantum **M**olecular **D**ynamics

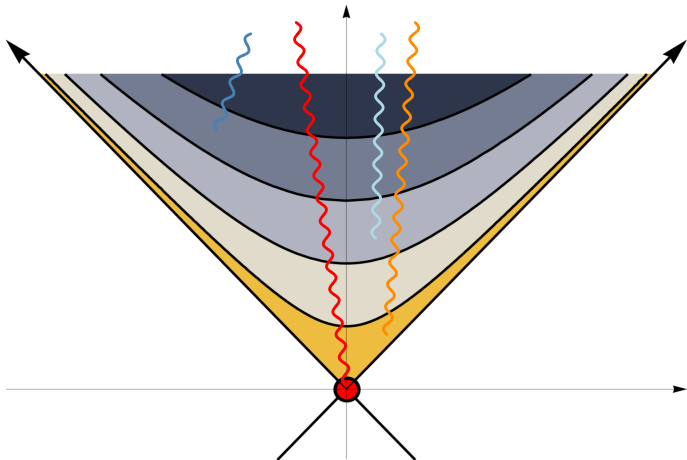
[Bleicher, et al. (1999)]

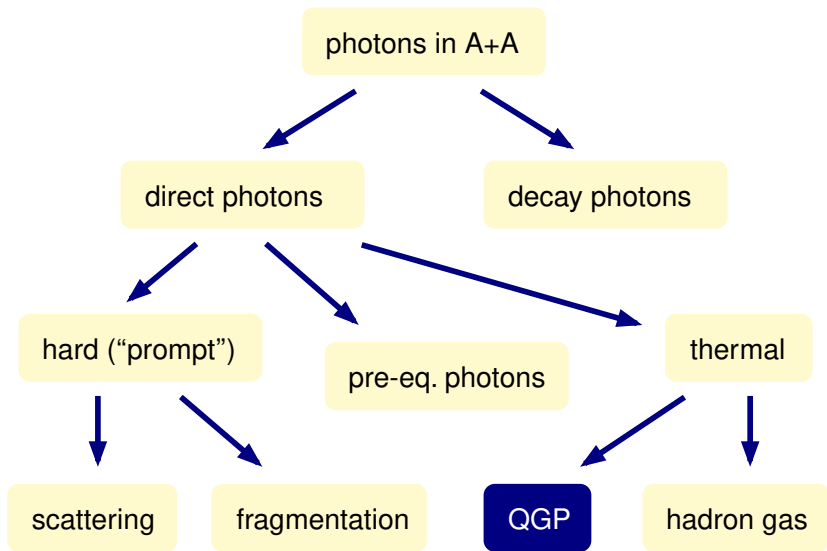
*SMASH* = **S**imulating **M**any **A**ccelerated **S**trongly interacting **H**adrons

[Weil, et al. (2016)]



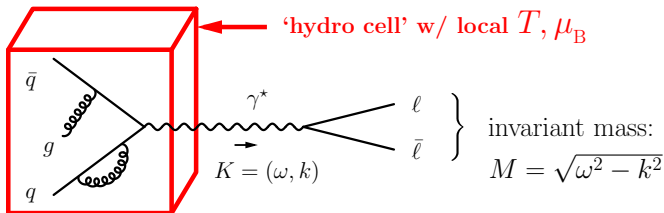
photons in A+A





## Basic relations from pert. theory

[McLerran, Toimela (1995)]  
[Gale, Kapusta (1991)]



Emission rate per unit volume,  $\Gamma_{\ell\bar{\ell}}$ , of an **equilibrated** QGP

$$\frac{d\Gamma_{\ell\bar{\ell}}}{d\omega d^3\mathbf{k}} = \frac{\alpha_{\text{em}}^2 \sum_{f=1}^{n_f} Q_f^2}{3\pi^3 M^2 (e^{\omega/T} - 1)} \times B\left(\frac{m_\ell^2}{M^2}\right) \times \rho_V(\omega, \mathbf{k})$$

- Quark charge-fractions:  $Q_f$  (in units of the electron charge)
- Kinematic factor:  $B(x) \equiv (1 + 2x)\Theta(1 - 4x)\sqrt{1 - 4x}$
- Spectral function  $\rho_V \equiv \rho_\mu^\mu$

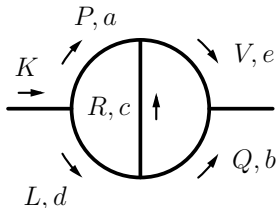
$$\rho_{\mu\nu}(\omega, \mathbf{k}) = \text{Im} \left[ \Pi_{\mu\nu}^{\text{ret}}(\omega + i0^+, \mathbf{k}) \right]$$

# QCD corrections

$$\Pi^{\mu\nu} = \left[ \sum_{l=0}^{\infty} g_s^{2l} \Pi_{(l)}^{\mu\nu} \right] + O(e^2); \quad \alpha_s = \frac{g_s^2}{4\pi}$$

$$= \text{---} \circ \text{---} + \text{---} \left( \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} \right) \text{---} + \text{---} \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right) \text{---} + \dots$$

Project 2-loop result onto 'basis' of **master diagrams** and evaluate:



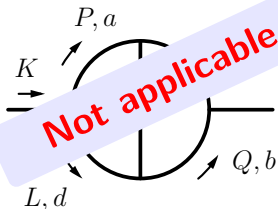
$$\rho_{abcde}^{(m,n)}(\omega, \mathbf{k}) \equiv \text{Im} \oint_{P,Q} \frac{p_0^m q_0^n}{P^{2a} Q^{2b} (K-P-Q)^{2c} (K-P)^{2d} (K-Q)^{2e}}$$

# QCD corrections

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$$= \text{tree} + \text{1-loop} + \text{2-loop} + \dots$$

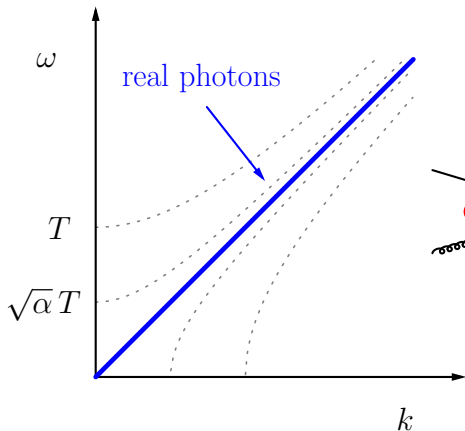
Project 2-loop result onto 'basis' of **master diagrams** to evaluate:



**Not applicable for  $K^2 \rightarrow 0!$**



$$\rho_{abcde}^{(m,n)}(\omega, \mathbf{k}) \equiv \text{Im} \oint_{P,Q} \frac{p_0^m q_0^n}{P^{2a} Q^{2b} (K-P-Q)^{2c} (K-P)^{2d} (K-Q)^{2e}}$$



## Thermal Screening

$$\frac{d\sigma}{dt} = \frac{-\pi \alpha_{em} \alpha_s}{3s^2} \frac{t^2 + s^2}{ts}$$



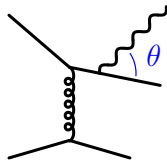
$$\sim \alpha_s T^2$$

[ Kapusta, Lichard, Seibert (1991) ]

[ Baier, et al (1992) ]

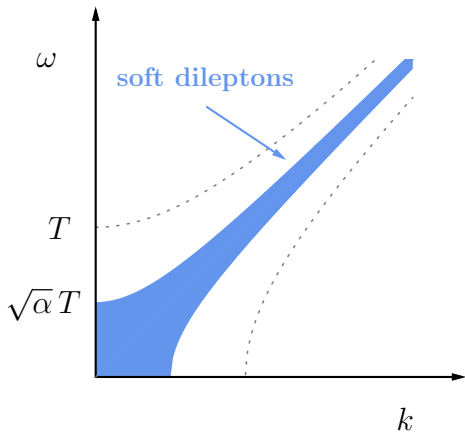
## Landau-Pomeranchuk-Migdal (LPM)

$$\int \frac{d \cos \theta}{E(1 - \cos \theta)} = \infty$$



LO: [ Arnold, Moore, Yaffe (2001) ] ,

NLO: [ Ghiglieri, et al (2013) ]



LO: [ Aurenche, et al (2002) ]

NLO: [ Ghiglieri, Moore (2014) ]

**light-like correlator**

↓ [Caron-Huot (2009)]

Effective  
Field Theory

'ladder diagrams' for  $M^2 \ll T^2 \rightarrow$  LPM effect + *Hard Thermal Loops*

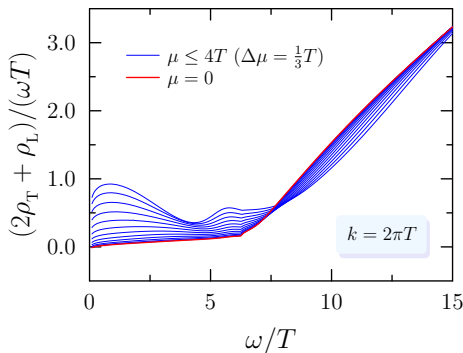
$$\rho_{\mu\nu}(\omega, \mathbf{k}) = \text{Im} \left[ \mu \text{---} \left( \text{---} \text{---} \text{---} \text{---} \text{---} \right) \text{---} \nu \right]$$



Combine LPM effect with strict 2-loop truncation,

[Ghisoiu, Laine (2014)]

$$\begin{aligned} \rho(\omega, k) \Big|_{\text{resummed}}^{\text{NLO}} &= \rho \Big|^{1\text{-loop}} + \rho \Big|^{2\text{-loop}} + \rho \Big|_{\text{beyond } O(g^2)}^{\text{LPM}} \\ &= \text{diagram 1} + \left( \text{diagram 2} + \text{diagram 3} \right) + \text{diagram 4} \end{aligned}$$



**NB:** Spectral fncs. can be checked with Euclidean corr. computable on the lattice:

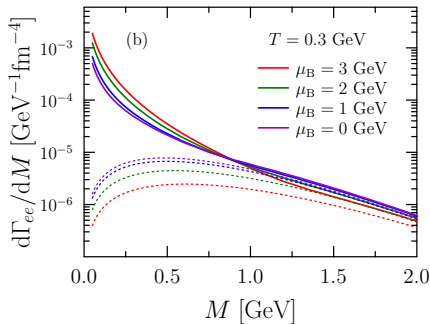
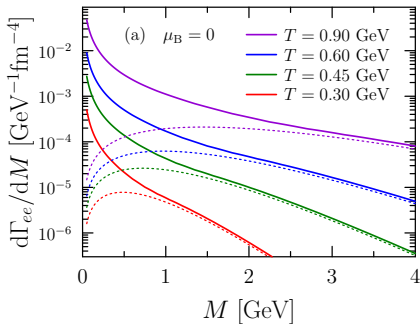
$$G(\tau) = \int_0^\infty d\omega \rho(\omega) \mathcal{K}(\omega, \tau)$$

⇒ see [GJ, Laine (2019)]  
and [Ali, et al. (2024)]

(  $\rho_V$  determined for  $\omega > k$  in [Laine (2013)] , and  $\omega < k$  in [GJ (2019)] )

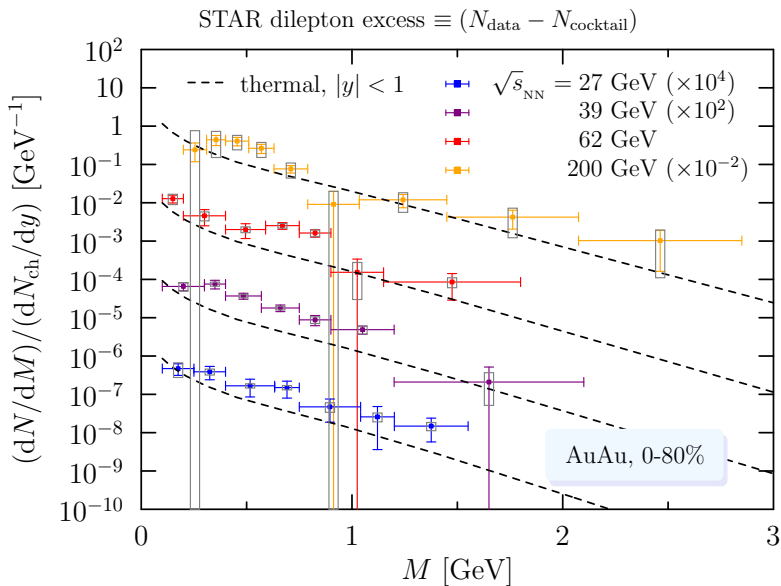
rate from a **static source**:

$$\frac{d\Gamma_{\ell\bar{\ell}}}{dM} = \int_{\mathbf{k}} \frac{M}{\sqrt{M^2 + k^2}} \frac{d\Gamma_{\ell\bar{\ell}}}{d\omega d^3\mathbf{k}}$$



- dileptons are a good thermometer!
- ... but a poor “baryometer”

\* in these, and subsequent, plots:  $\alpha_s = 0.3$



\*see also: [\[Burnier, Gastaldi \(2015\)\]](#) (LHC energies)

for large  $M \gg T$  and  $\mu_B$ :

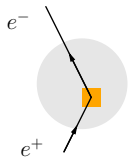
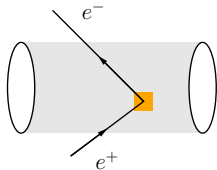
$$\frac{d\Gamma_{e\bar{e}}}{dM} \propto (MT_{\text{eff}})^{3/2} \exp(-M/T_{\text{eff}})$$

$\Rightarrow$  determine  $T_{\text{eff}}$  from the 'inverse slope' of the spectrum

What physics does this *effective* temperature represent?

in simulations we can access the full history, so the method can be tested!

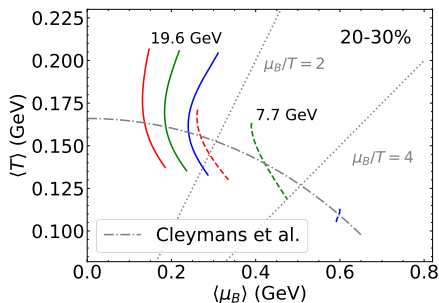
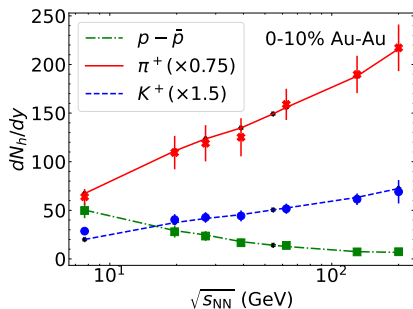
$$\frac{dN_{e\bar{e}}}{d^4K'} = \int dt \int d^3\mathbf{x} \left. \frac{d\Gamma_{e\bar{e}}}{d\omega d^3\mathbf{k}} \right|_{K^\mu = \Lambda^{\mu\nu} K'_\nu}$$



**note:**  $\rho_V(\omega, k)$  evaluated at  $\omega = K'_\mu u^\mu$ ,  $k = \sqrt{(K'_\mu u^\mu)^2 - M^2}$

# Calibrating the thermometer

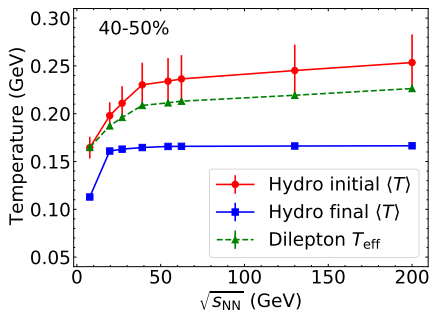
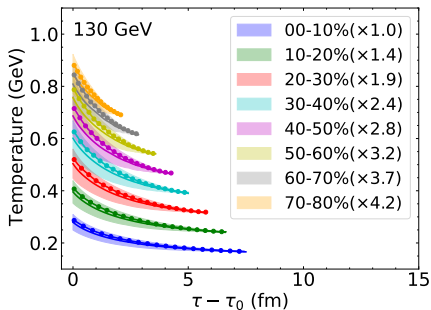
- MC-Glauber initial conditions (at finite  $\mu_B$ )
- Hydro with MUSIC (including viscous corrections)
- Equation of state: NEOS-B (neglects strangeness and  $\mu_e$ )
- Hydro stops at  $e_{fo} = 0.26 \text{ GeV}/\text{fm}^3$
- Freeze-out (iS3D) and hadronic scatterings w/ UrQMD



code public: <https://github.com/LipeiDu/DileptonEmission>

# Calibrating the thermometer

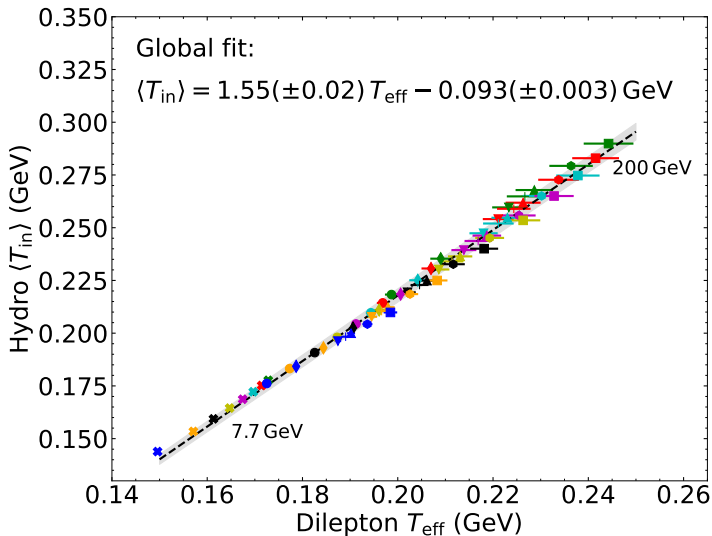
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$T_{\text{eff}}$  represents the *initial* temperature!

[Churchill, et al. (2024)]



... predicted over 40 years ago (!!)

[Kajantie, Miettinen (1981)]

# Summary

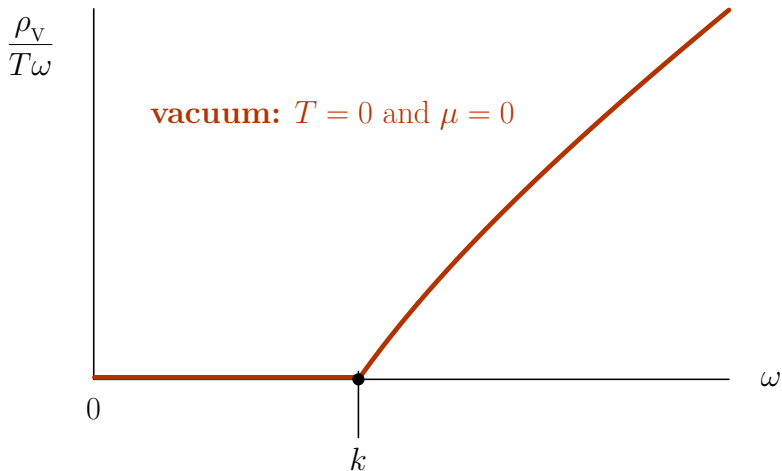
Arxiv: 2211.09575  
2311.06675  
2311.06951

- thermal dilepton yields at NLO+LPM  
⇒ predicted from first principles, at finite  $T$  and  $\mu_B$
- extracted 'effective' temperature  
⇒ linear relationship between  $T_{\text{eff}}$  and  $T_{\text{in}}$

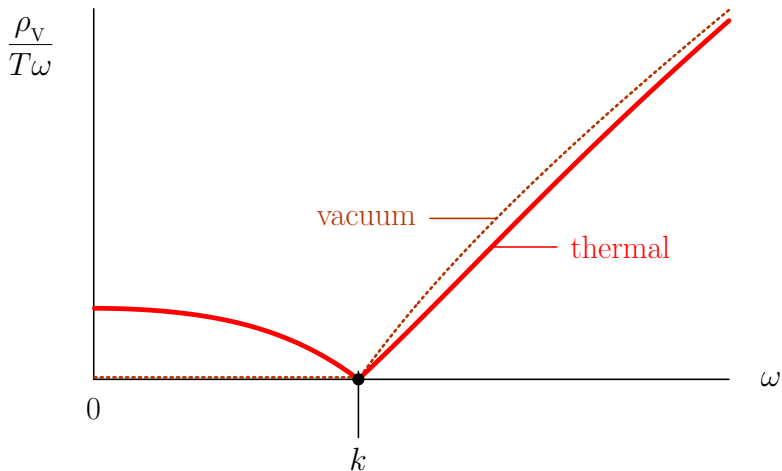




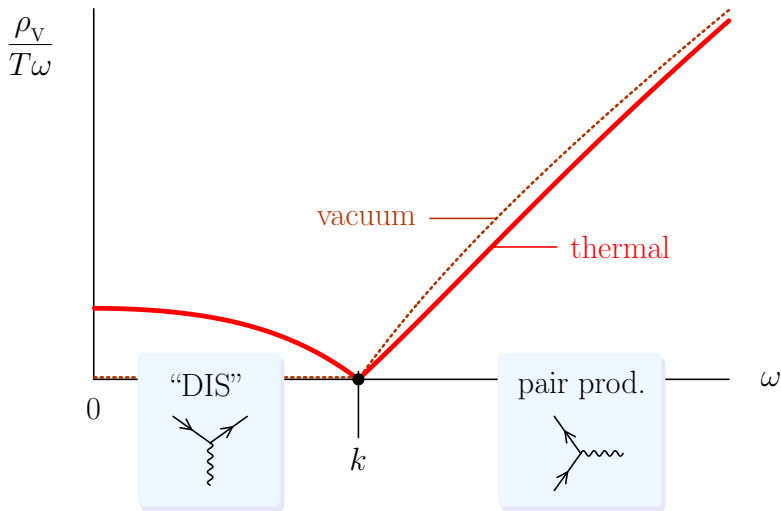
$$\text{Im} \left[ \text{Diagram} \right] \xrightarrow{T=0} \frac{NK^2}{4\pi} \Theta(K^2)$$



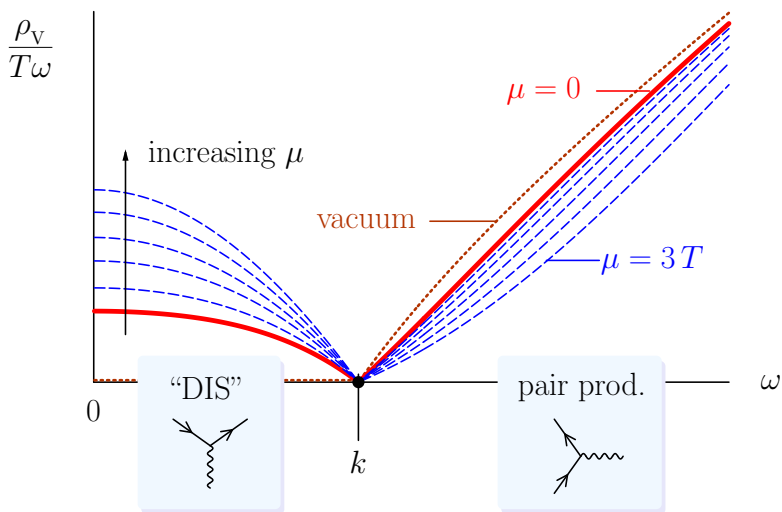
$$\text{Im} \left[ \text{Diagram} \right] \xrightarrow{T>0} \frac{NK^2}{4\pi} \left\{ \frac{2T}{k} \log \left[ \frac{1 + e^{-\frac{1}{2}(\omega+k)/T}}{1 + e^{-\frac{1}{2}|\omega-k|/T}} \right] + \Theta(K^2) \right\}$$

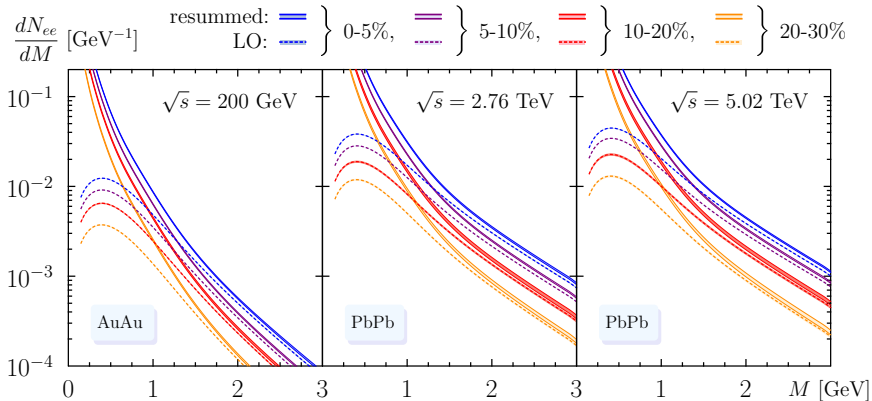


$$\text{Im} \left[ \text{Diagram} \right] \xrightarrow[\mu=0]{T>0} \frac{NK^2}{4\pi} \left\{ \frac{2T}{k} \log \left[ \frac{1 + e^{-\frac{1}{2}(\omega+k)/T}}{1 + e^{-\frac{1}{2}|\omega-k|/T}} \right] + \Theta(K^2) \right\}$$



$$\text{Im} \left[ \text{Diagram} \right] = \frac{NK^2}{4\pi} \left\{ \frac{T}{k} \sum_{\nu=\pm\mu} \log \left[ \frac{1 + e^{(\nu - \frac{1}{2}(\omega+k))/T}}{1 + e^{(\nu - \frac{1}{2}|\omega-k|)/T}} \right] + \Theta(K^2) \right\}$$





classical YM action in 2D

→ saturation scale  $Q_s$

IP-Glasma w/  $b = 0 \dots 20$  fm

thermal dileptons during *local* equilib

→ specified by viscous EM tensor...

$$T^{\mu\nu} = \varepsilon u^\mu u^\nu - (g^{\mu\nu} - u^\mu u^\nu)(p + \Pi) + \pi^{\mu\nu}$$

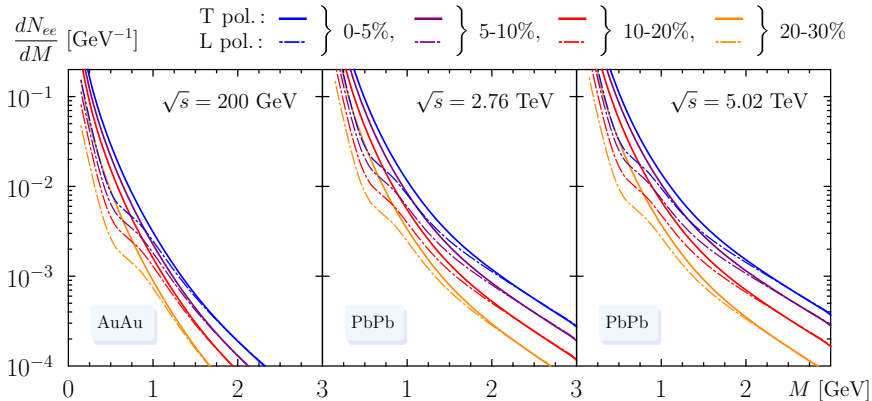
hydro: MUSIC 2 + 1D

Cooper-Frye (iS3D) → UrQMD

CGC ●

$\tau_0 \simeq 0.4$  fm

$\tau_{\text{freezeout}} @ T = 145$  MeV



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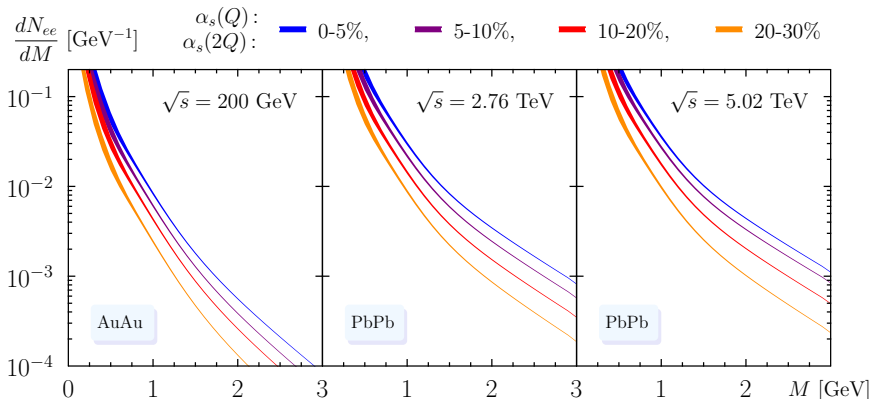
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# Considerations for non-zero $\mu_B$

- chemical equilibrium  $\Rightarrow \mu \equiv \mu_q = \frac{1}{3}\mu_B$
- Debye mass  $m_D$  and the 'asymptotic' quark mass  $m_\infty$

$$m_D^2 \equiv g^2 \left[ \left( \frac{1}{2} n_f + N \right) \frac{T^2}{3} + n_f \frac{\mu^2}{2\pi^2} \right]$$

$$m_\infty^2 \equiv g^2 \frac{C_F}{4} \left( T^2 + \frac{\mu^2}{\pi^2} \right)$$



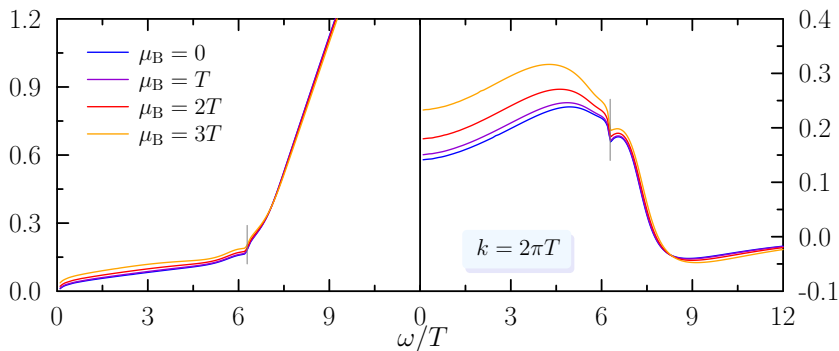
- large frequency limit: enhancement ↘

$$\rho_V \simeq \frac{NM^2}{4\pi} + 4g^2 C_F N \left\{ \frac{3M^2}{4(4\pi)^3} + \frac{\pi(\omega^2 + \frac{k^2}{3})}{36M^4} \left( T^4 + \frac{6}{\pi^2} T^2 \mu^2 + \frac{3}{\pi^4} \mu^4 \right) \right\}$$

**NEW RESULTS:** the **full** effect of  $\mu_B$  on  $\rho(\omega, k) \Big|_{\text{resummed}}^{\text{NLO}} \dots$

# What we find

pQCD spectral functions  $T = 280$  MeV and  $n_f = 3$ :



Spectral functions,

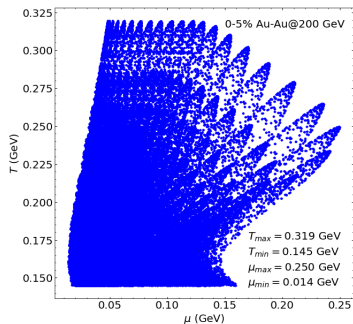
$$\text{Left: } \rho_V/(\omega T) = (2\rho_T + \rho_L)/(\omega T)$$

$$\text{Right: } \rho_H/(\omega T) = 2(\rho_T - \rho_L)/(\omega T)$$

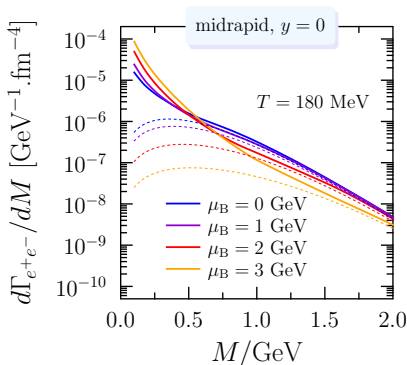
# Impact on yield (non-zero $\mu_B$ )

BES  $\Rightarrow$  probe baryon rich region work w/ Churchill, Du, Gale, Jeon

MUSIC: [Schenke, Jeon, Gale (2010)]



( $\uparrow$  figure by Lipai Du)

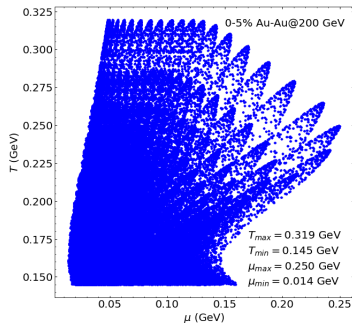


$\Rightarrow$  compensation of LO suppression & NLO enhancement! ...

# Impact on yield (non-zero $\mu_B$ )

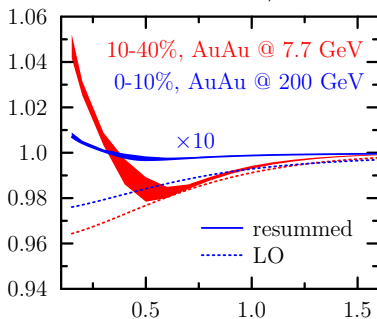
BES  $\Rightarrow$  probe baryon rich region work w/ Churchill, Du, Gale, Jeon

MUSIC: [Schenke, Jeon, Gale (2010)]



( $\uparrow$  figure by Lipei Du)

$$\text{ratio: } \frac{dN_{e^+e^-}/dM|_{\mu>0}}{dN_{e^+e^-}/dM|_{\mu=0}}$$



smooth MC-Glauber initial conditions + baryon diffusion